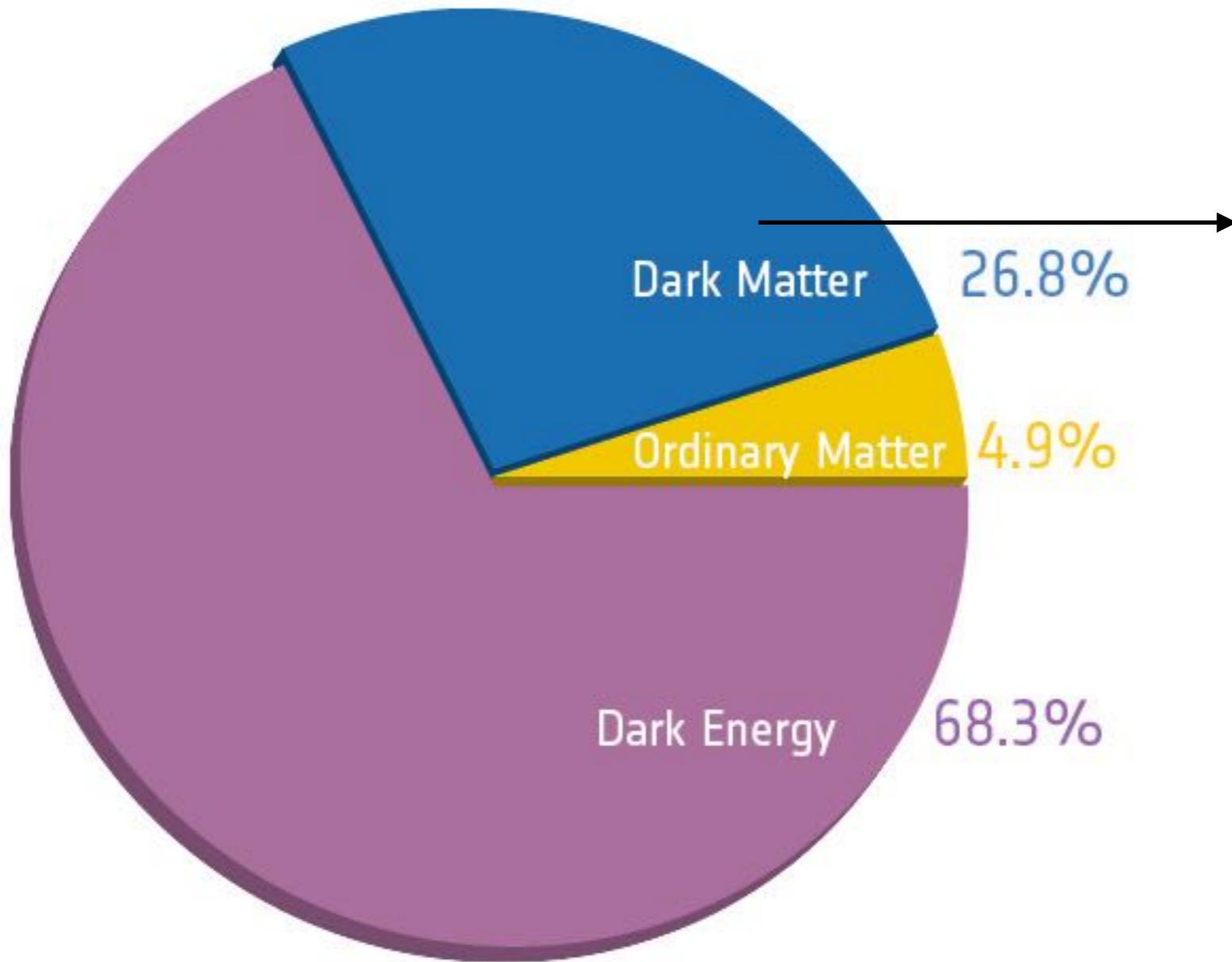


Distinguishing the spin of DM using dilepton events

Antonio Delgado
University of Notre Dame

- Introduction
- DM in the compressed region: challenges
- Kinematical variables
- Conclusions

Work done in collaboration with Rodolfo Capdevilla, Nirmal Raj and Adam Martin
[arXiv:1709.00439](https://arxiv.org/abs/1709.00439)



What is this?

- Is DM its own antiparticle?
- Does it carry spin?
- What is its mass?
- How does it couple to the SM if at all?
- I am going to present a way of ‘partially’ answering this questions using LHC data with fully reconstructed final states.

- Lots of models of DM require extra degrees of freedom apart from the neutral state itself.
- That is needed in order to have renormalizable couplings with SM (with the exception of the Higgs portal)
- I am going to assume a simplified model including our DM and two messengers

$$\mathcal{L} \supset -\sqrt{2}(\lambda_{\tilde{Q}} \tilde{Q} \chi_B^\dagger q^\dagger + \lambda_{\tilde{L}} \tilde{L} \chi_B^\dagger \ell^\dagger) + \text{H.c.}$$

- I am going to assume that DM is either spin 1/2 or 0 and the messenger therefore will have the other spin to be able to write the previous coupling
- The mass terms can be written as (δm is needed to avoid problems with large DD cross-sections)

$$\mathcal{L}_{\text{mass}} = (\chi_A \ \chi_B) \begin{pmatrix} \delta m & m_\chi \\ m_\chi & \delta m' \end{pmatrix} \begin{pmatrix} \chi_A \\ \chi_B \end{pmatrix} + \text{H.c.}$$

Spin 1/2

$$\mathcal{L}_{\text{mass}} = \frac{1}{2} (\phi_\chi \ \phi_\chi^\dagger) \begin{pmatrix} \delta m^2 & m_\chi^2 \\ m_\chi^2 & \delta m'^2 \end{pmatrix} \begin{pmatrix} \phi_\chi \\ \phi_\chi^\dagger \end{pmatrix}$$

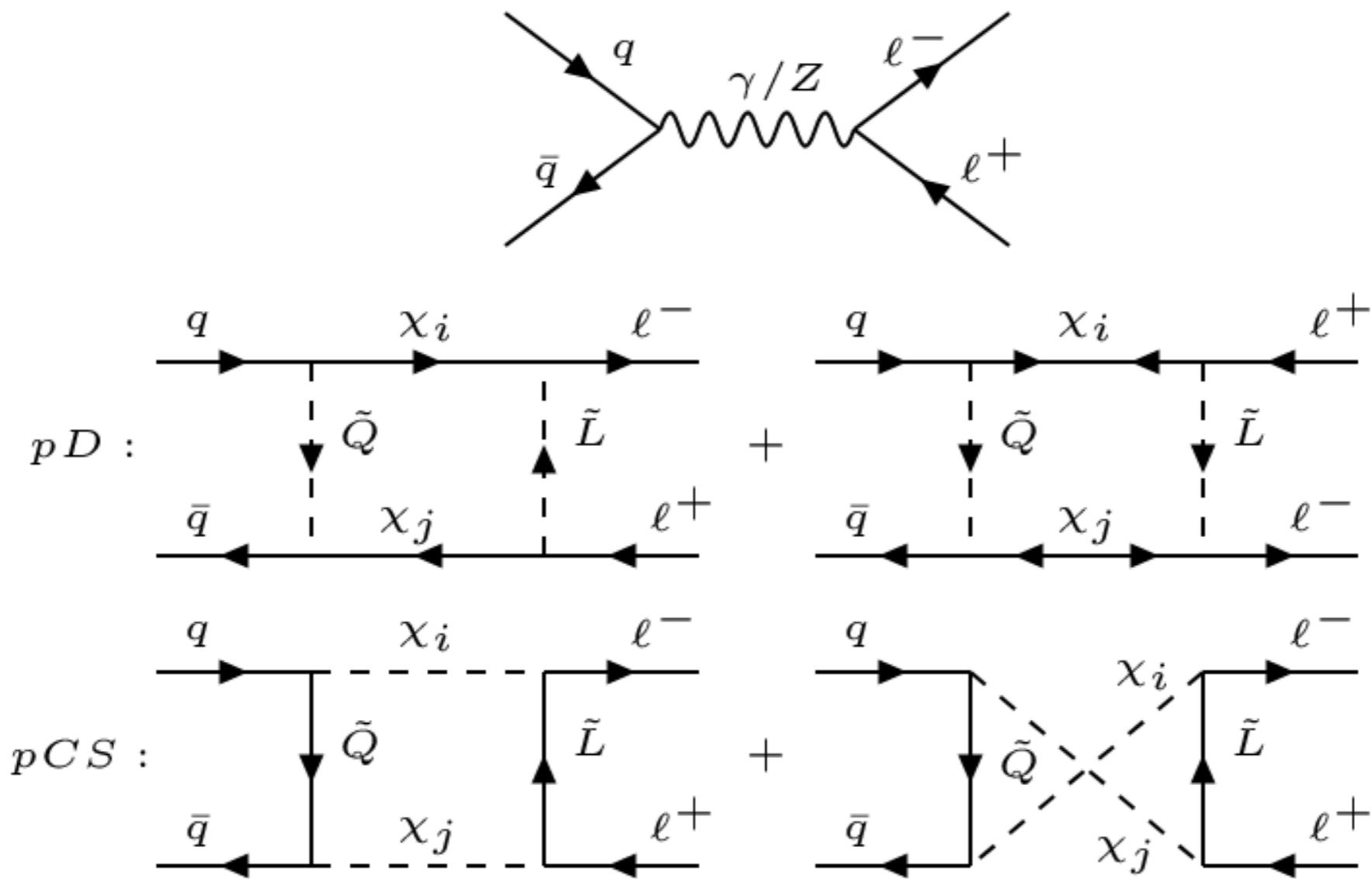
Spin 0

- Varying the spin and quantum numbers of the messengers one can classify the different scenarios that are going to be analyzed:

Model	χ spin	\tilde{Q}, \tilde{L} spin	\tilde{Q} under G_{SM}	\tilde{L} under G_{SM}
pD_{RR}^u	1/2	0	(3, 1, 2/3)	(1, 1, -1)
pD_{RL}^u	1/2	0	(3, 1, 2/3)	(1, 2, -1/2)
pCS_{RR}^u	0	1/2	(3, 1, 2/3)	(1, 1, -1)
pCS_{RL}^u	0	1/2	(3, 1, 2/3)	(1, 2, -1/2)
pD_{RR}^d	1/2	0	(3, 1, -1/3)	(1, 1, -1)
pD_{RL}^d	1/2	0	(3, 1, -1/3)	(1, 2, -1/2)
pCS_{RR}^d	0	1/2	(3, 1, -1/3)	(1, 1, -1)
pCS_{RL}^d	0	1/2	(3, 1, -1/3)	(1, 2, -1/2)

- The usual strategy to discover a model like the one proposed would be to produce the **color** companion and use the standard search of jets+MET
- But in the compressed case, needed in some scenarios to reproduce the **correct relic abundance**, the amount of MET may be small so the search may not be completely effective.
- Even if one can discover the messenger using this channel, it is difficult to extract information on the nature of DM since the final state is not fully reconstructed.

- In the paper we propose to use the following alternative fully reconstructed signal:

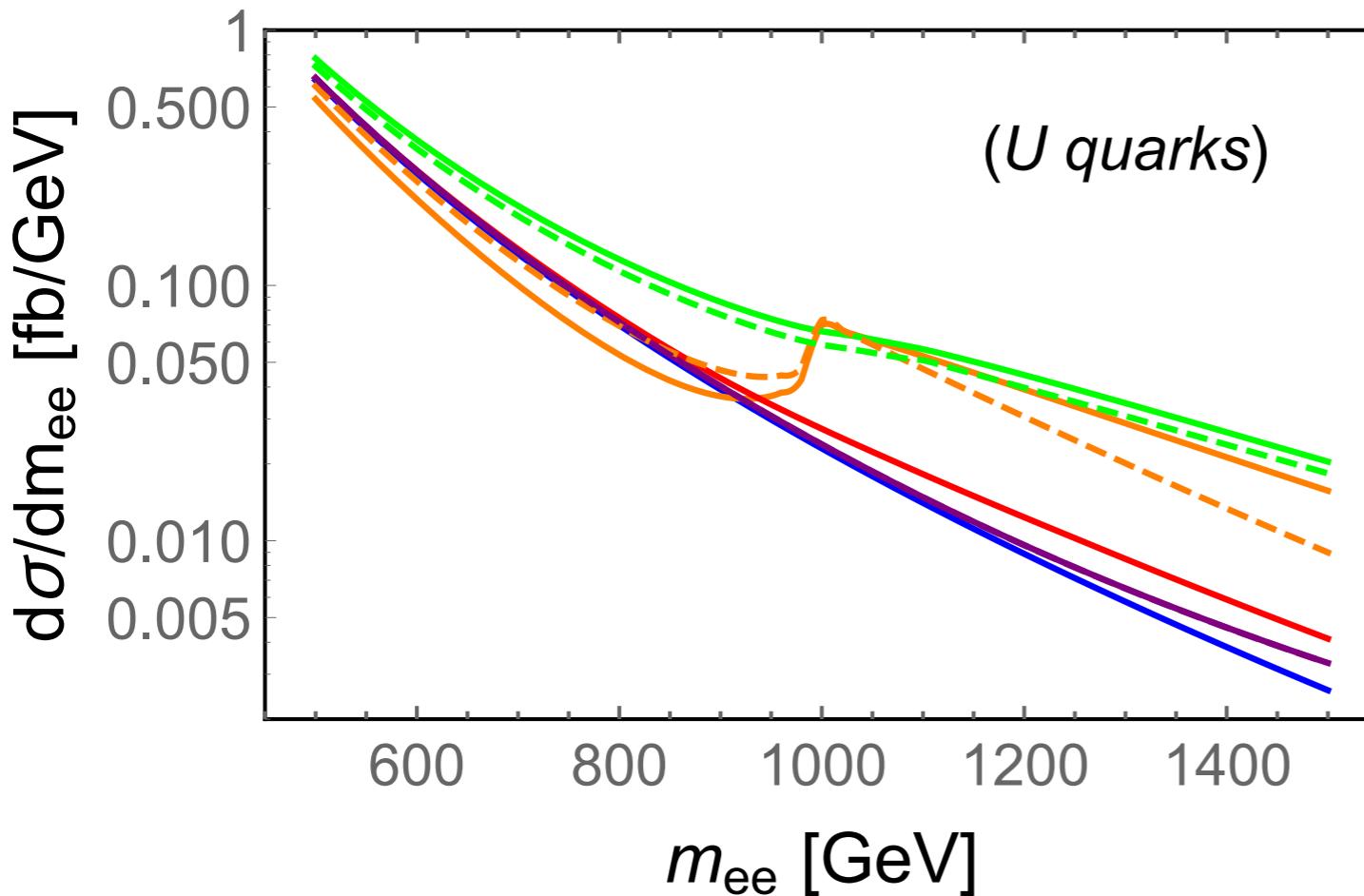


- We first calculate analytically the LO cross-section for both SM and new physics:

$$\begin{aligned} d\sigma_{\text{tot}} &\equiv \frac{d^2\sigma_{\text{tot}}}{d\cos\theta \ dm_{\ell\ell}} \\ &= d\sigma_{\text{SM}} + d\sigma_{\text{int}} + d\sigma_{\chi} \end{aligned}$$

$$\begin{aligned} d\sigma_{\text{SM}} &= \frac{1}{32\pi m_{\ell\ell}^2 N_c} \sum_{\text{spins}} |\mathcal{M}_{\text{SM}}|^2 , \\ d\sigma_{\text{int}} &= \frac{1}{32\pi m_{\ell\ell}^2 N_c} \sum_{\text{spins}} 2\text{Re}(\mathcal{M}_{\text{SM}} \mathcal{M}_{\chi}^*) , \\ d\sigma_{\chi} &= \frac{1}{32\pi m_{\ell\ell}^2 N_c} \sum_{\text{spins}} |\mathcal{M}_{\chi}|^2 , \end{aligned}$$

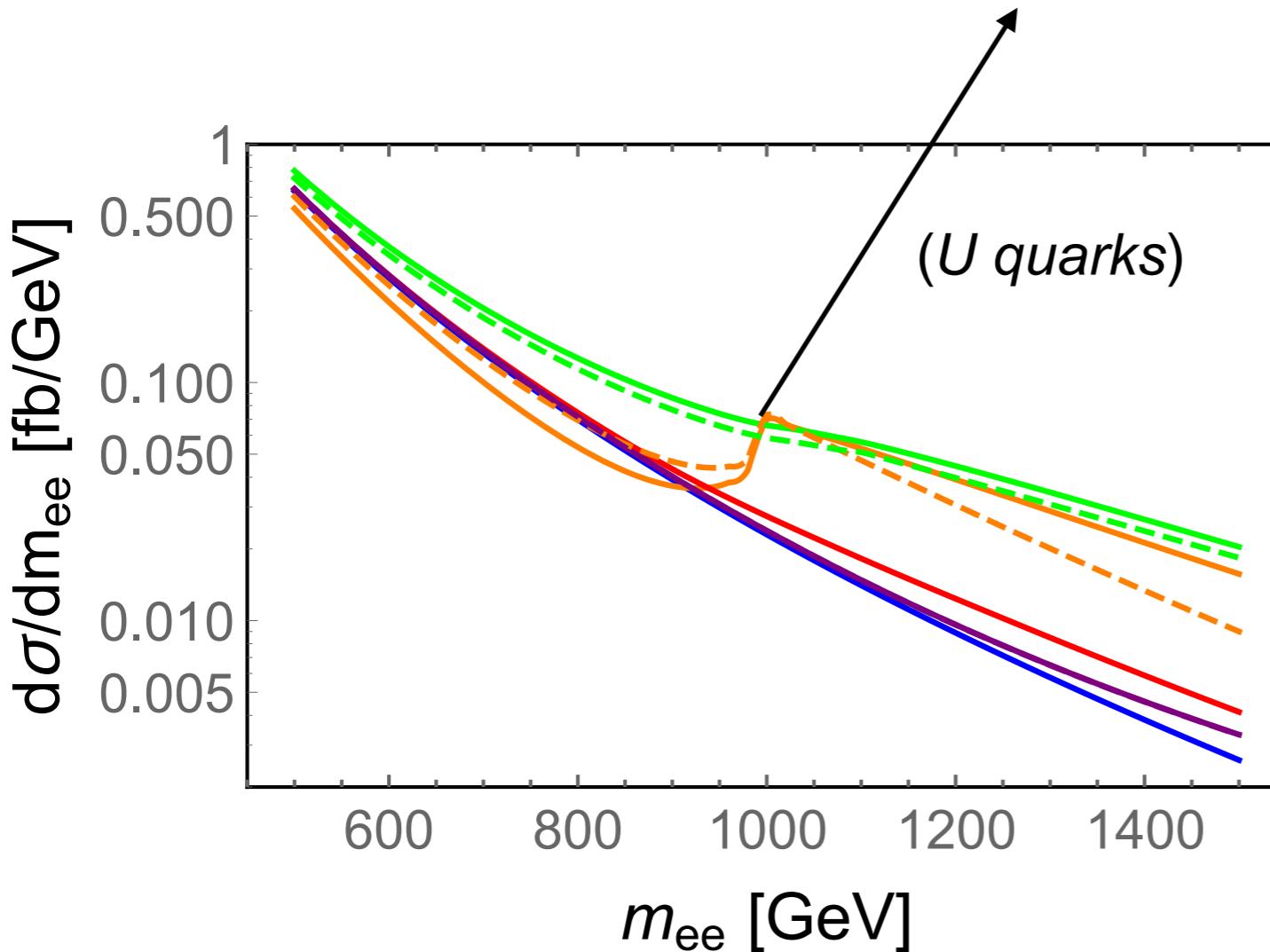
**To the σ_{SM} we add the NLO contribution from QCD
MCMF 8.0**



- Blue is the SM background
- Orange pseudo Dirac
- Green pseudo Complex
- Red is Majorana
- Purple real scalar
- Solid is RR
- Dashed is RL

$$\lambda = 2 \text{ m}_\chi = 500 \text{ GeV } m_\phi = 550 \text{ GeV}$$

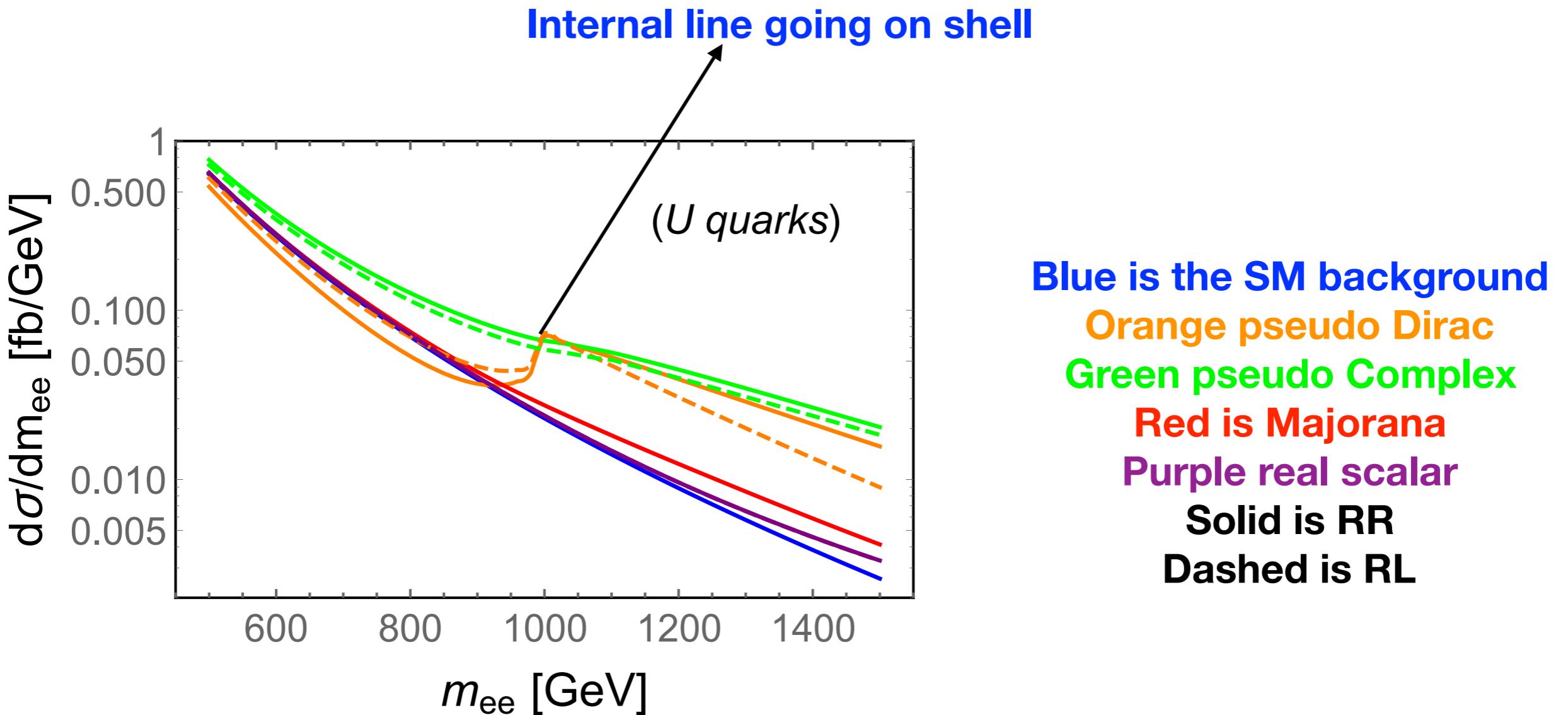
**Majorana fermion DM and Real scalar (self-conjugate)
do not give enough signal!!!!**



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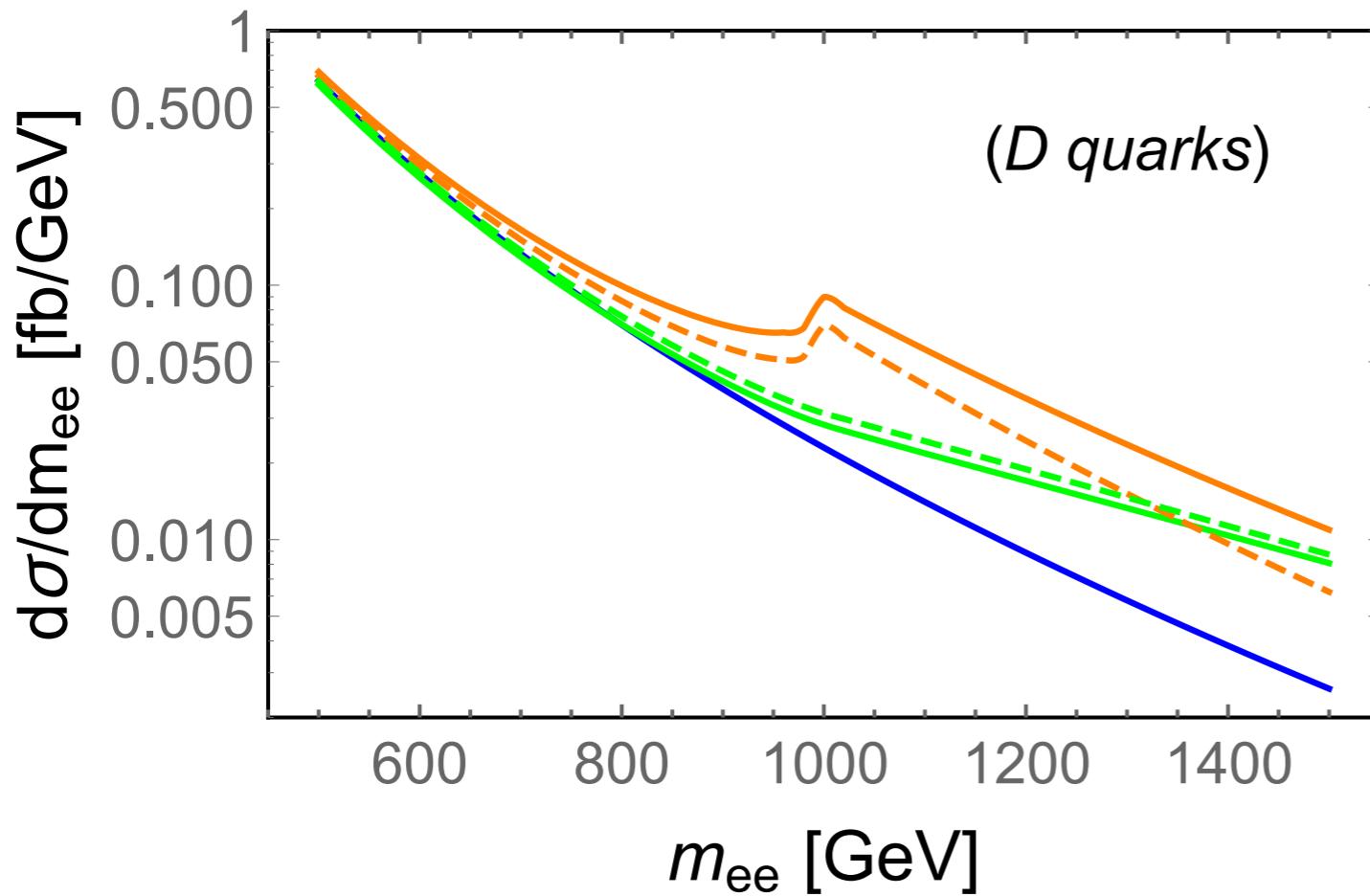
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$$\lambda = 2 \quad m_\chi = 500 \text{ GeV} \quad m_\phi = 550 \text{ GeV}$$

The differences between U and D quarks come from some interference terms

- We are going to use the following kinematical variables to discriminate the different models:

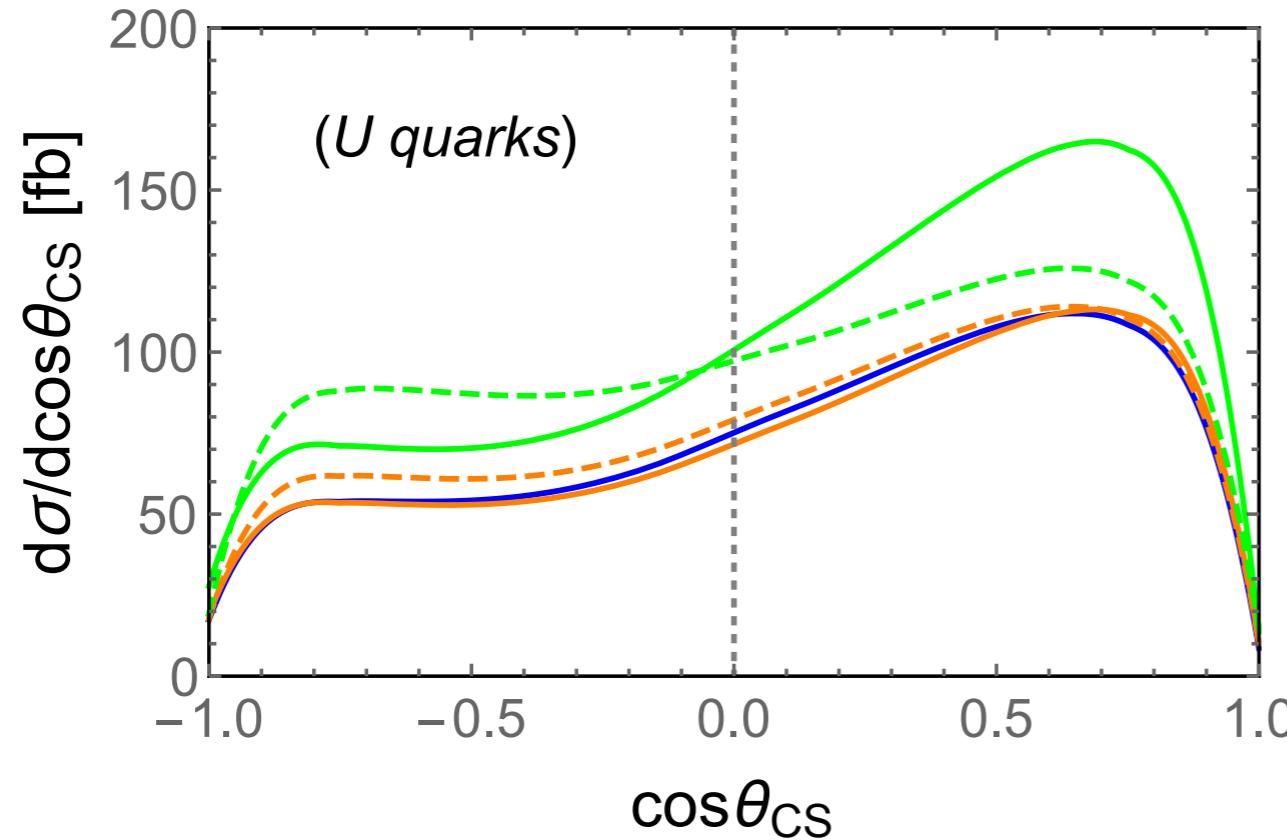
$$\cos \theta_{\text{CS}} = \frac{Q_z}{|Q_z|} \frac{2(p_1^+ p_2^- - p_1^- p_2^+)}{|Q| \sqrt{Q^2 + Q_T^2}}$$

**Q is the total dilepton momentum
p's are the light-cone momenta**

$$A_{\text{CE}} \equiv \frac{N(|\cos \theta| < \cos \theta_0) - N(|\cos \theta| > \cos \theta_0)}{N(|\cos \theta| < \cos \theta_0) + N(|\cos \theta| > \cos \theta_0)}$$

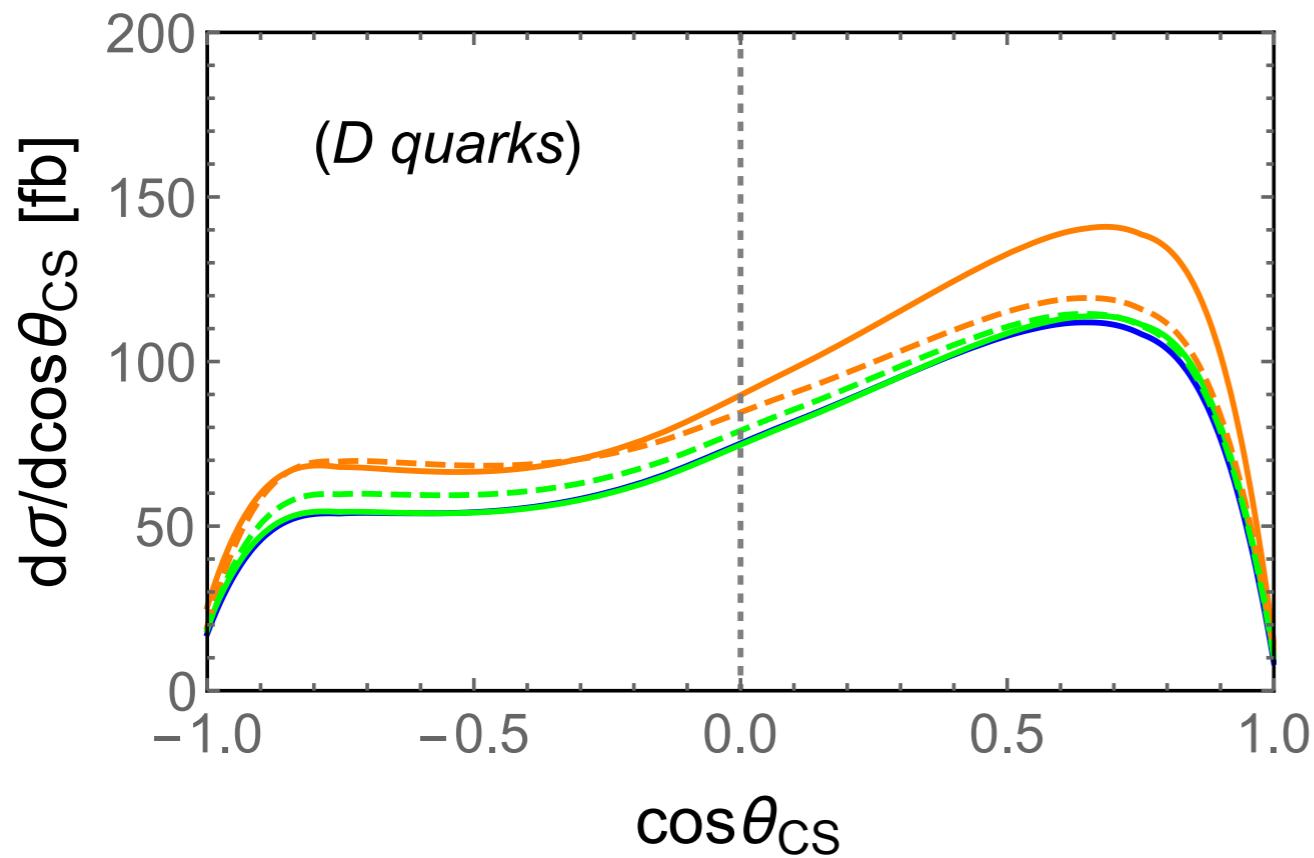
$$A_{\text{FB}} \equiv \frac{N(\cos \theta > 0) - N(\cos \theta < 0)}{N(\cos \theta > 0) + N(\cos \theta < 0)}$$

$$|\eta_{\ell^\pm}| \leq 2.4 , \quad p_T^{\ell^\pm} \geq 40 \text{ GeV}$$



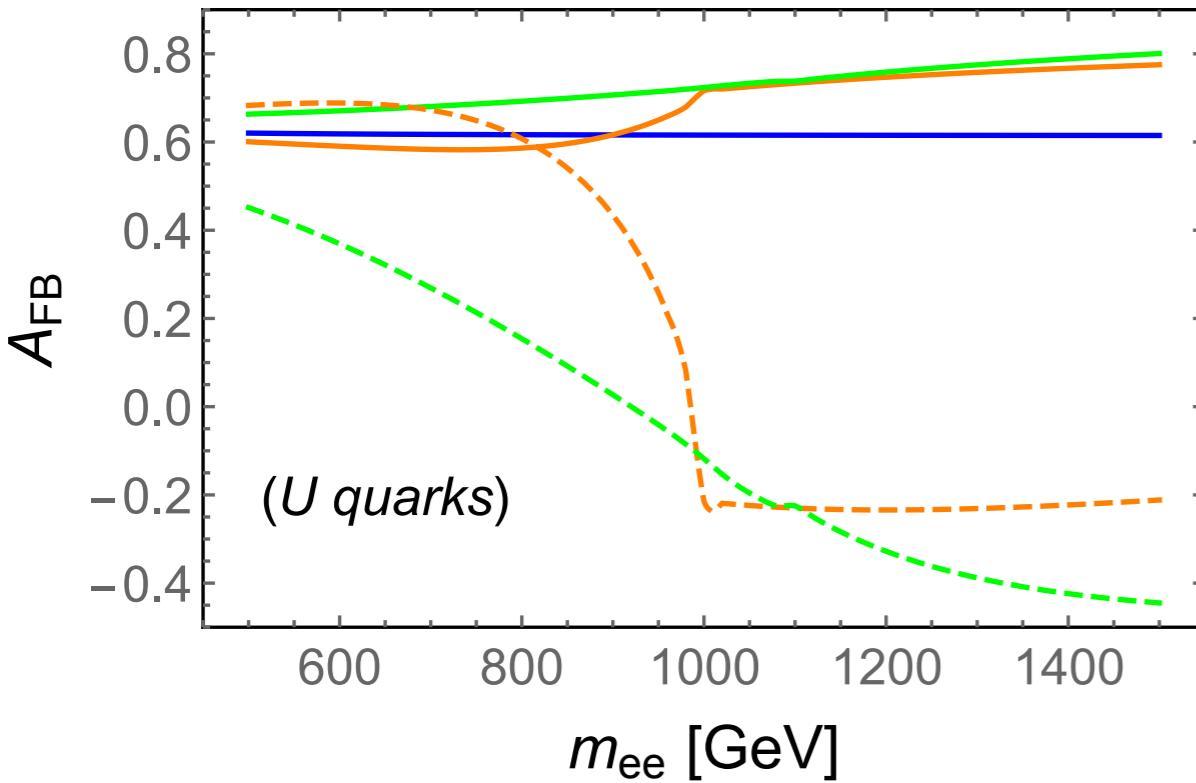
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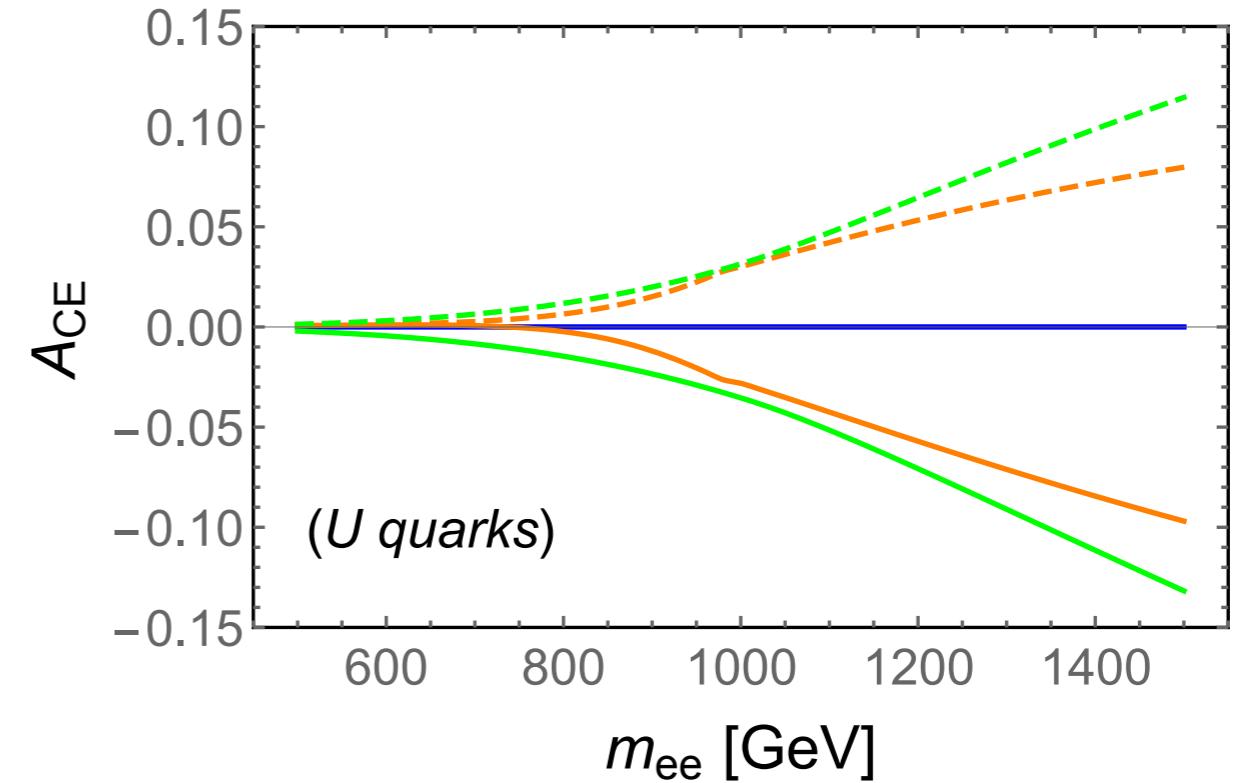


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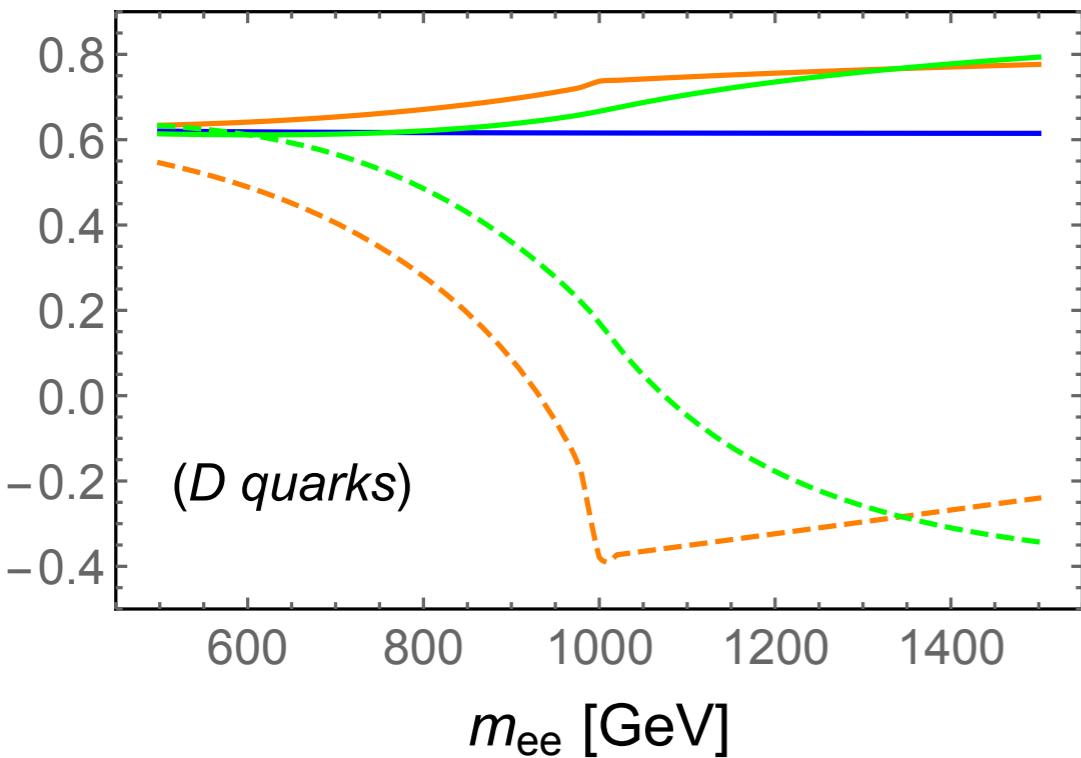
$\lambda = 2$ $m_\chi = 500$ GeV $m_\phi = 550$ GeV



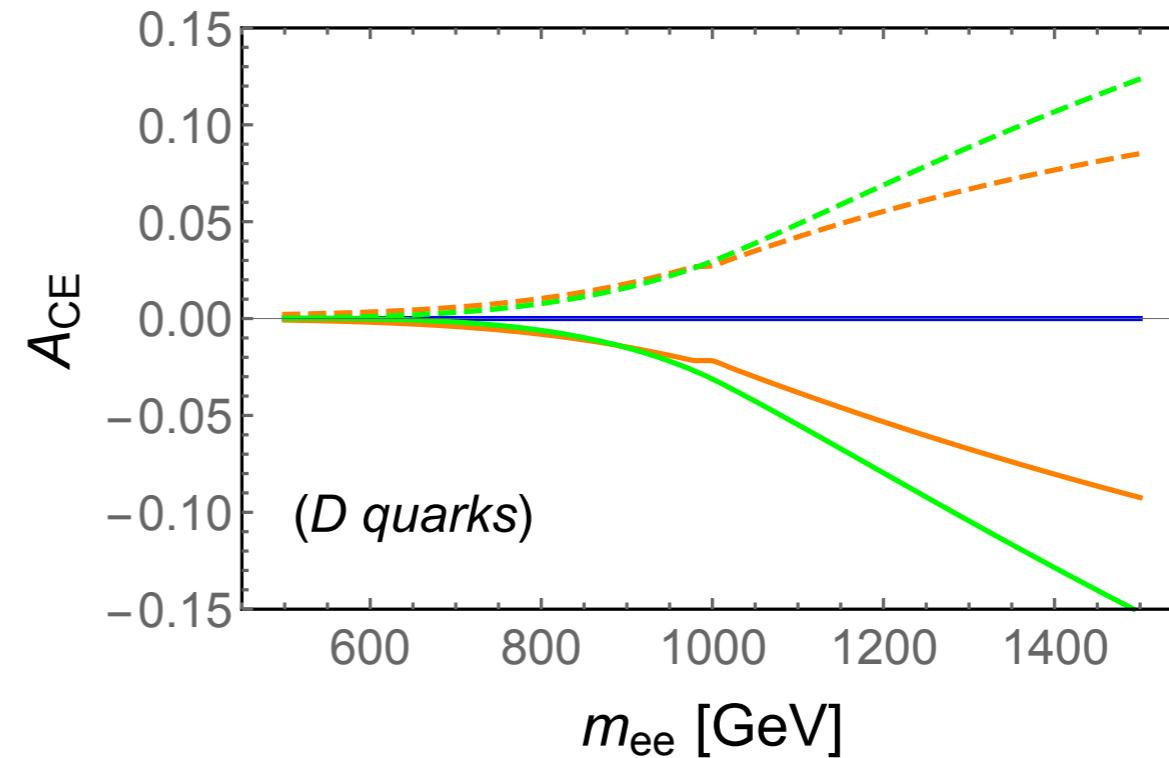
$\lambda = 2$ $m_x = 500$ GeV $m_\phi = 550$ GeV



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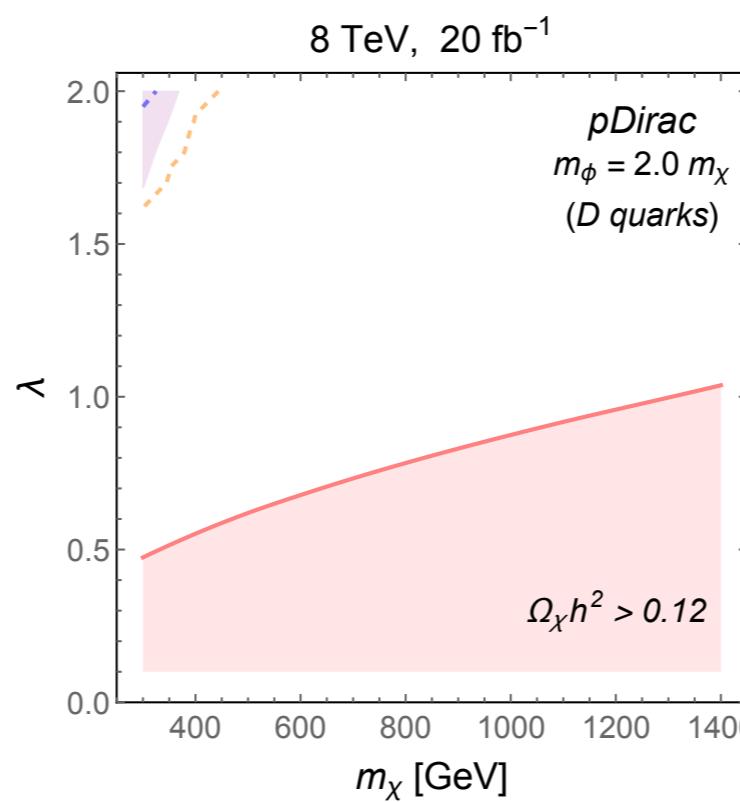
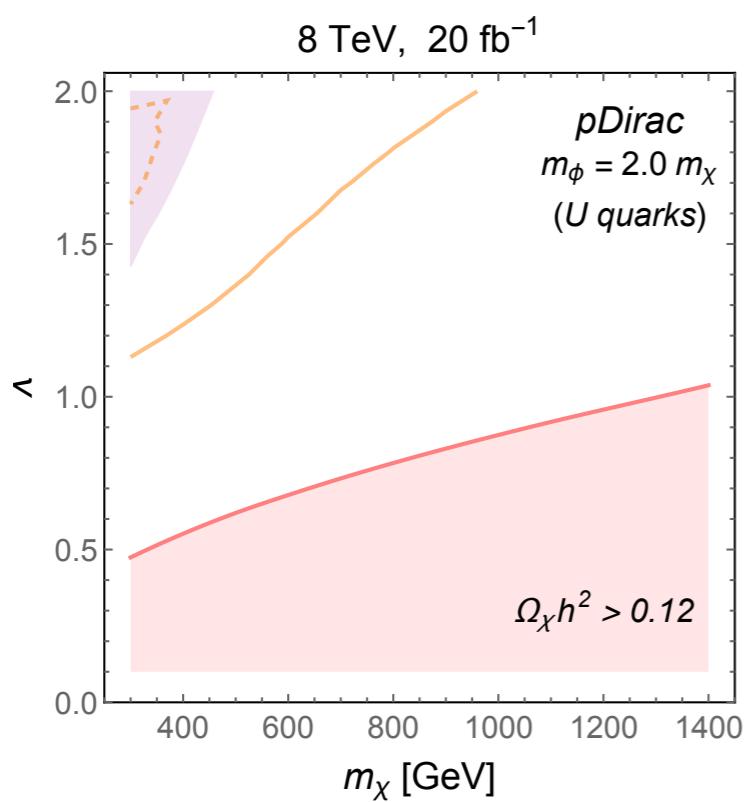
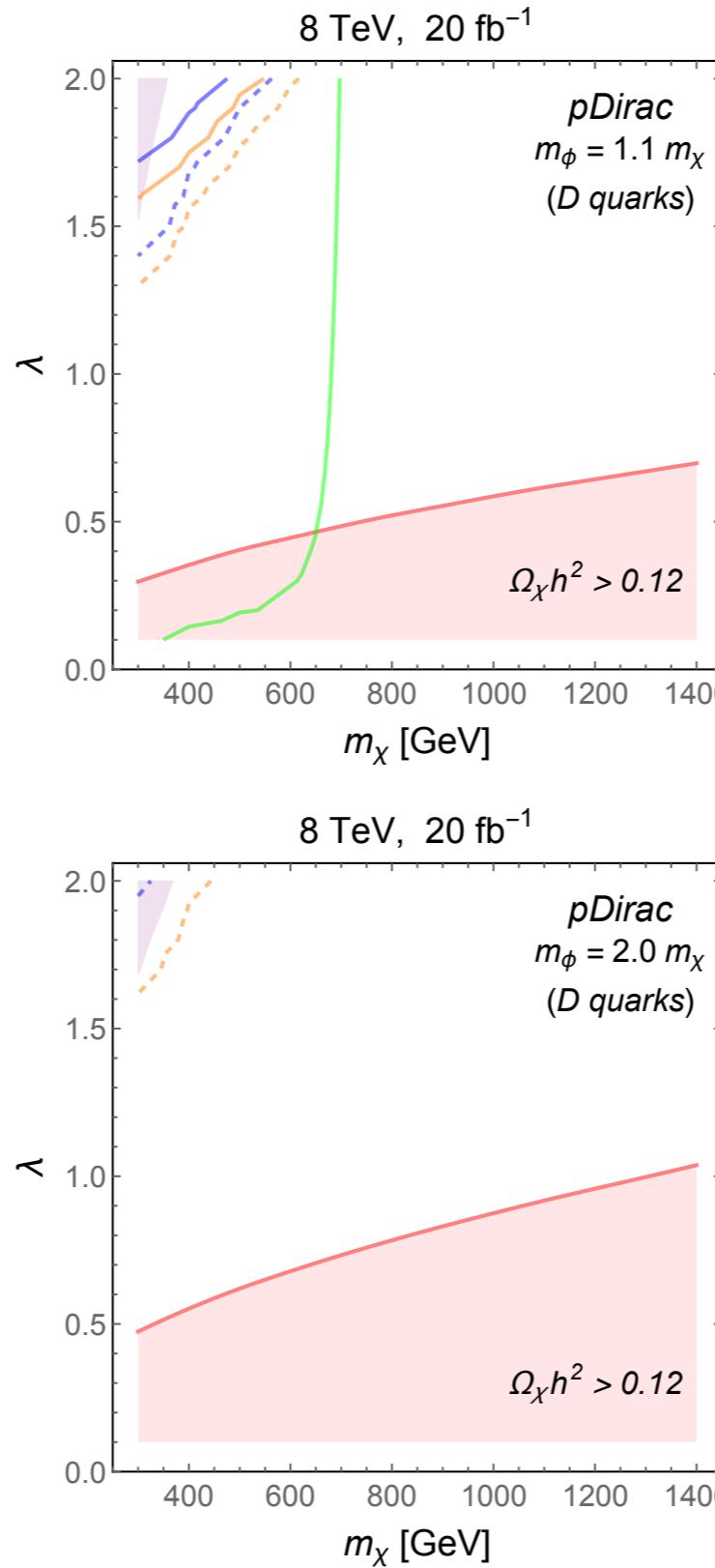
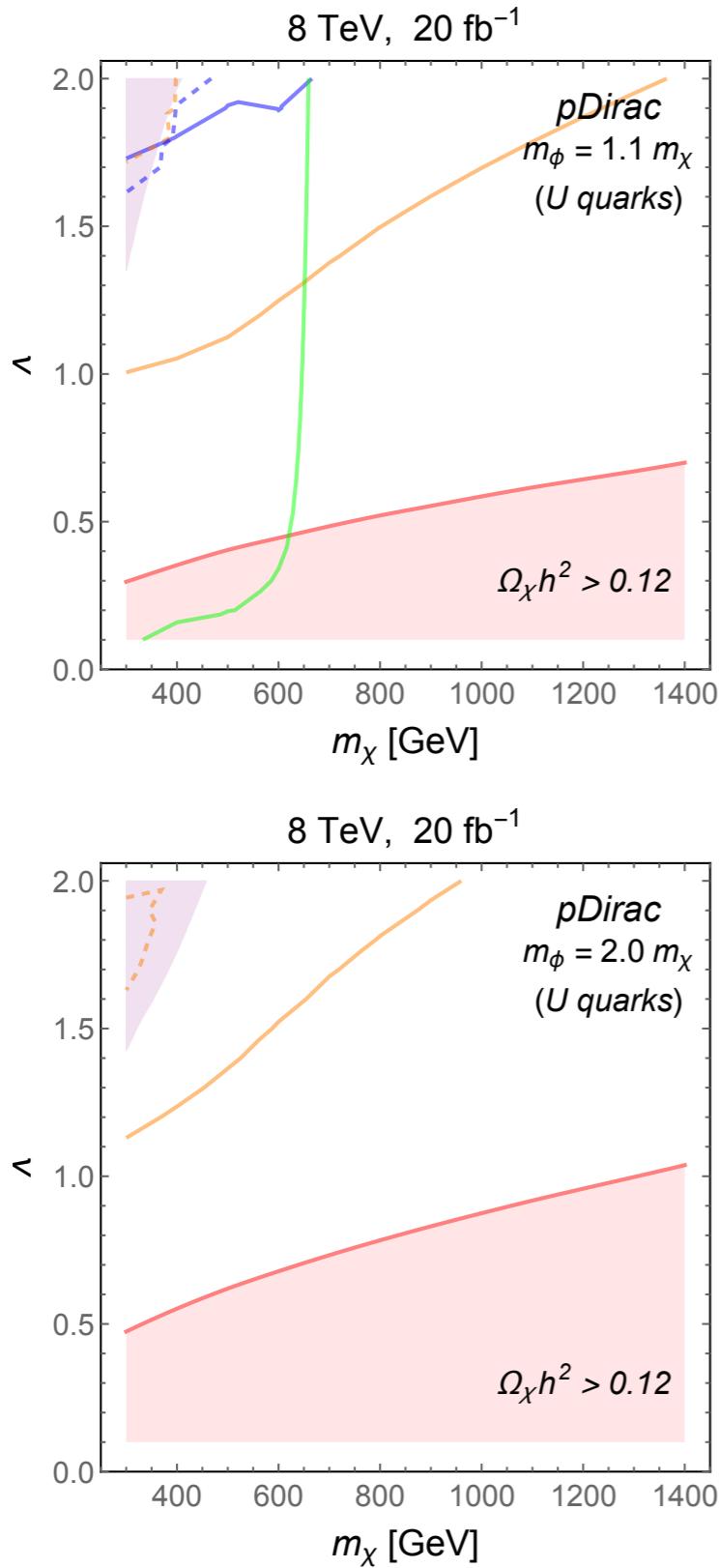


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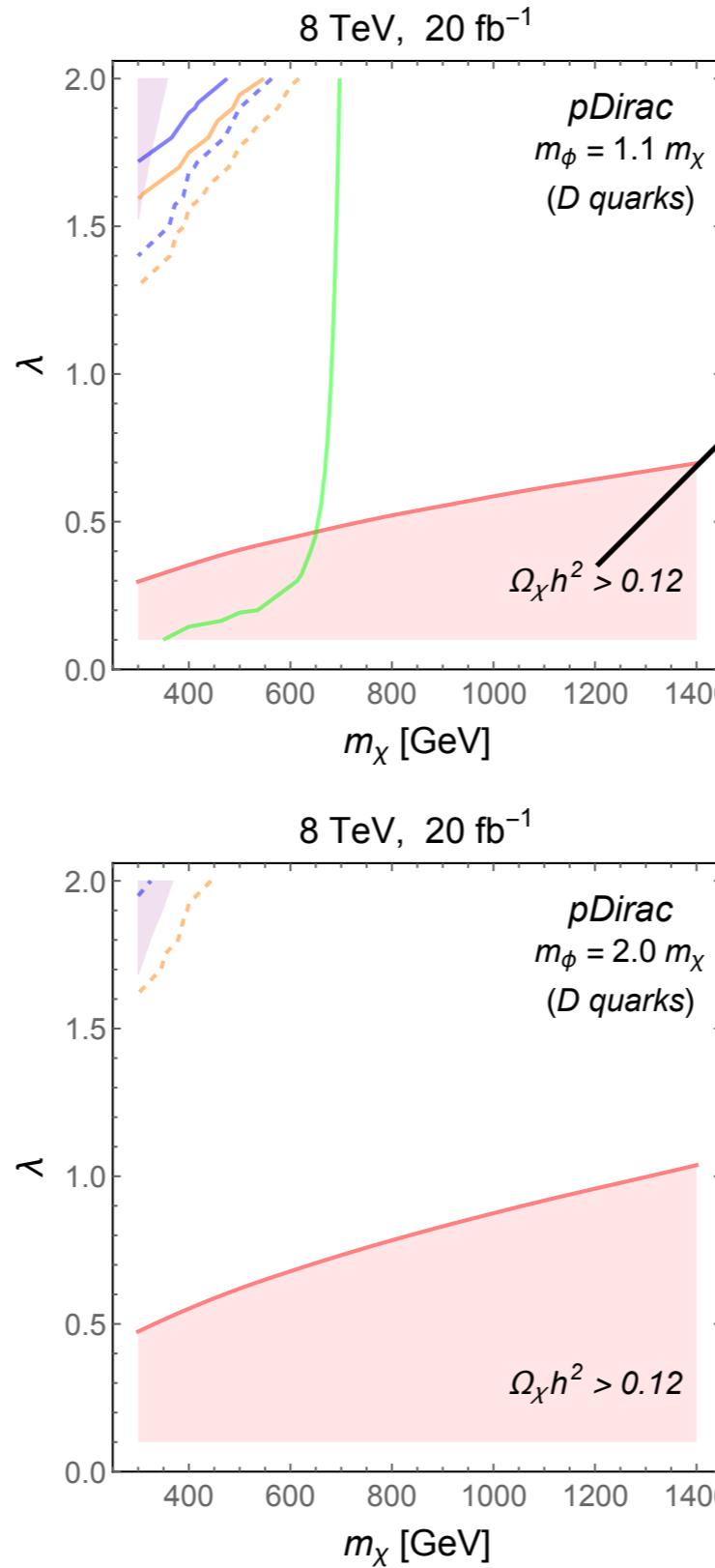
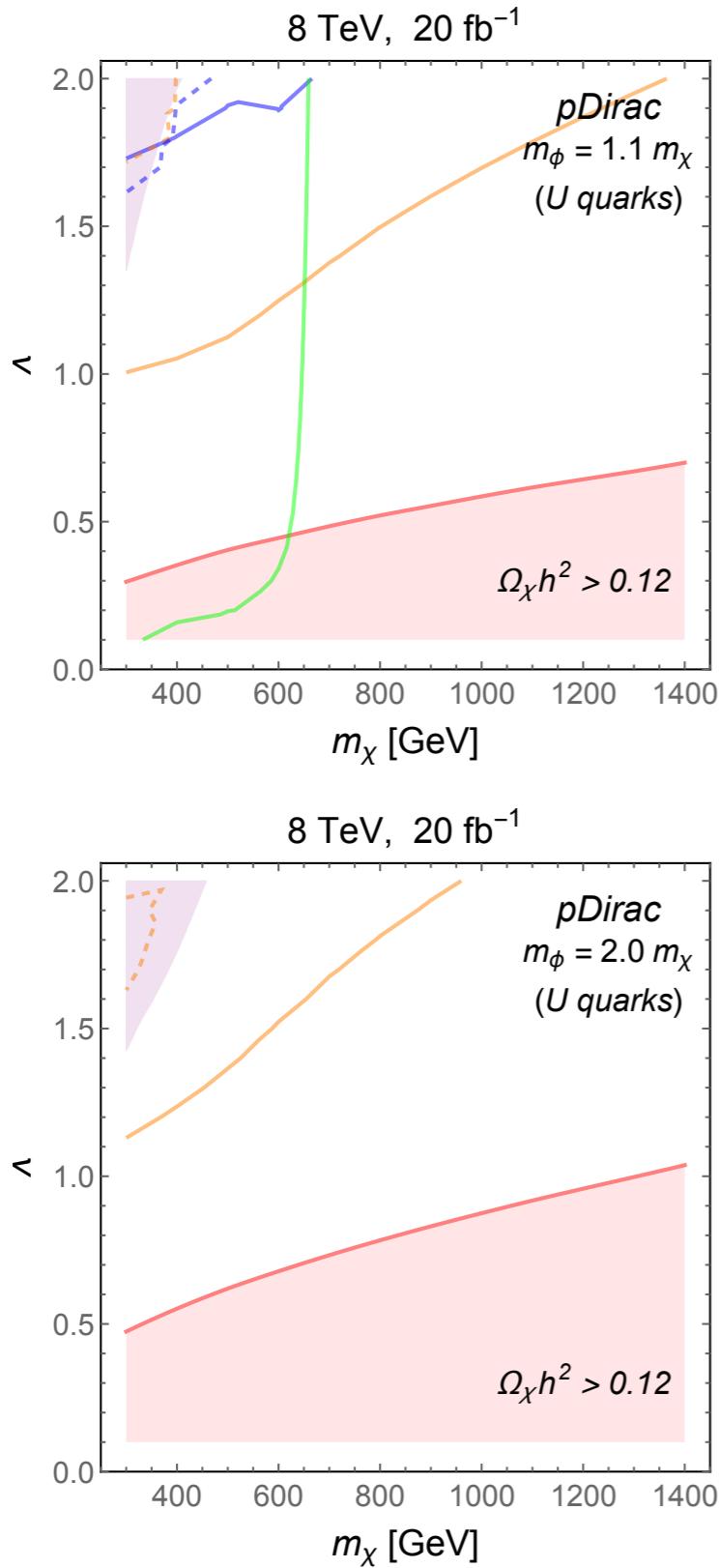
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- **Bounds from 8 TeV:**



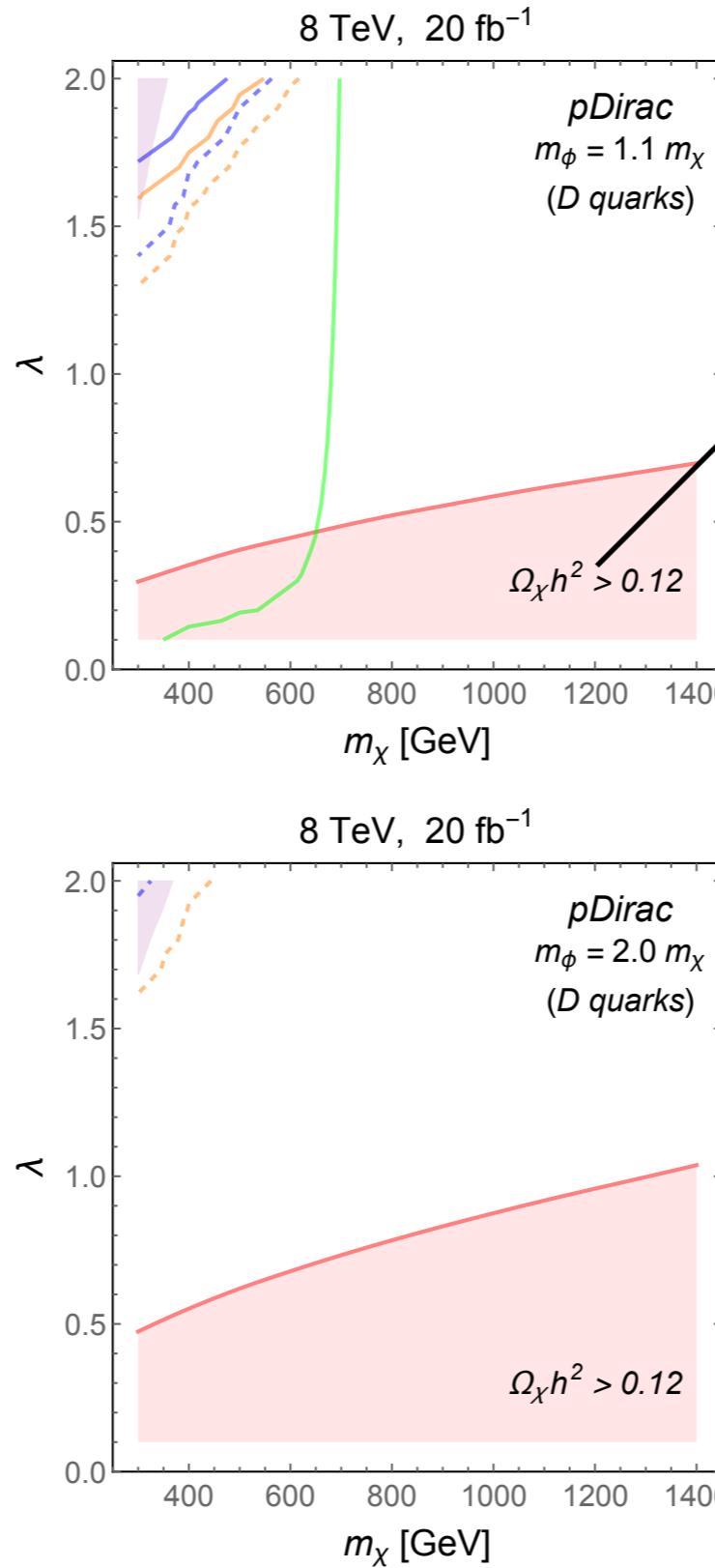
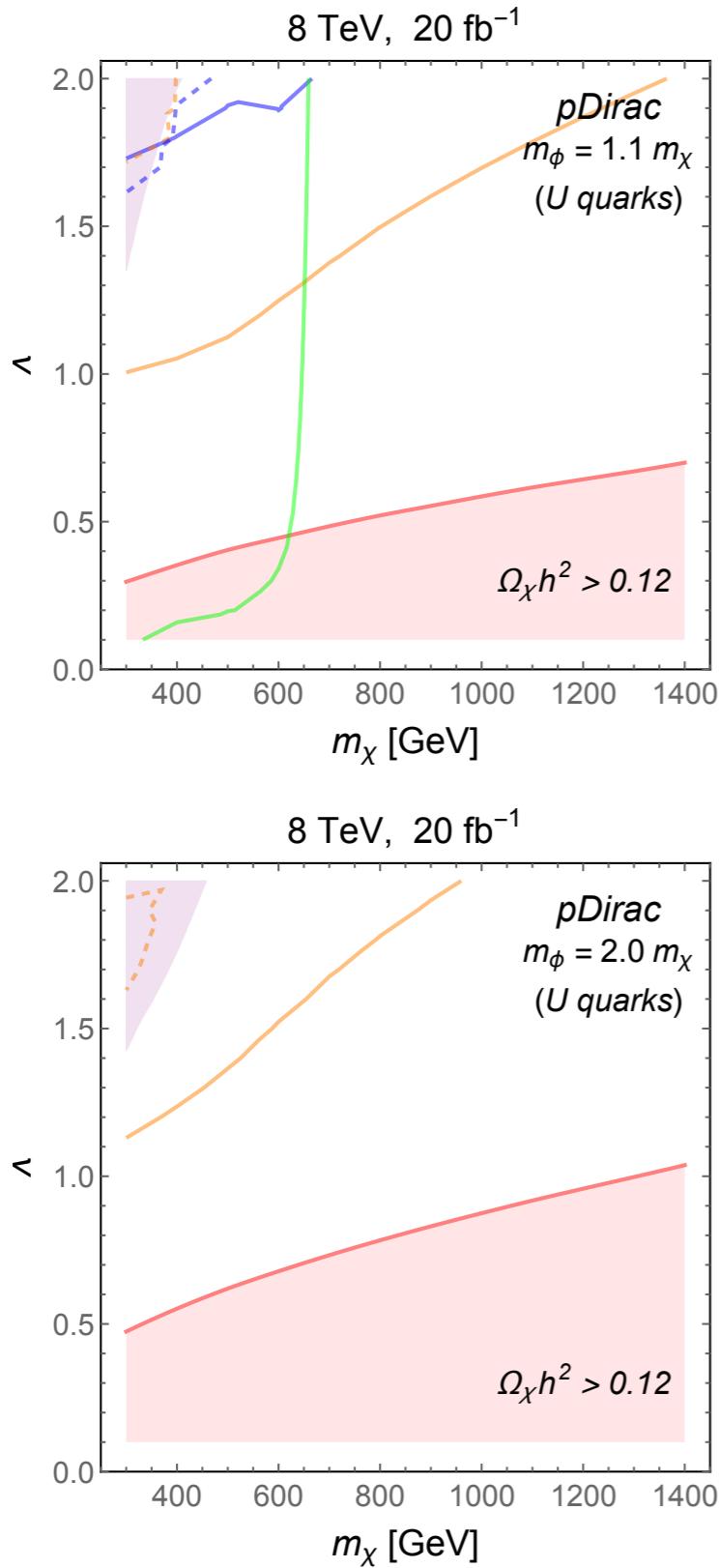
Purple jets+MET bound
Blue is RL model
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- **Bounds from 8 TeV:**



MicrOmegas

Purple jets+MET bound

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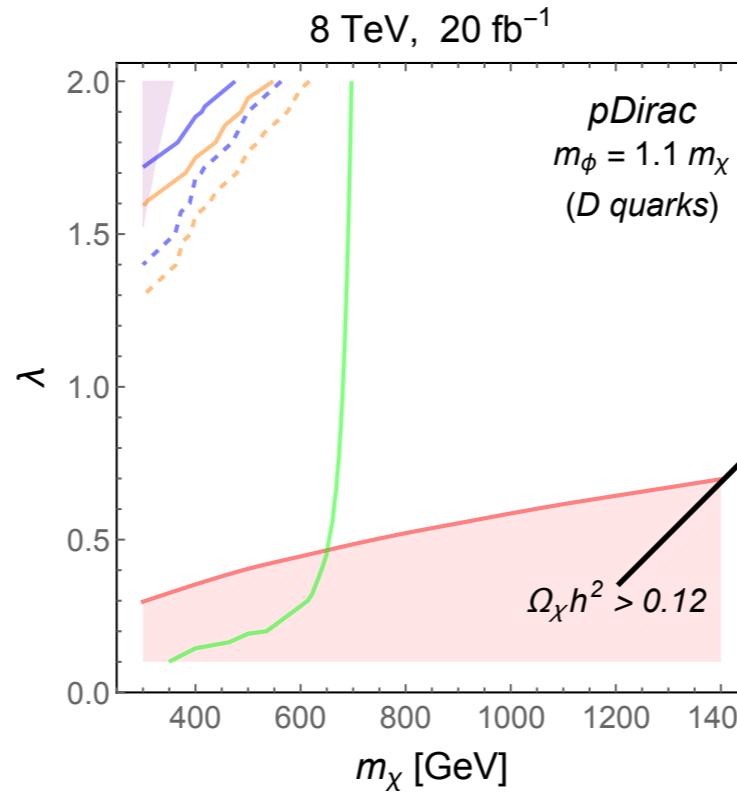
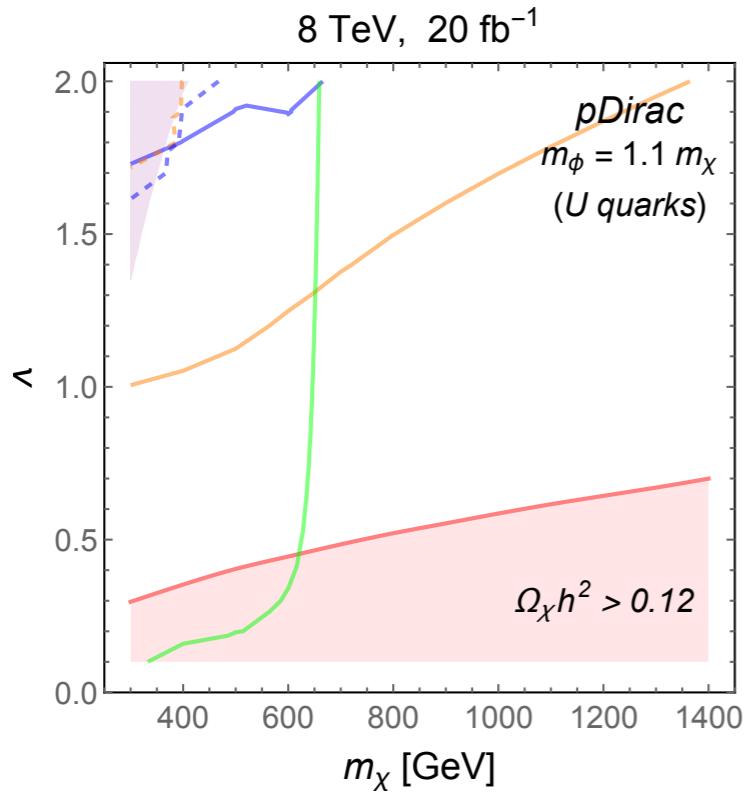
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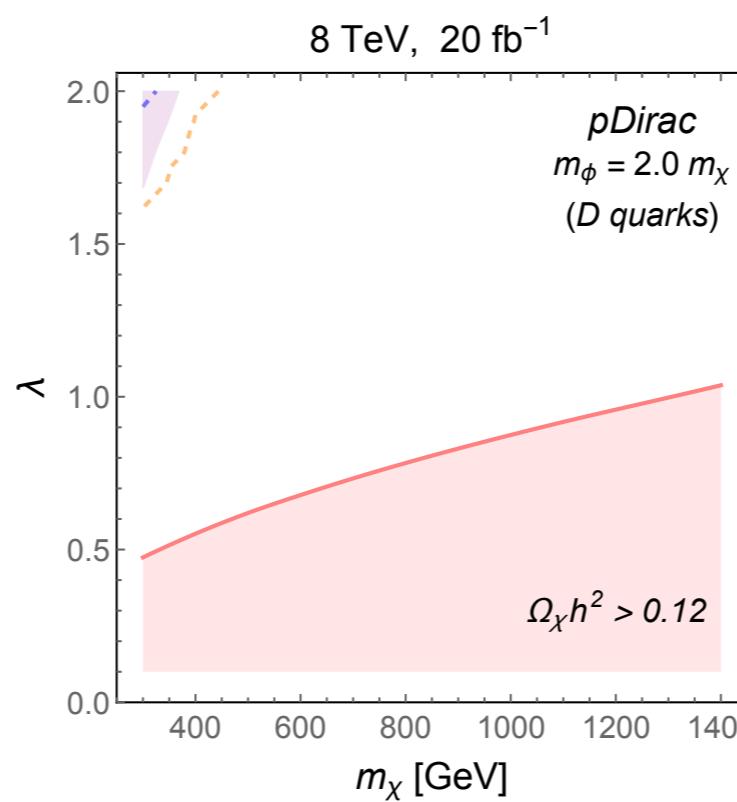
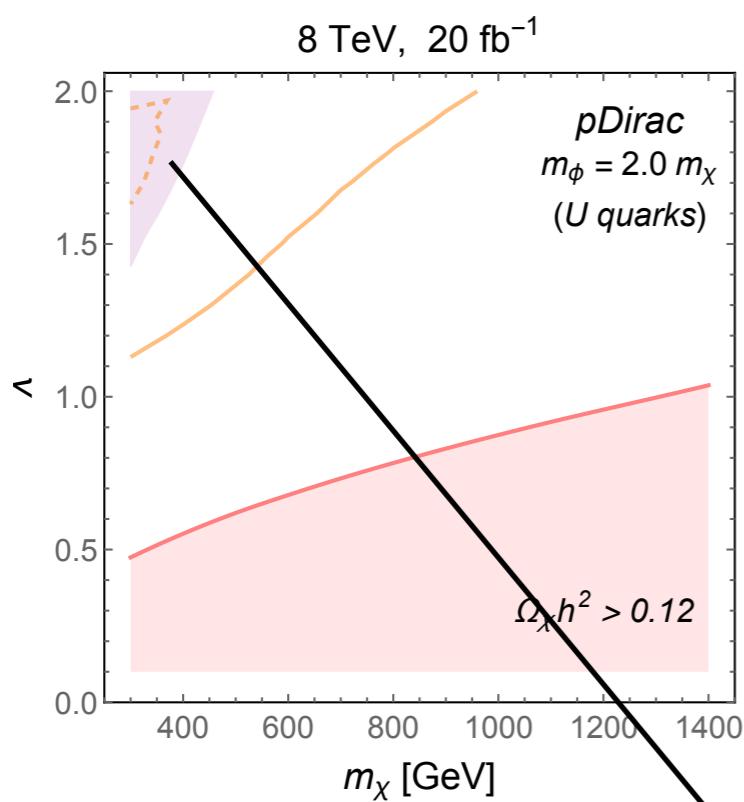
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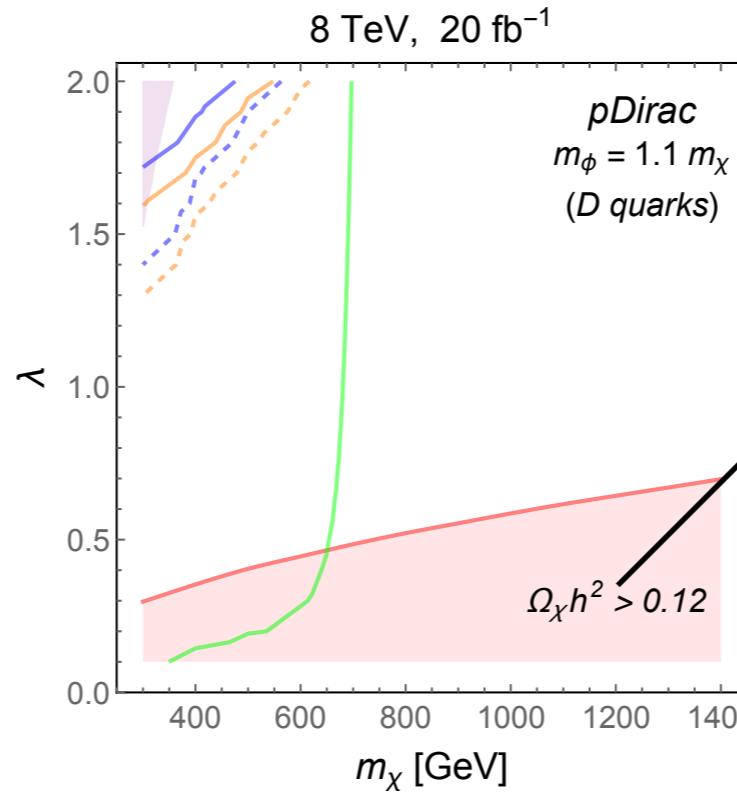
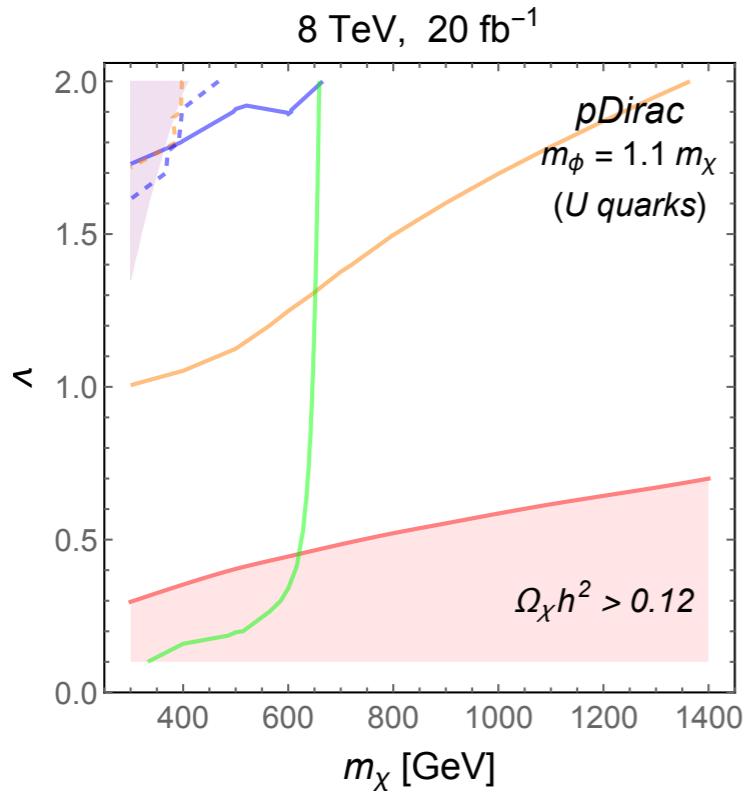


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- **Bounds from 8 TeV:**



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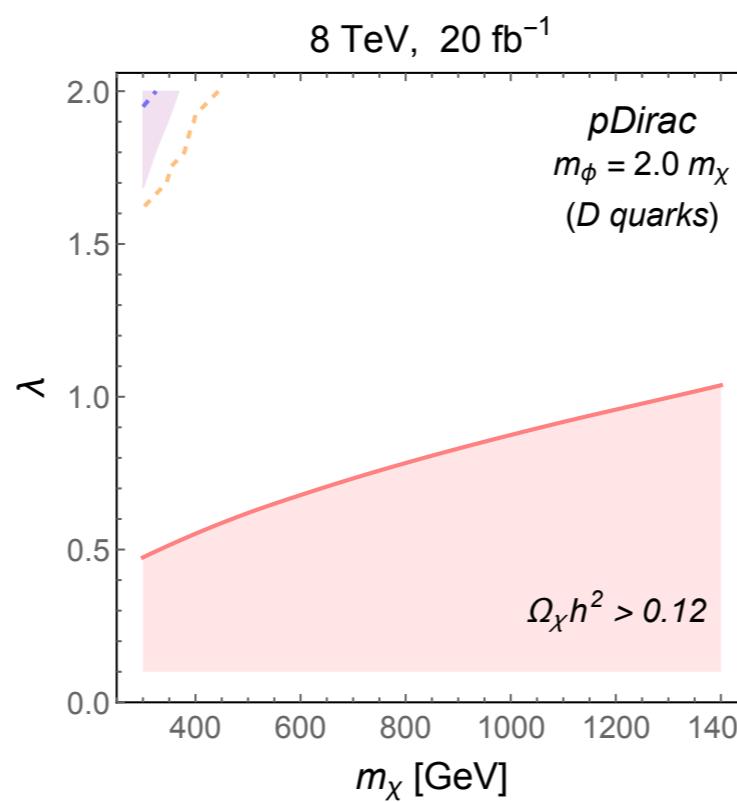
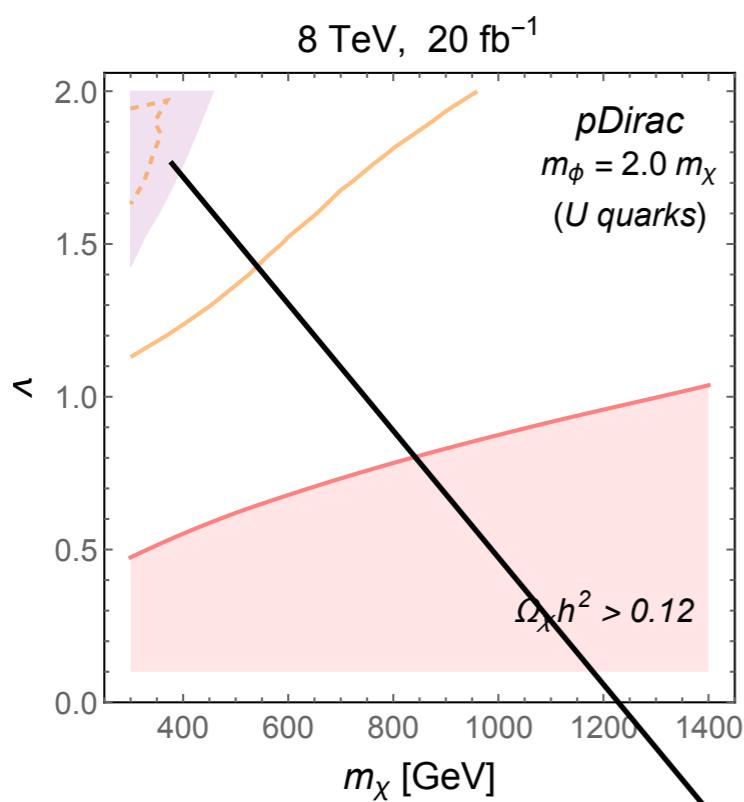
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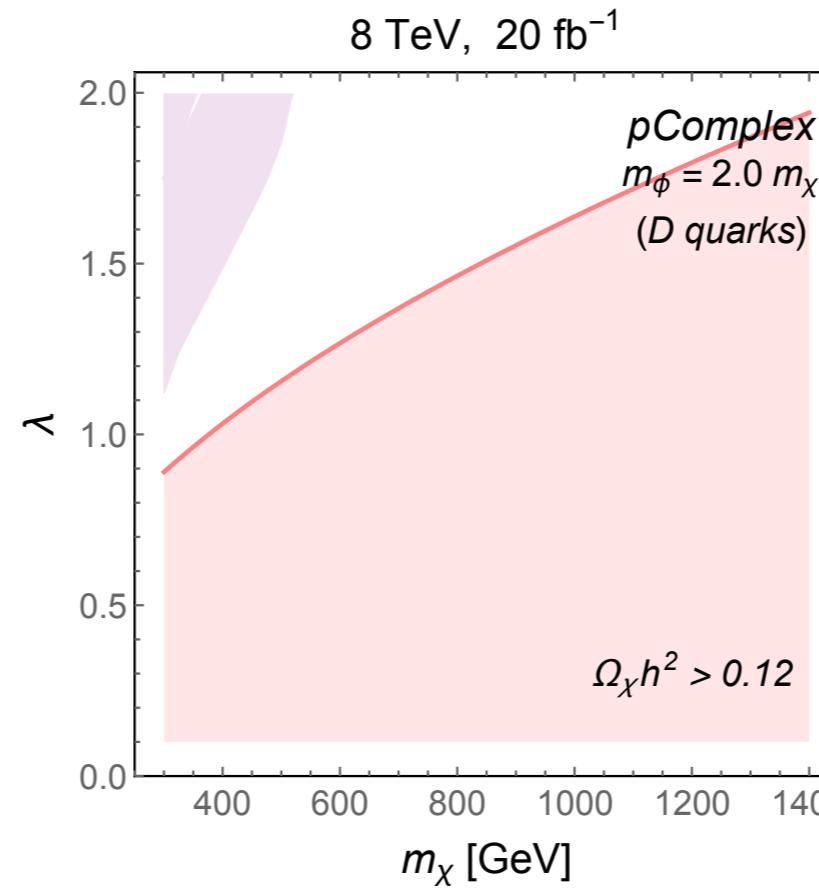
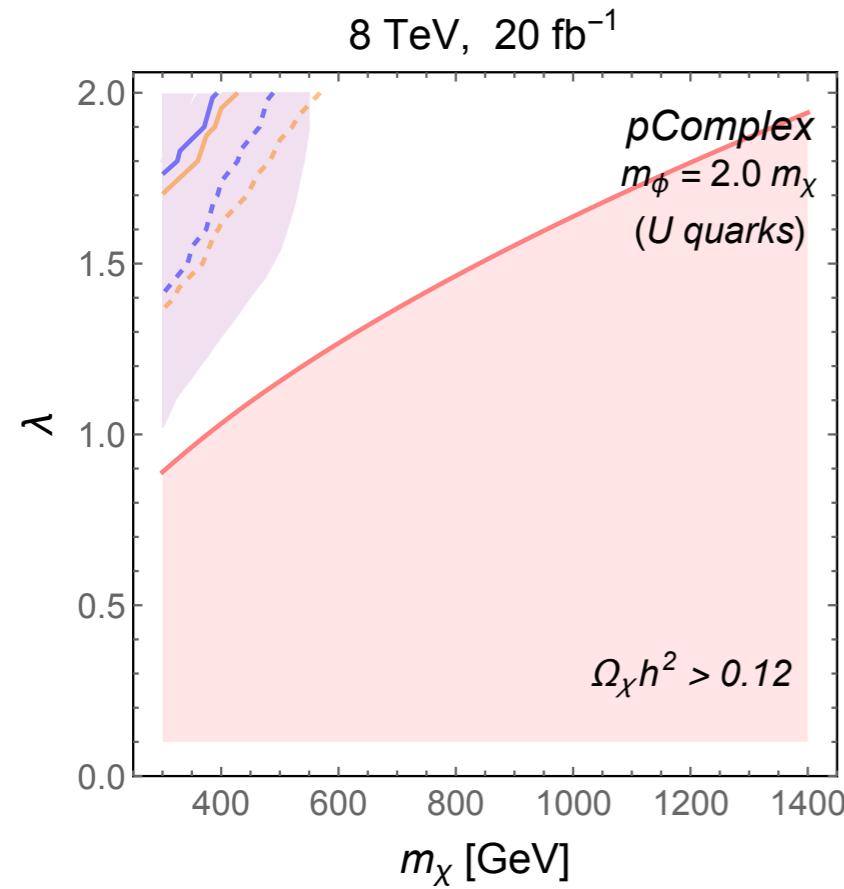
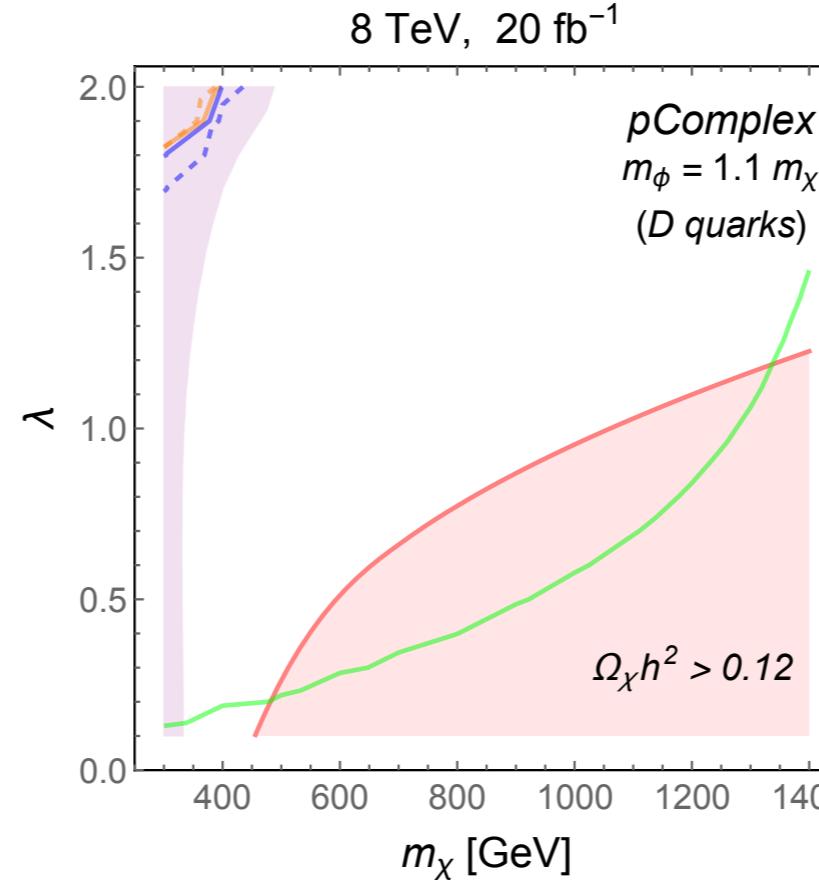
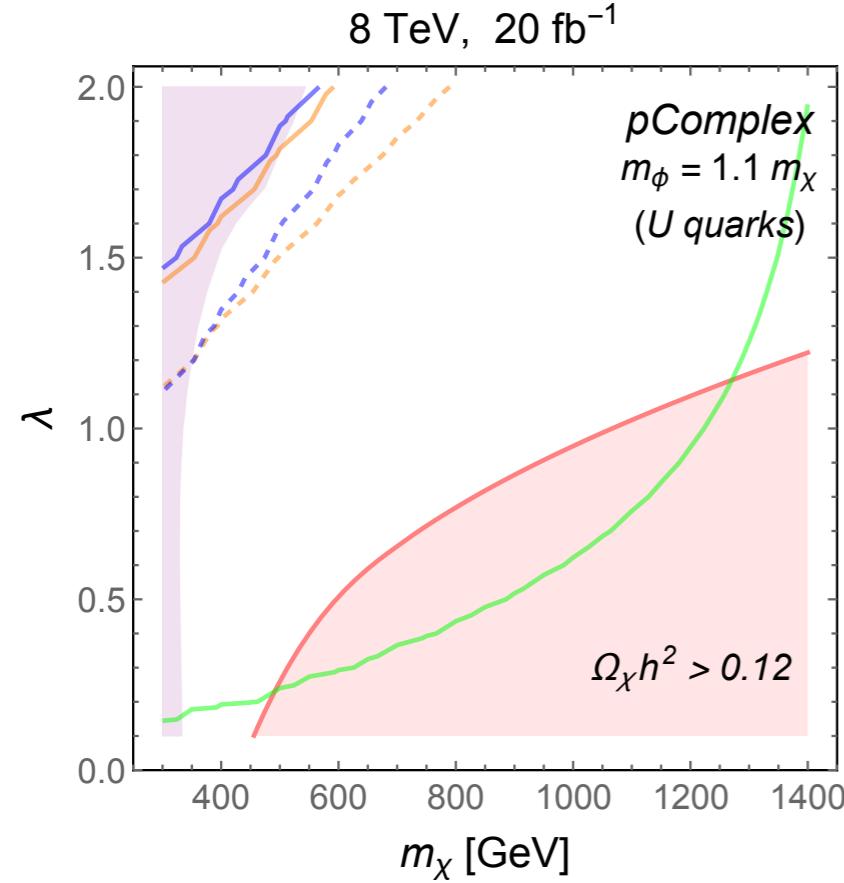
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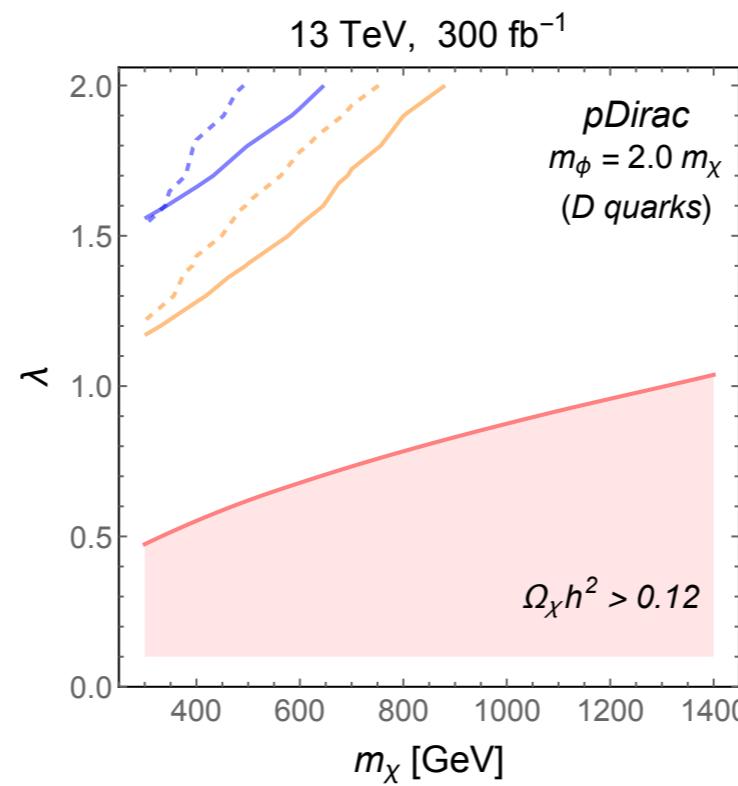
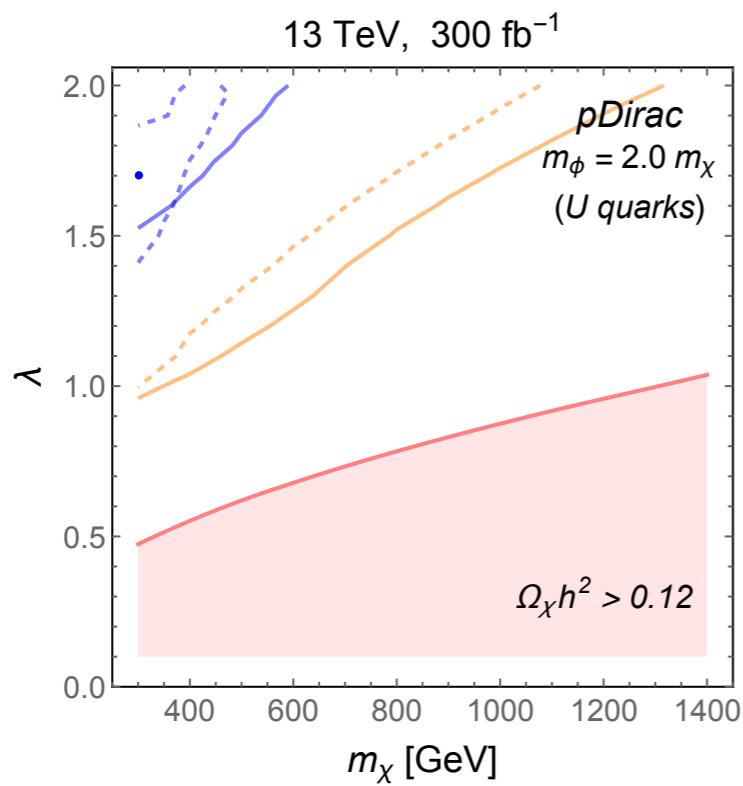
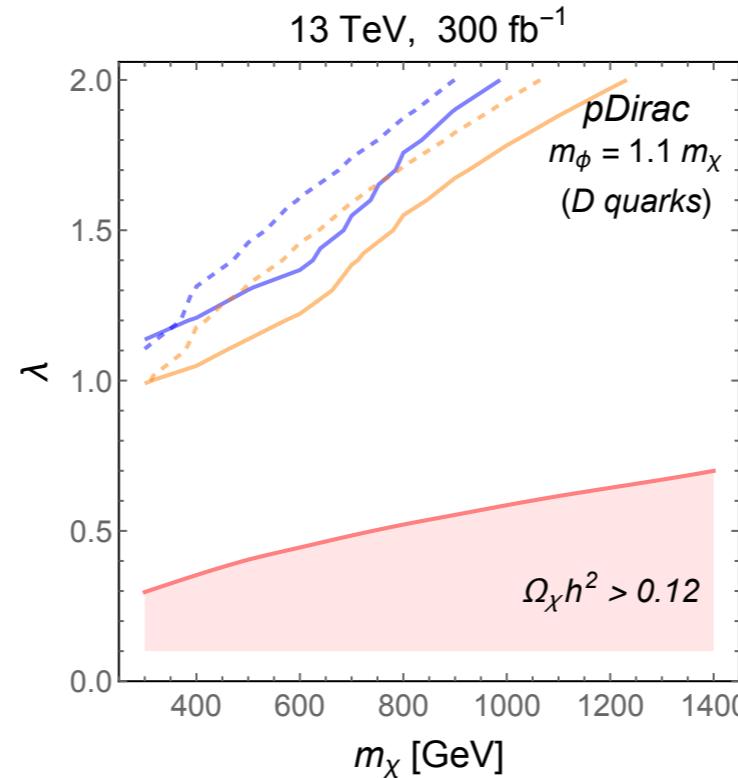
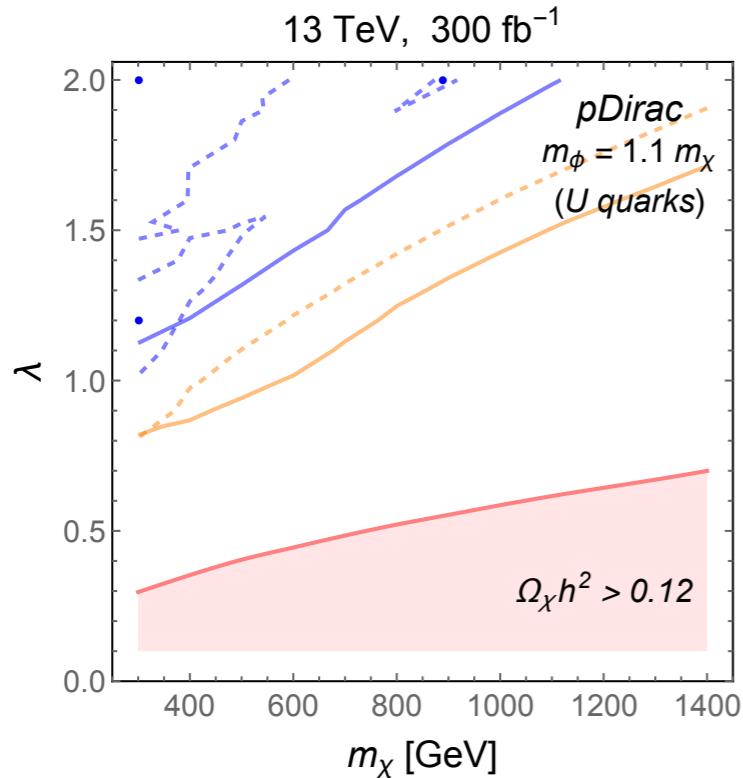


Reinterpretation of T2qq



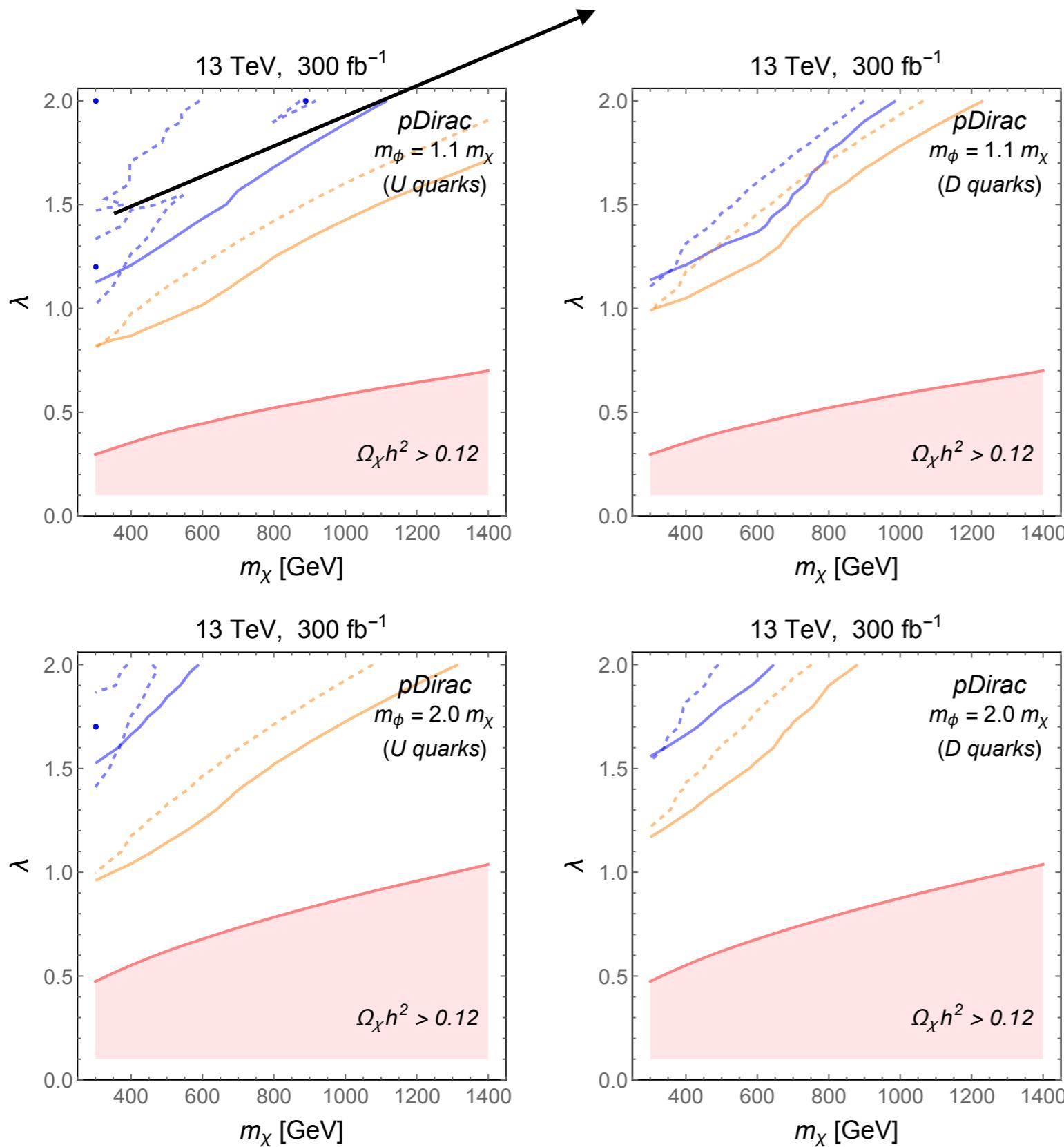
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- Prospect for 13 TeV:



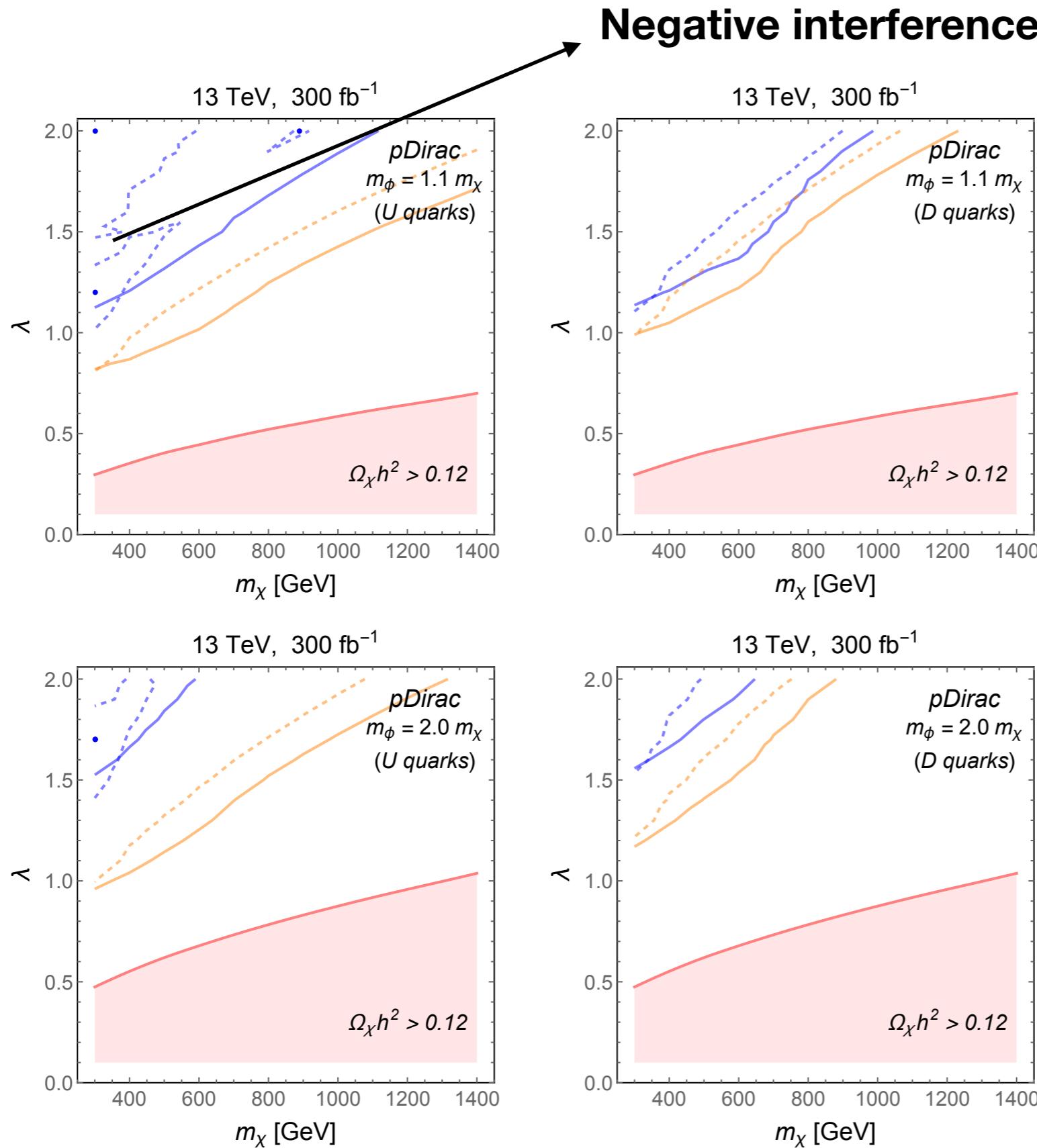
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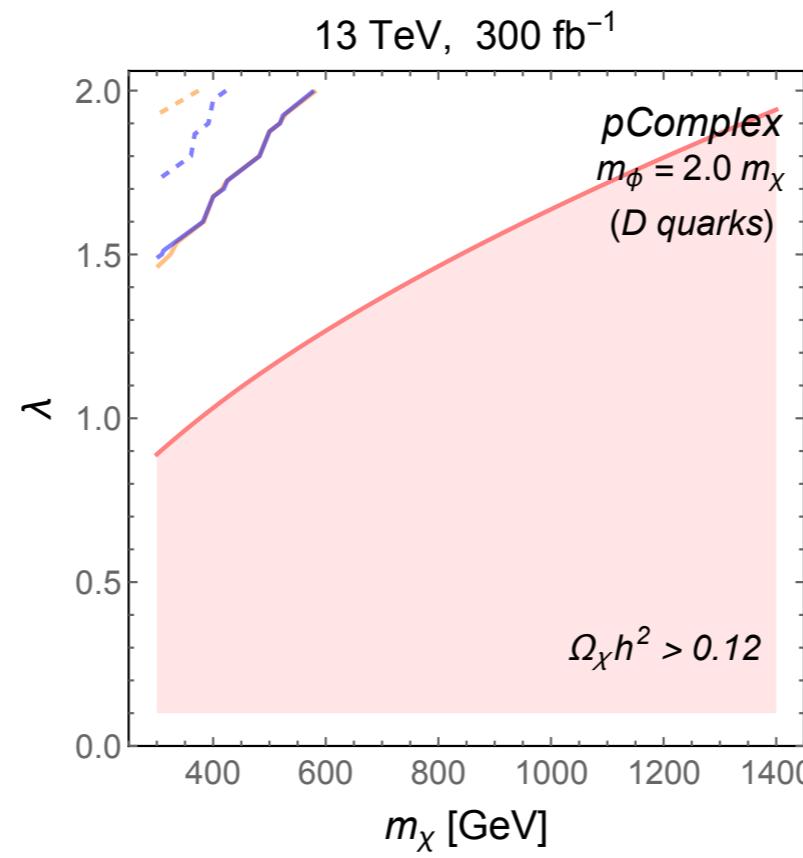
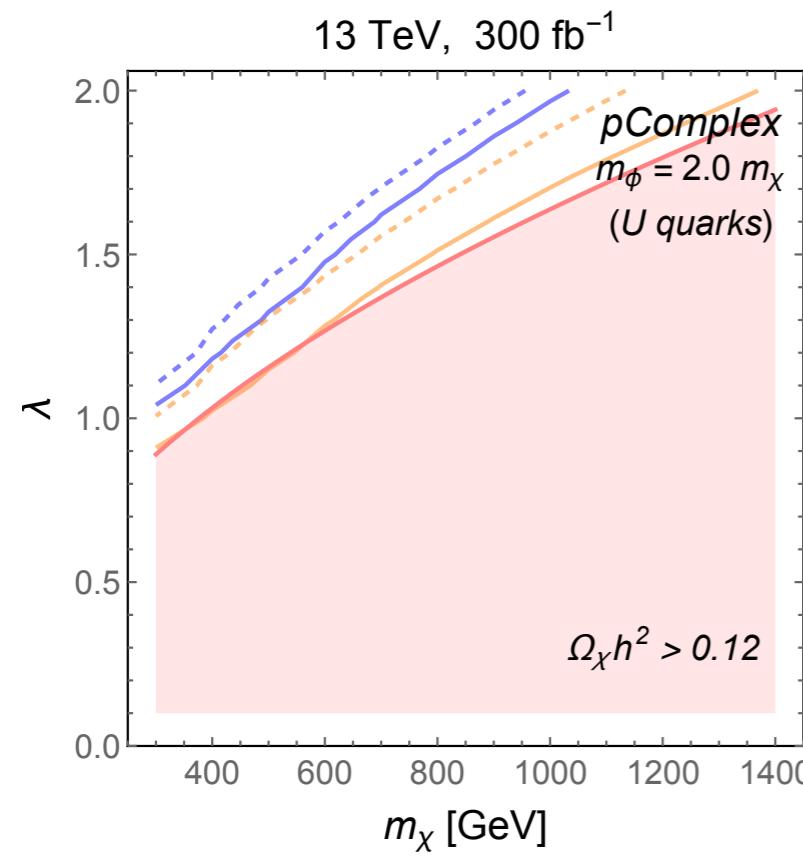
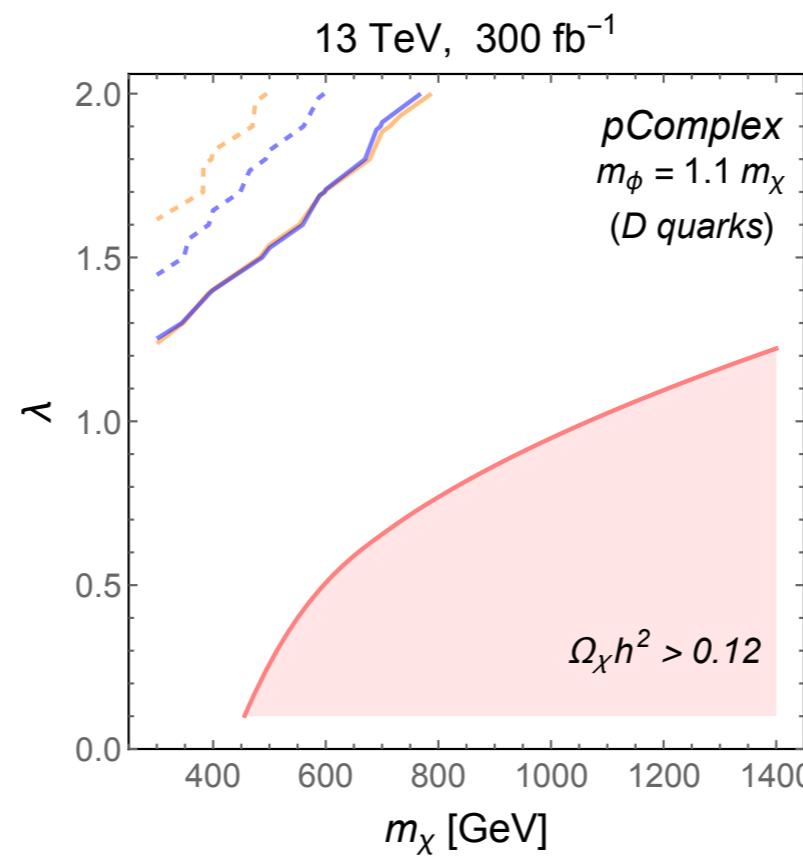
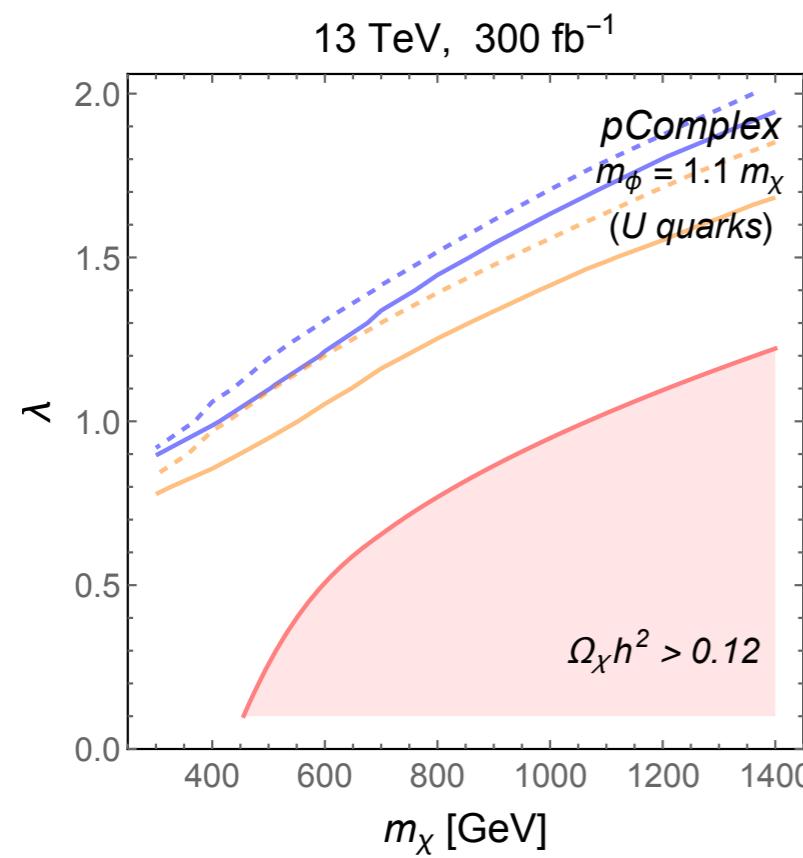


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Conclusions

- DM is one of the reasons that we need physics BSM.
- It is one that has some experimental evidence (although only gravitational)
- In models where DM couples to the SM one can use a mono-X channel to discover it....
- or use a cascade decay into DM (susy like.....)

- In this talk I have shown an alternative way of discovering DM using dilepton events
- If there are color and uncolor messengers (à la susy) one can produce dilepton via loops.
- Interference of those loops with the SM DY production can be the handle to discover new physics at the LHC.
- Angular correlations depend heavily on the spin of DM.
- The same approach can be use for LQ or Z'.