



Bounds on Dark Matter Interactions with Electroweak Gauge Bosons

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1210.0525

Dark Matter @LHC
KICP 2013

Dark Matter EFTs

We want to use _____ to bound DM interactions with _____ and
(an experiment) (some SM particles)

to compare these limits with those from _____ on these same interactions.
(another experiment)

joke stolen from R. Cotta

Dark Matter EFTs

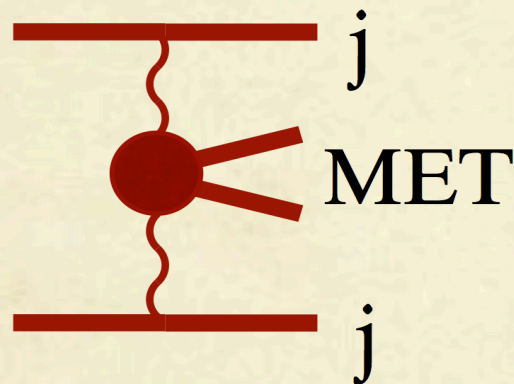
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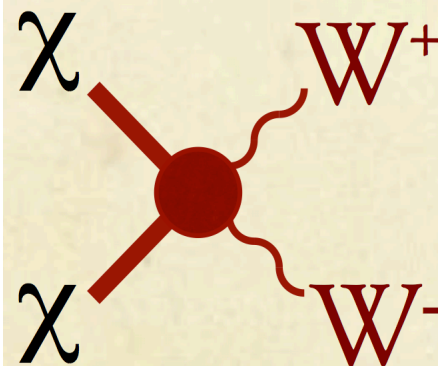
joke stolen from R. Cotta

We want to use $2j + \text{MET (VBF)}$ to bound DM interactions with γ, Z^0, W^\pm and
(an experiment) (some SM particles)

to compare these limits with those from $\text{Dwarf Galaxy Observations}$ on these same interactions.
(another experiment)



Vector Boson
fusion @LHC



Cosmic
Annihilation
products

Why VBF?

- Probes $\chi\chi$ interactions with EW gauge bosons
- Relate to indirect DM searches



» However, expect weaker constraints than mono-jet signatures

Assumptions:

- Respect $U(1)_{\text{em}}$ gauge invariance
- χ is Dirac fermion
- UV completions will result in non-trivial combinations of operators for collider subprocesses WW , ZZ , $\gamma\gamma$, γZ
 - » Weight processes and add incoherently
 - » Assume most simple structure

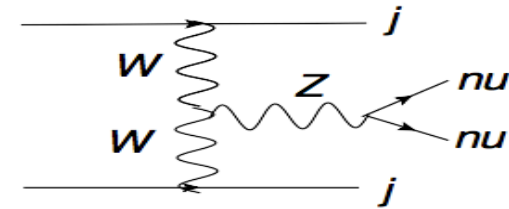
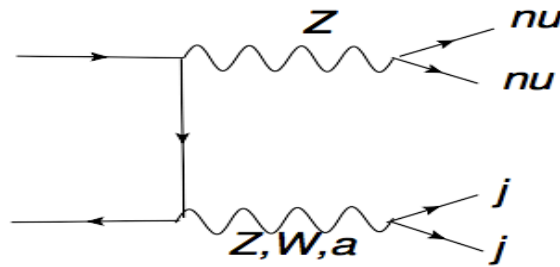
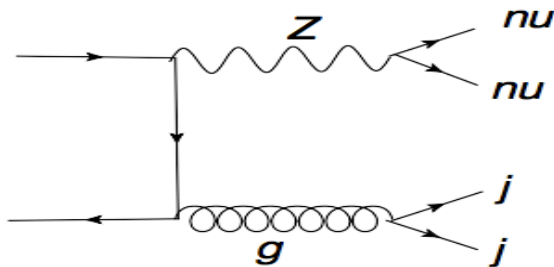
Set of $\chi\chi V$ and $\chi\chi VV$ operators

| Name | Expression | Norm. | Vertices | Sub-Proc. | Ann. |
|-----------------|--|----------------|----------|------------------|------------------------------|
| <i>dim</i> = 5: | | | | | |
| D5a | $\bar{\chi}\chi V^{a\mu}V_{\mu}^a$ | Λ^{-1} | 4pt | ZZ, WW | v^2 |
| D5b | $\bar{\chi}i\gamma_5\chi V^{a\mu}V_{\mu}^a$ | Λ^{-1} | 4pt | ZZ, WW | 1 |
| D5c | $\bar{\chi}\sigma_{\mu\nu}t^a\chi V^{a\mu\nu}$ | Λ^{-1} | 3/4pt | A, Z, WW | 1 |
| D5d | $\bar{\chi}\sigma_{\mu\nu}t^a\chi\tilde{V}^{a\mu\nu}$ | Λ^{-1} | 3/4pt | A, Z, WW | 1 (VV), v^2 ($f\bar{f}$) |
| <i>dim</i> = 6: | | | | | |
| D6a | $\bar{\chi}\gamma_{\mu}t^aD_{\nu}\chi V^{a\mu\nu}$ | Λ^{-2} | 3/4pt | A, Z, WW | 1 |
| D6b | $\bar{\chi}\gamma_{\mu}\gamma_5t^aD_{\nu}\chi V^{a\mu\nu}$ | Λ^{-2} | 3/4pt | A, Z, WW | 1 (VV), v^2 ($f\bar{f}$) |
| <i>dim</i> = 7: | | | | | |
| D7a | $\bar{\chi}\chi V^{\mu\nu}V_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | v^2 |
| D7b | $\bar{\chi}i\gamma_5\chi V^{\mu\nu}V_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | 1 |
| D7c | $\bar{\chi}\chi V^{\mu\nu}\tilde{V}_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | v^2 |
| D7d | $\bar{\chi}i\gamma_5\chi V^{\mu\nu}\tilde{V}_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | 1 |

Set of $\chi\chi V$ and $\chi\chi VV$ operators

| Name | Expression | Norm. | Vertices | Sub-Proc. | Ann. |
|------------|--|----------------|----------|------------------|--|
| $dim = 5:$ | | | | | |
| D5a | $\bar{\chi}\chi V^{a\mu}V_{\mu}^a$ | Λ^{-1} | 4pt | ZZ, WW | scalar mediator |
| D5b | $\bar{\chi}i\gamma_5\chi V^{a\mu}V_{\mu}^a$ | Λ^{-1} | 4pt | ZZ, WW | pseudoscalar mediator |
| D5c | $\bar{\chi}\sigma_{\mu\nu}t^a\chi V^{a\mu\nu}$ | Λ^{-1} | 3/4pt | A, Z, WW | Dipole operators $\bar{\chi}\sigma_{\mu\nu}\chi V^{a\mu\nu}$ |
| D5d | $\bar{\chi}\sigma_{\mu\nu}t^a\chi\tilde{V}^{a\mu\nu}$ | Λ^{-1} | 3/4pt | A, Z, WW | |
| $dim = 6:$ | | | | | |
| D6a | $\bar{\chi}\gamma_{\mu}t^aD_{\nu}\chi V^{a\mu\nu}$ | Λ^{-2} | 3/4pt | A, Z, WW | vector mediator $\bar{\chi}\gamma_{\mu}\chi V^{a\mu\nu}D_{\nu}$ |
| D6b | $\bar{\chi}\gamma_{\mu}\gamma_5t^aD_{\nu}\chi V^{a\mu\nu}$ | Λ^{-2} | 3/4pt | A, Z, WW | |
| $dim = 7:$ | | | | | |
| D7a | $\bar{\chi}\chi V^{\mu\nu}V_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | arise @ 1-loop $\bar{\chi}\chi V^{\mu\nu}V_{\mu\nu}$ |
| D7b | $\bar{\chi}i\gamma_5\chi V^{\mu\nu}V_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | |
| D7c | $\bar{\chi}\chi V^{\mu\nu}\tilde{V}_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | |
| D7d | $\bar{\chi}i\gamma_5\chi V^{\mu\nu}\tilde{V}_{\mu\nu}$ | Λ^{-3} | 4pt | AA, AZ, ZZ, WW | |

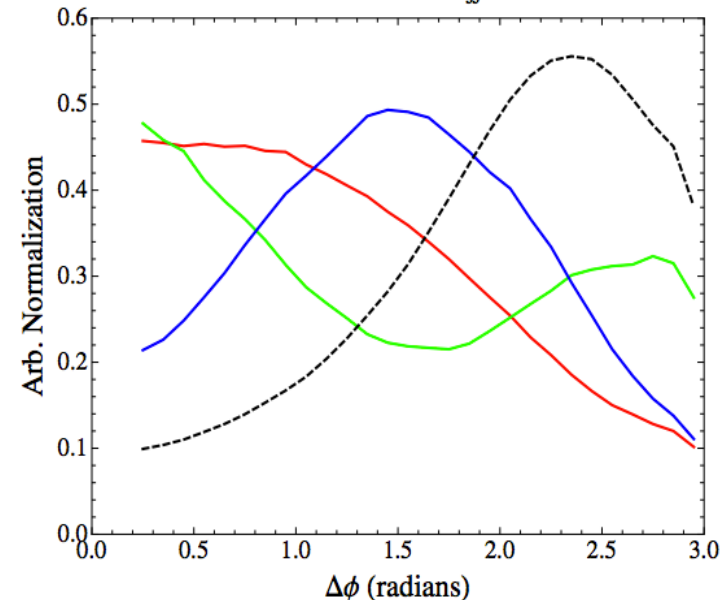
VBF Backgrounds @LHC



Piggy-back on invisible Higgs studies: devise set of cuts

Eboli, Zeppenfeld: hep-ph/0009158

Signal $\Delta\phi$ Distributions,
Signal RGB and $\sigma_{EW Zjj}$ BG (black).

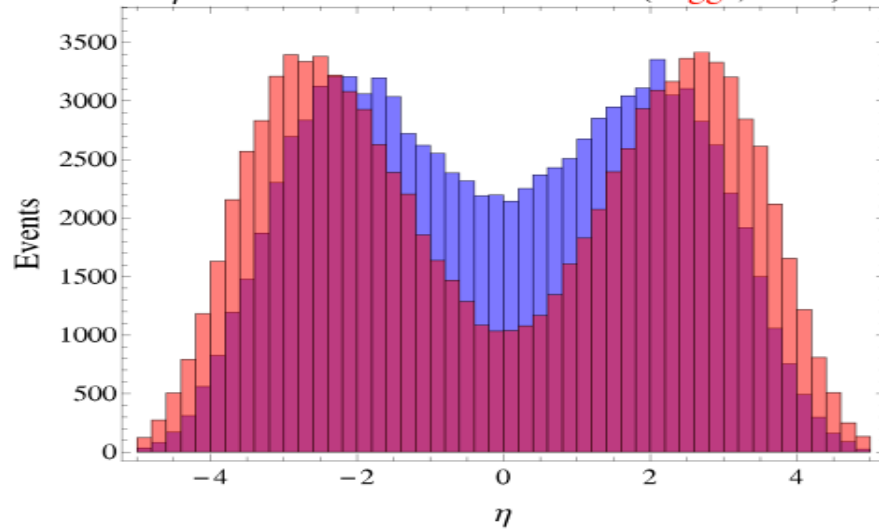


| σ (fb) | QCD Zjj | | QCD Wjj | | EW Zjj | | EW Wjj | | Total | |
|---------------------|-----------|------|-----------|------|----------|------|----------|------|-------|------|
| | [51] | Here | [51] | Here | [51] | Here | [51] | Here | [51] | Here |
| Eqs. (6-8) | 1254 | 1055 | 1284 | 906 | 151 | 148 | 101 | 85 | 2790 | 2194 |
| Eqs. (6-9) + C.J.V. | 71.8 | 56.6 | 70.2 | 47.3 | 14.8 | 14.6 | 9.9 | 8.2 | 167 | 127 |

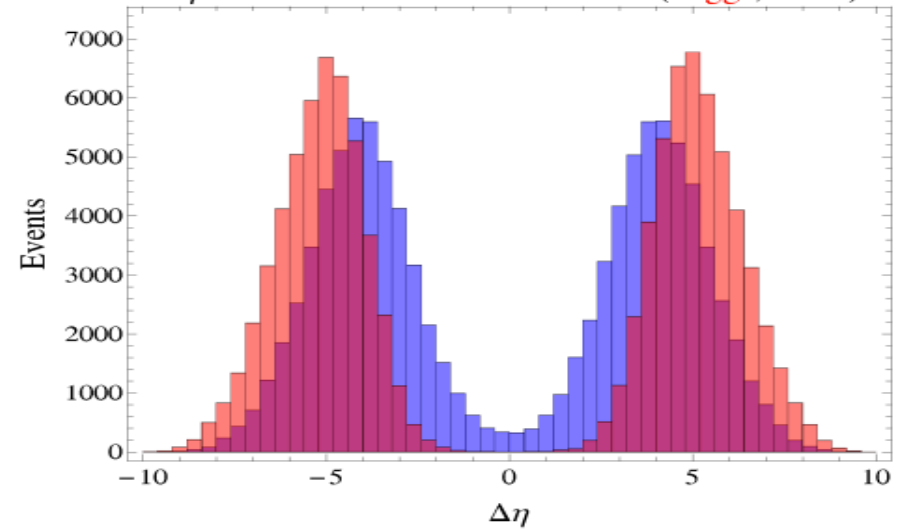
Pseudorapidity Dependence of signal and background



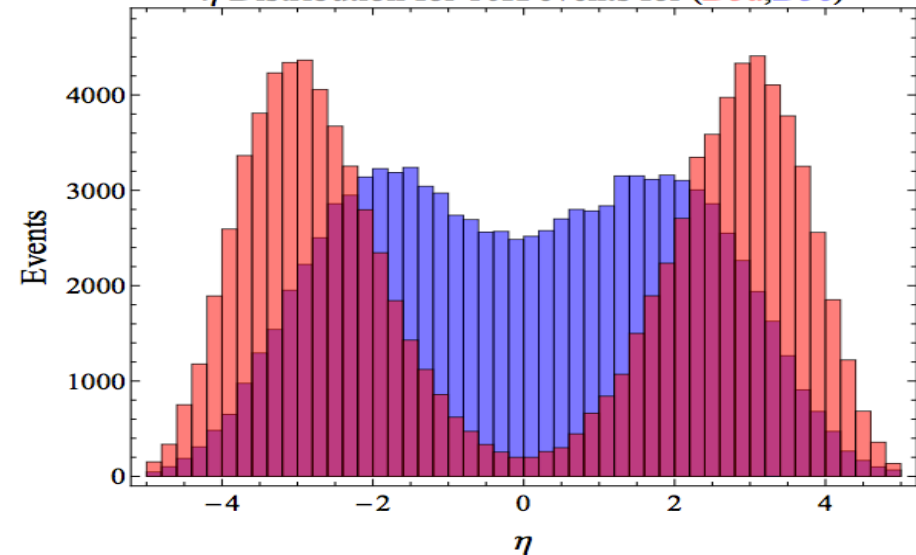
η Distribution for 10K events for (Higgs, ZEW)



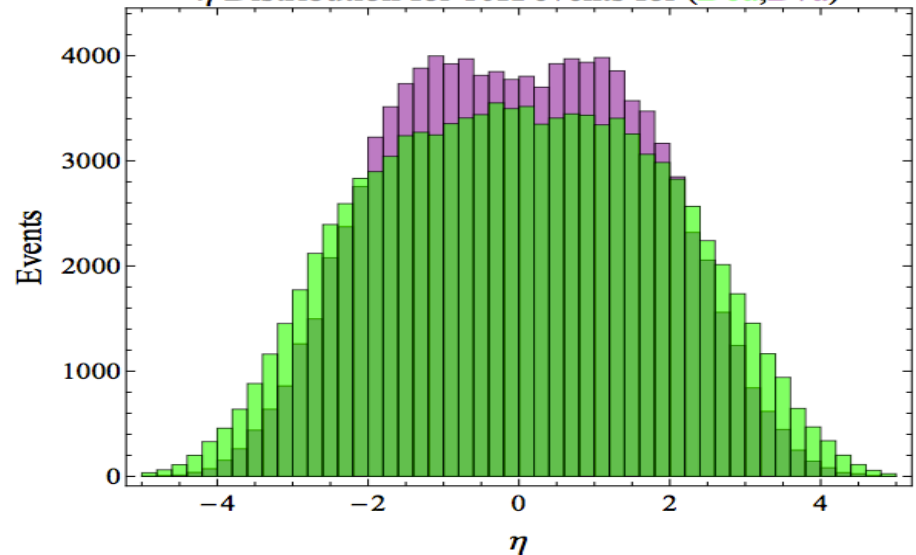
$\Delta\eta$ Distribution for 10K events for (Higgs, ZEW)



η Distribution for 10K events for (D5a, D5c)

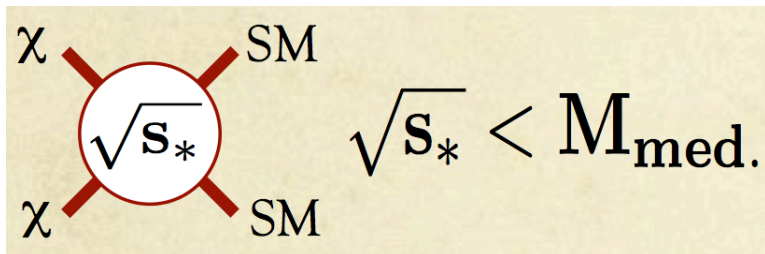


η Distribution for 10K events for (D6a, D7a)



Validity of Effective Field Theory

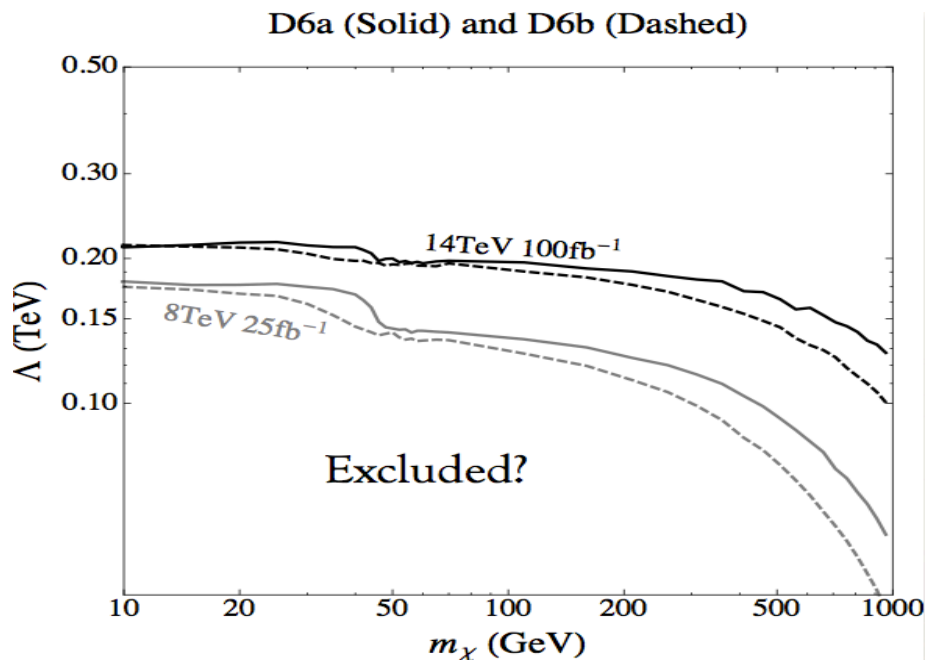
EFT is valid when:



Scattering: $M_{\text{med.}} > |q|, (\ll m_\chi)$

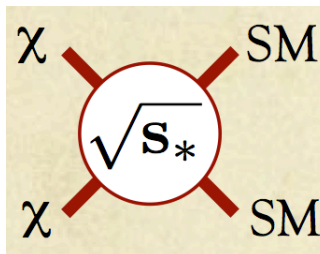
Annihilation: $\sim m_\chi$

Production: $\sim p_T$



Validity of Effective Field Theory

EFT is valid when:

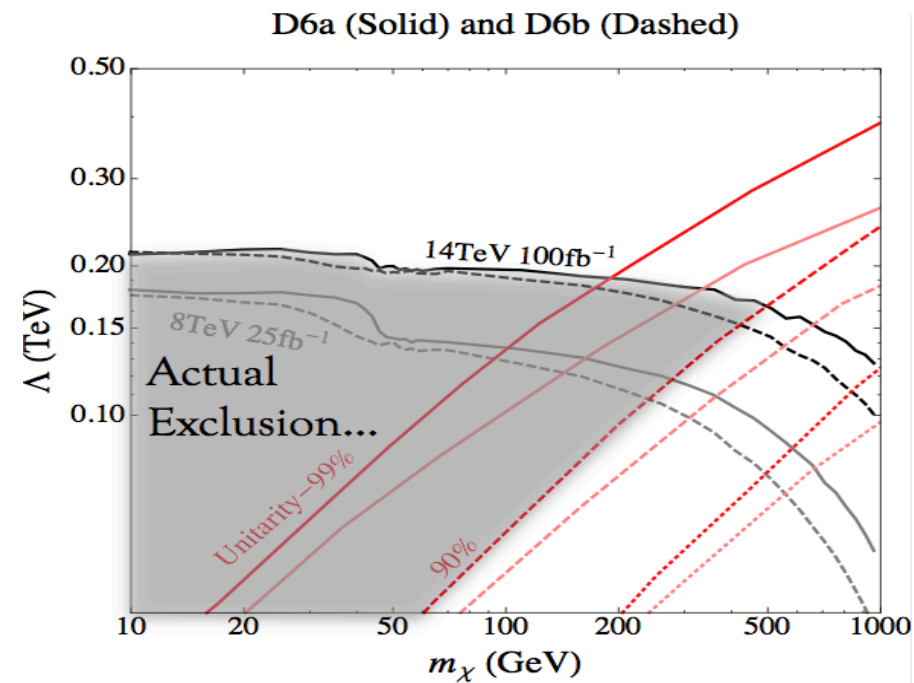
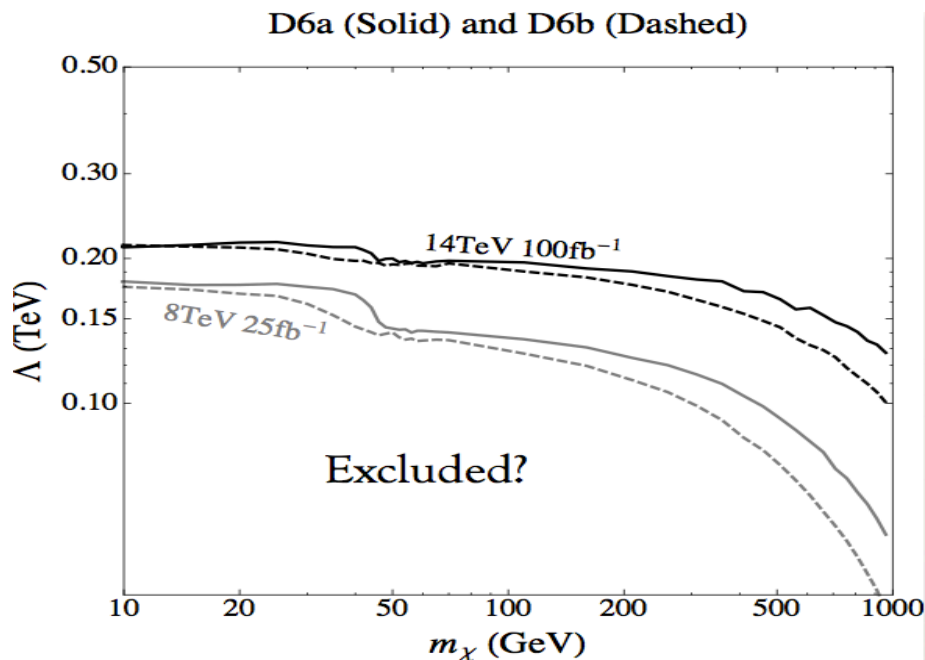


$$\sqrt{s_*} < M_{\text{med.}}$$

Scattering: $M_{\text{med.}} > |q|, (\ll m_\chi)$
 Annihilation: $\sim m_\chi$
 Production: $\sim p_T$

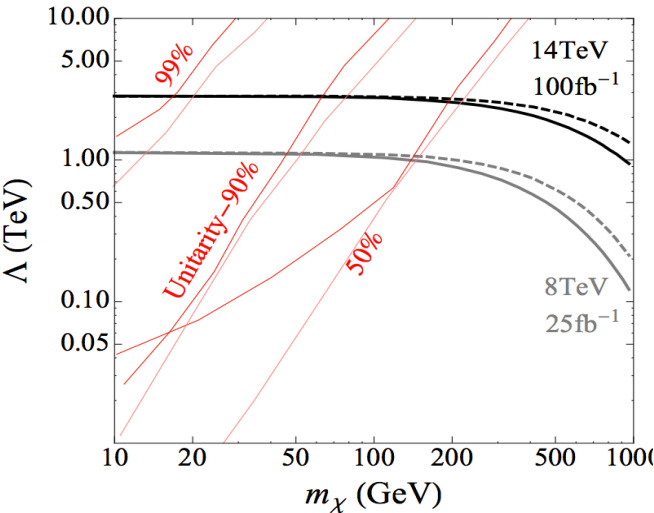
Employ partial wave analysis: $|a_0(qq \rightarrow \chi\chi)| < 1/2$

99%, 90%, 50% VBF events obey this limit

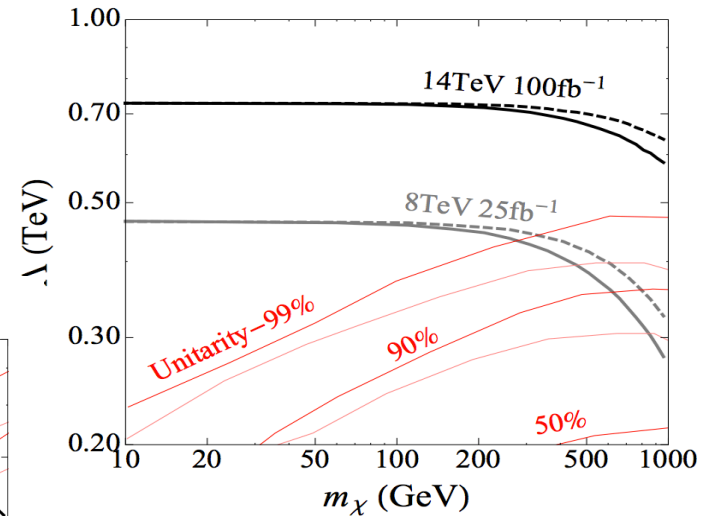


LHC VBF Results

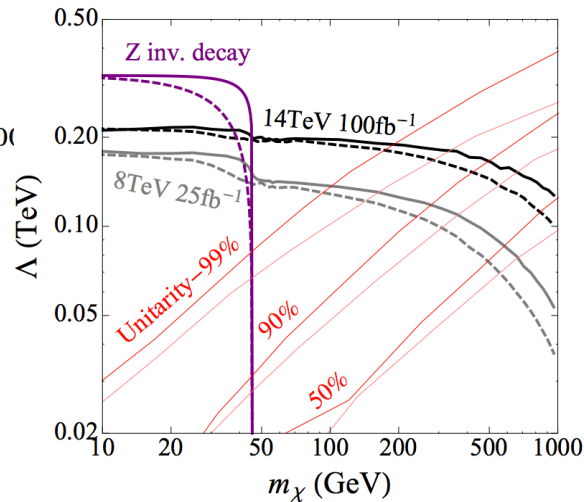
D5a (Solid) and D5b (Dashed)



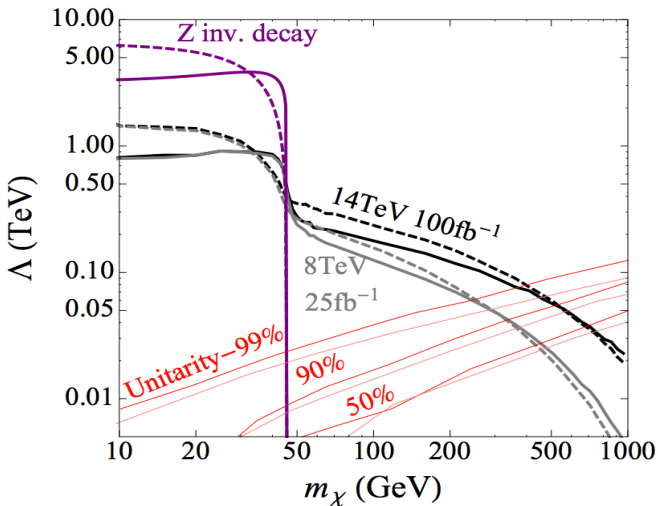
D7a (Solid) and D7b (Dashed)



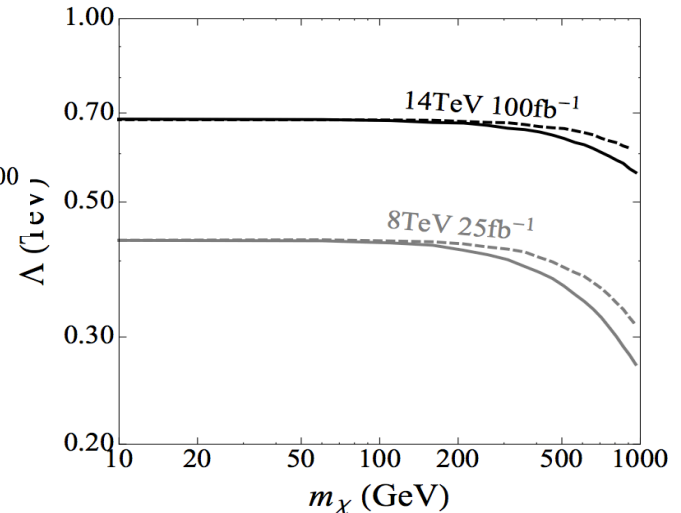
D6a (Solid) and D6b (Dashed)



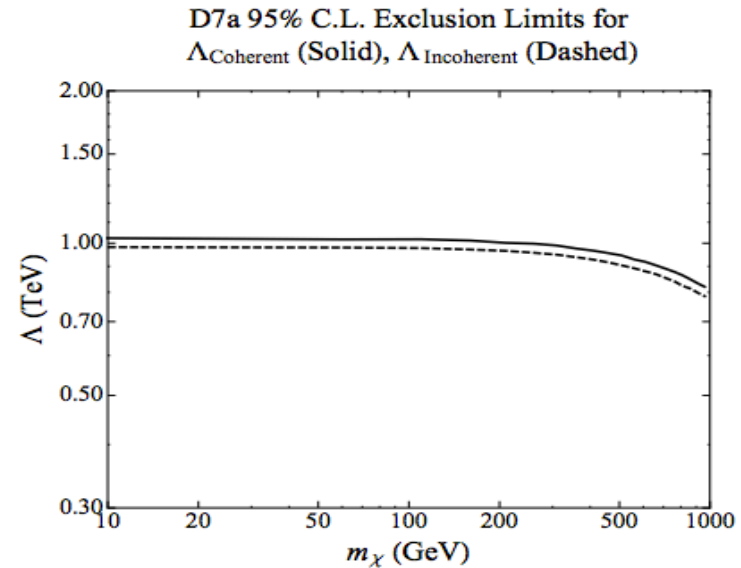
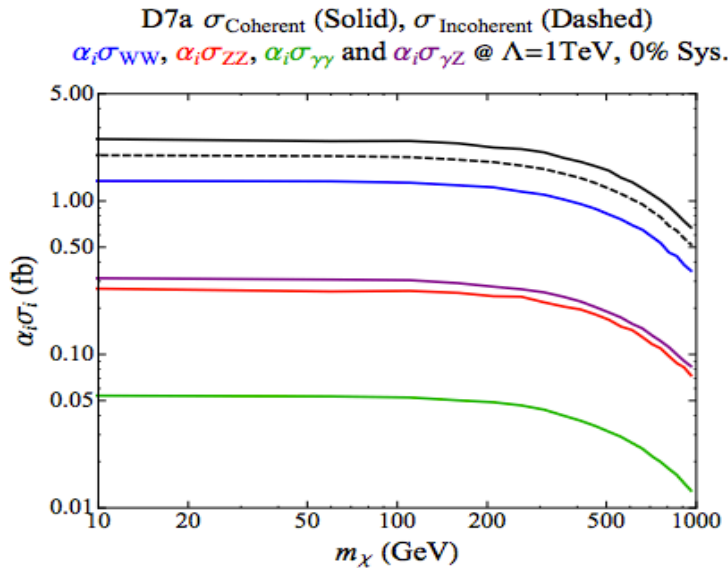
D5c (Solid) and D5d (Dashed)



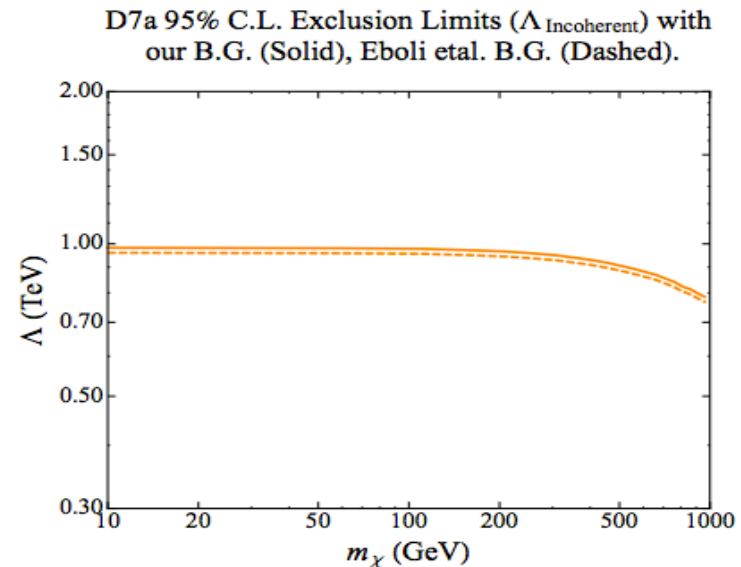
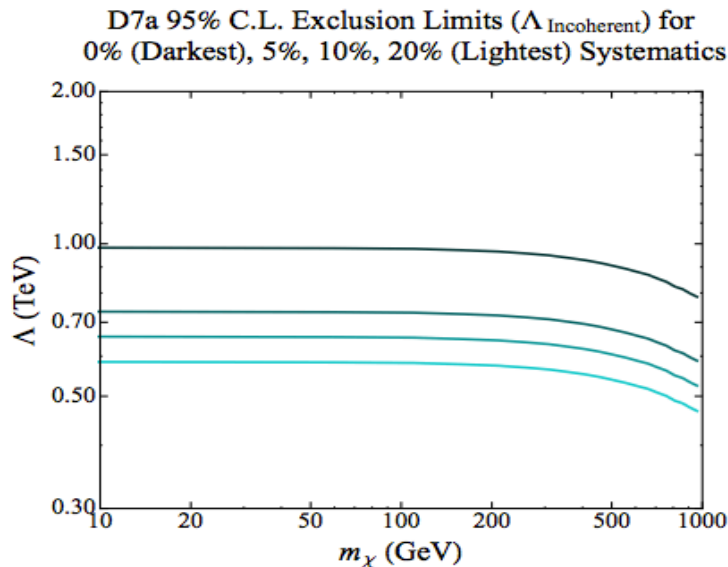
D7c (Solid) and D7d (Dashed)



Effects of Uncertainties on LHC Constraints



14 TeV
 100 fb⁻¹



Include data from:

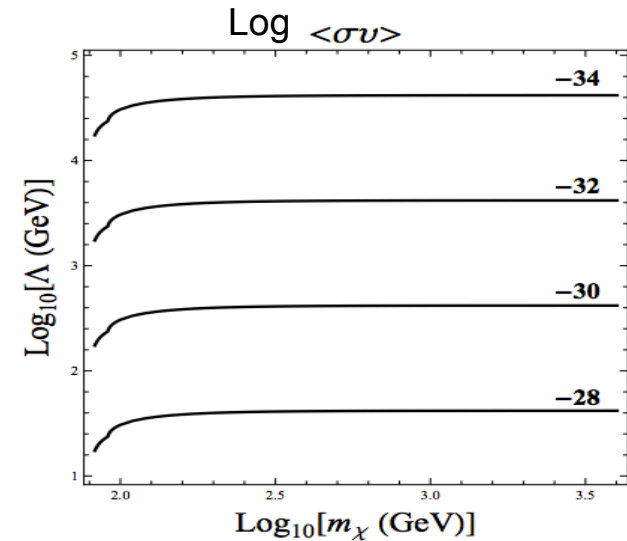
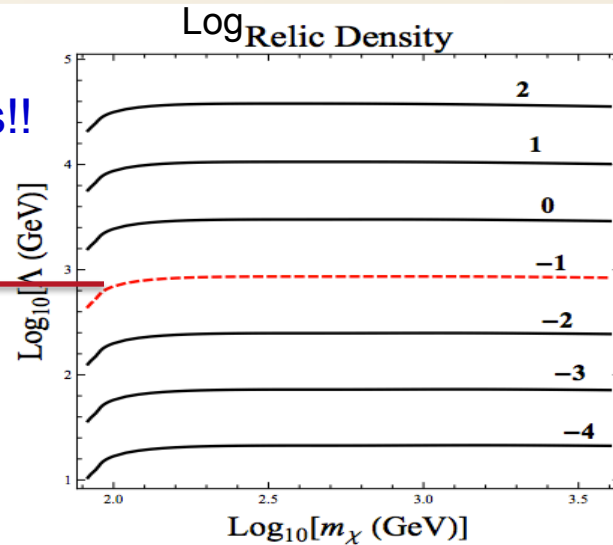
- Continuum γ -spectra from dwarf spheroidal galaxies for annihilation into SM final states
 - » (Fermi 10 MW dwarf spheroids data + Veritas Segue I)
- γ -ray lines for annihilations into $\gamma\gamma$ and γZ (Fermi)
 - » Assume NFW profile (bounds weaken $\sim 30\text{-}40\%$ w/ isothermal profile)
- Ratio of anti-proton/proton cosmic ray flux (Pamela)
 - » Assume NFW profile

Strength of indirect limits depends on velocity suppression of operator

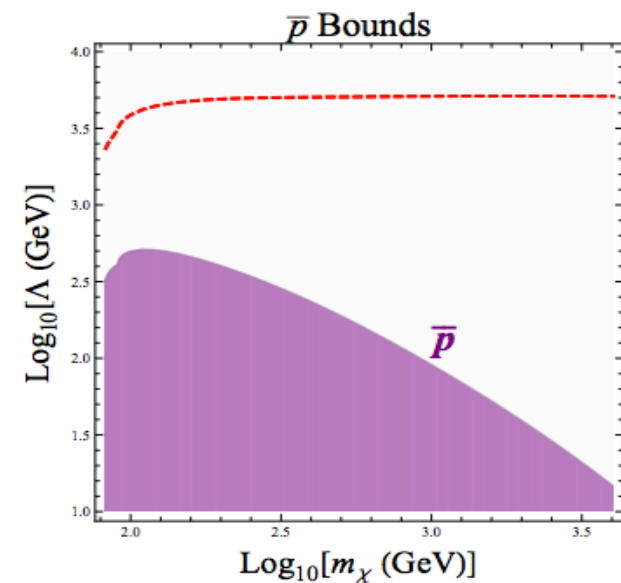
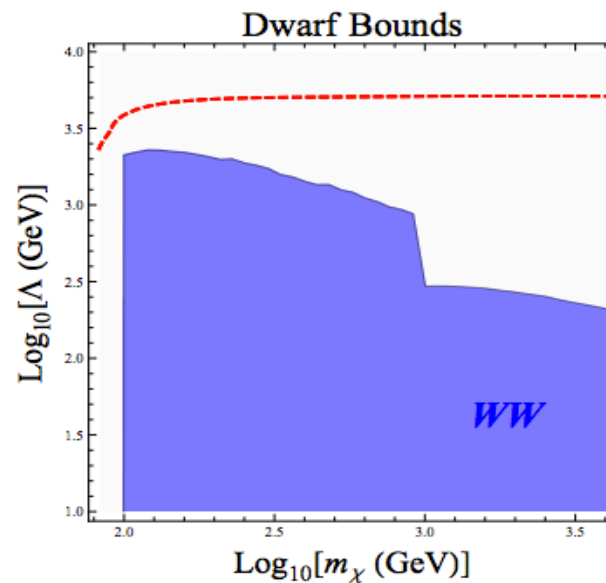
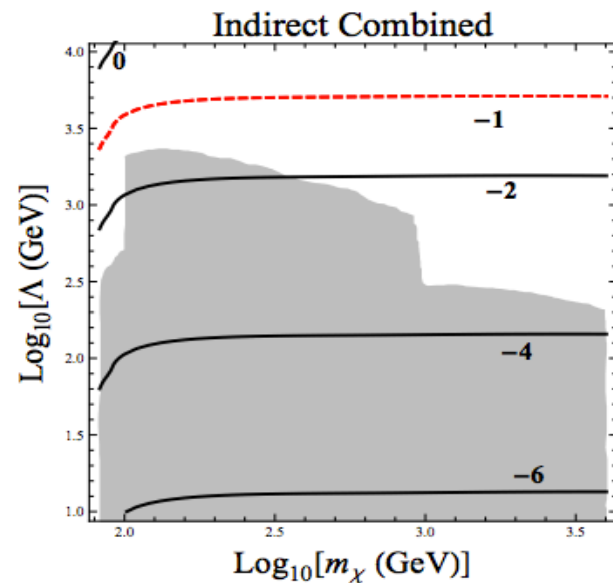
Indirect Search Results

Operator D5a
No indirect constraints!!

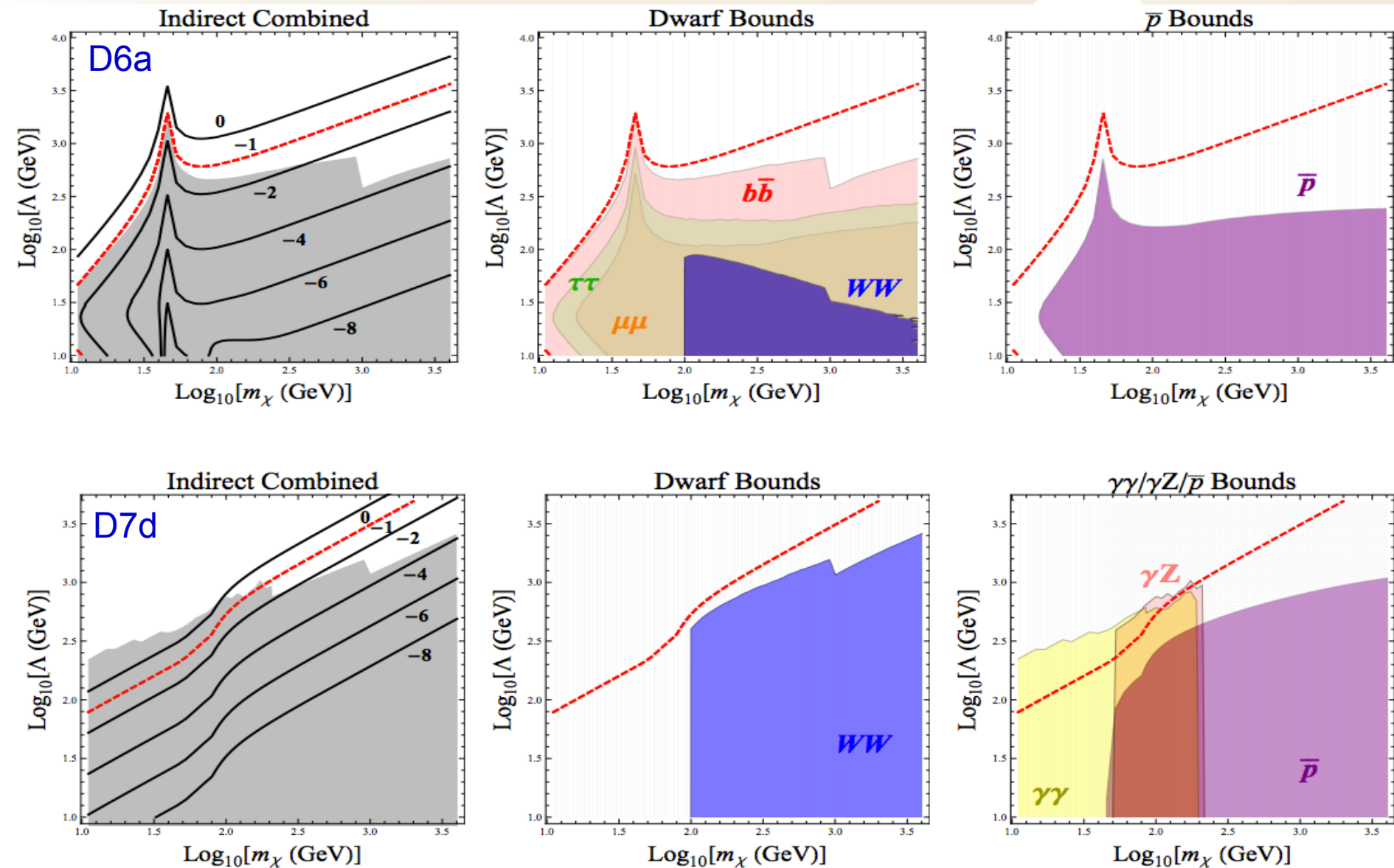
WMAP measurement
Thermal cosmology



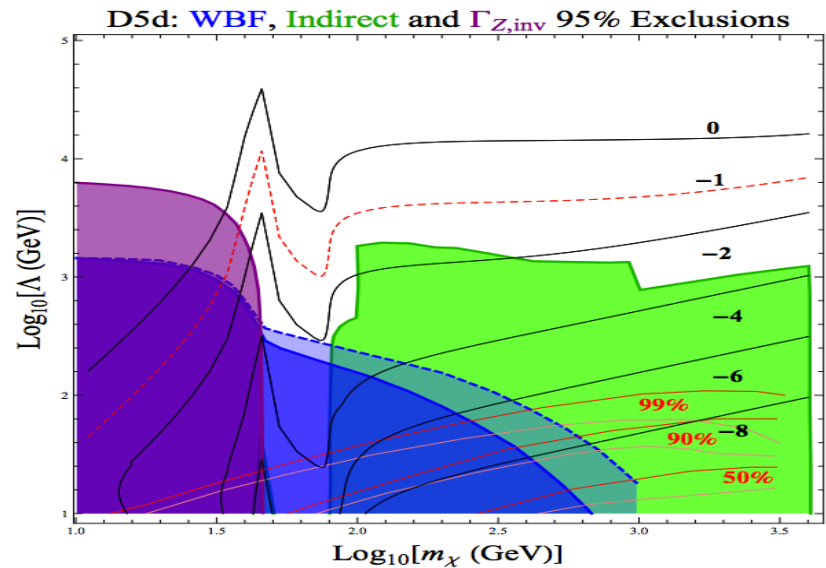
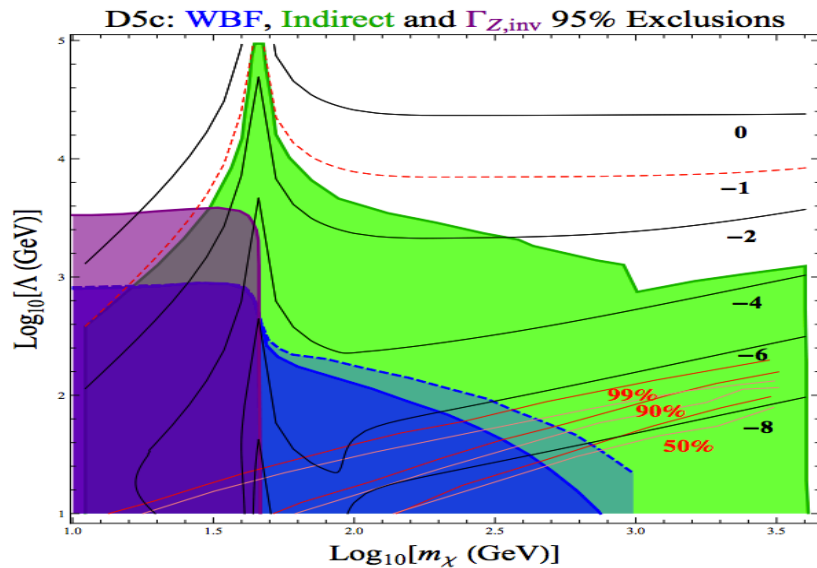
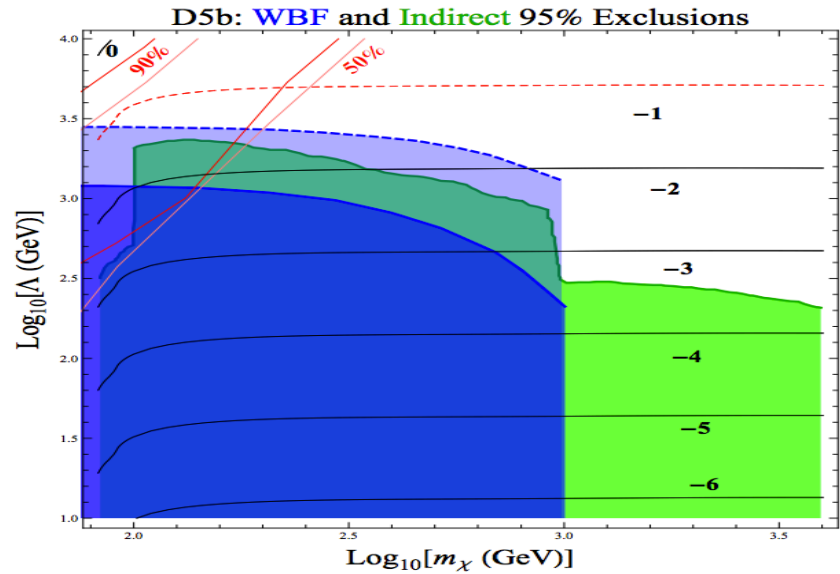
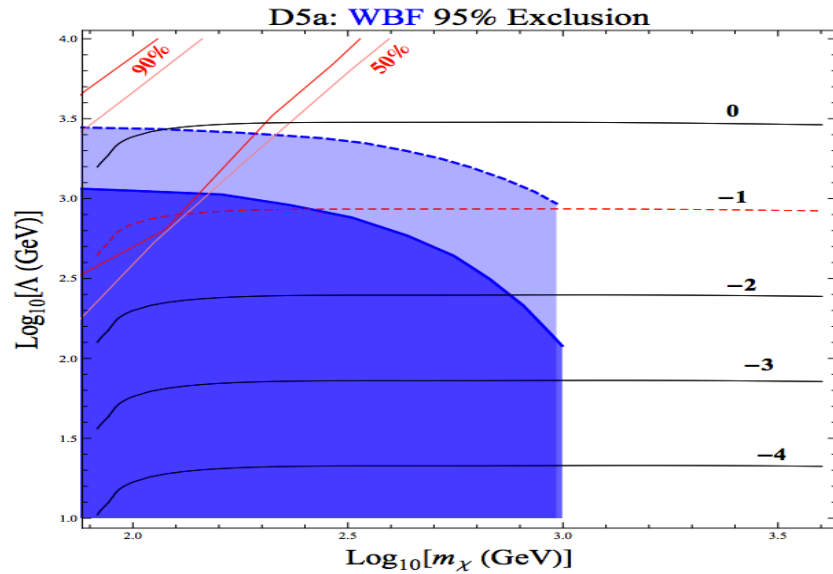
Operator D5b



Indirect Search Results

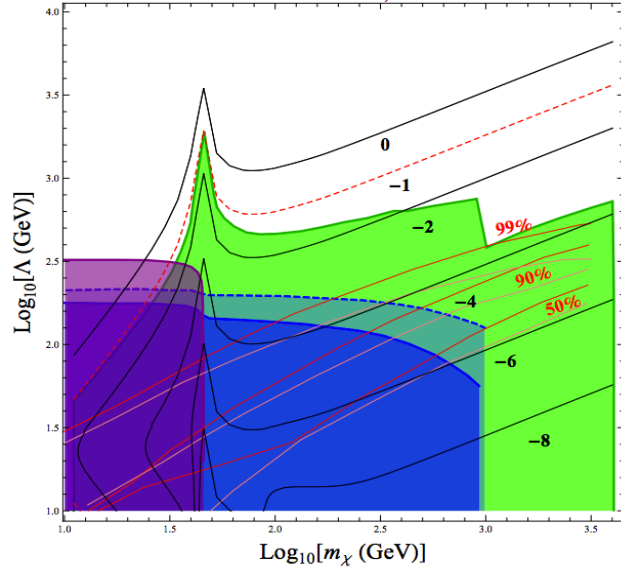


Combined Constraints: D5a-d

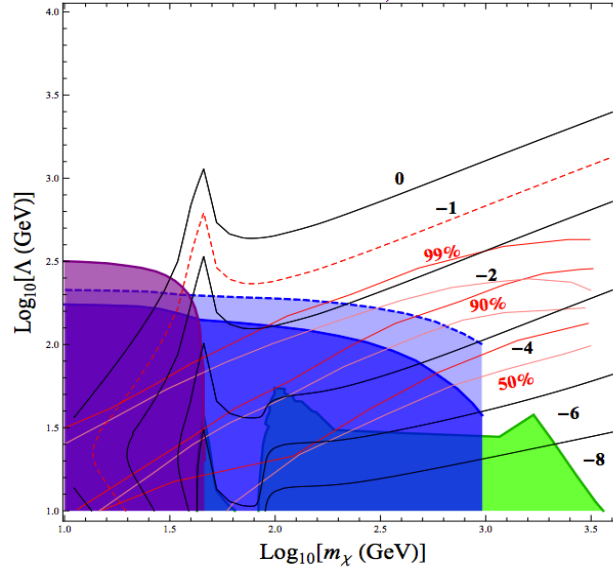


Combined Constraints: D6-D7

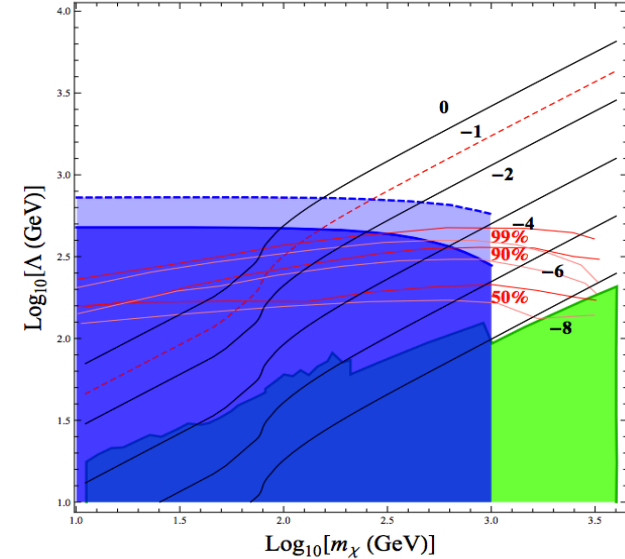
D6a: WBF, Indirect and $\Gamma_{Z,\text{inv}}$ 95% Exclusions



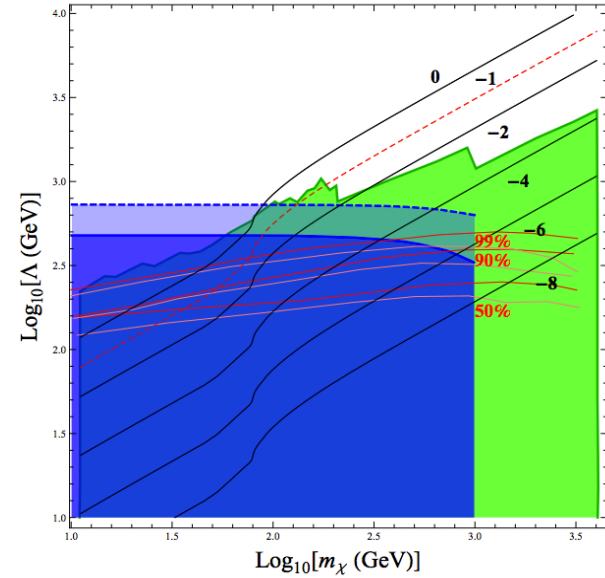
D6b: WBF, Indirect and $\Gamma_{Z,\text{inv}}$ 95% Exclusions



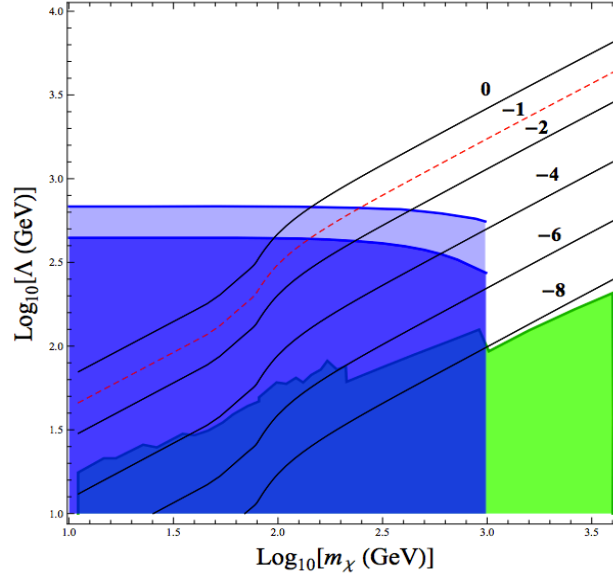
D7a: WBF and Indirect 95% Exclusions



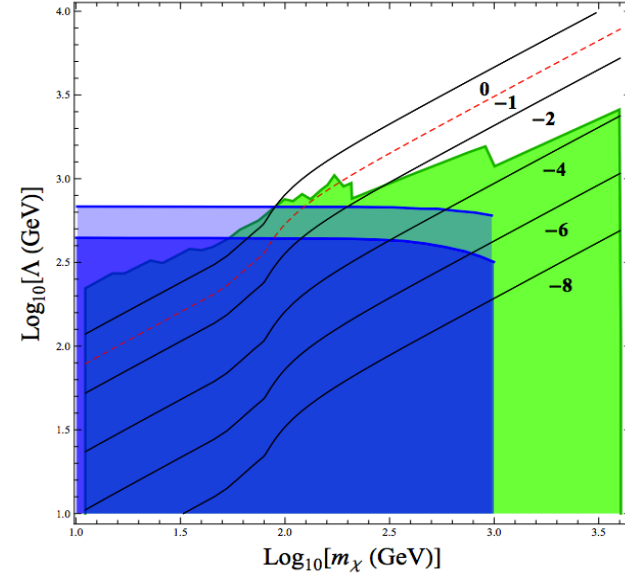
D7b: WBF and Indirect 95% Exclusions



D7c: WBF and Indirect 95% Exclusions



D7d: WBF and Indirect 95% Exclusions



Summary: EW boson probes of dark matter

- Provides complement to monjet and direct detection searches
- Interactions with EW gauge bosons probed for cut-offs in the weak-scale values, ~ 100 's GeV to few TeV
- Relatively light dark matter requires additional structure in UV theory to avoid overclosing the universe

BACKUP

Assumed weighted combination for collider subprocesses

$$\text{D5a :} \quad \frac{1}{\Lambda} \bar{\chi} \chi \left(\frac{Z^\mu Z_\mu}{2} + W^{+\mu} W_\mu^- + h.c. \right)$$

$$\text{D5b :} \quad \frac{1}{\Lambda} \bar{\chi} i \gamma_5 \chi \left(\frac{Z^\mu Z_\mu}{2} + W^{+\mu} W_\mu^- + h.c. \right)$$

$$\text{D5c :} \quad \frac{g_w}{\Lambda} \left(\bar{\chi} \sigma_{\mu\nu} t^3 \chi W^{3\mu\nu} + \frac{s_w}{c_w} \frac{Y}{2} \bar{\chi} \sigma_{\mu\nu} \chi B^{\mu\nu} \right)$$

$$\text{D5d :} \quad \frac{g_w}{\Lambda} \left(\bar{\chi} \sigma_{\mu\nu} t^3 \chi \widetilde{W}^{3\mu\nu} + \frac{s_w}{c_w} \frac{Y}{2} \bar{\chi} \sigma_{\mu\nu} \chi \widetilde{B}^{\mu\nu} \right)$$

$$\text{D6a :} \quad \frac{g_w}{\Lambda^2} \left(\bar{\chi} \gamma_\mu t^3 D_\nu \chi W^{3\mu\nu} + \frac{s_w}{c_w} \frac{Y}{2} \bar{\chi} \gamma_\mu D_\nu \chi B^{\mu\nu} \right)$$

$$\text{D6b :} \quad \frac{g_w}{\Lambda^2} \left(\bar{\chi} \gamma_5 \gamma_\mu t^3 D_\nu \chi W^{3\mu\nu} + \frac{s_w}{c_w} \frac{Y}{2} \bar{\chi} \gamma_5 \gamma_\mu D_\nu \chi B^{\mu\nu} \right)$$

$$\text{D7a :} \quad \frac{1}{\Lambda^3} \bar{\chi} \chi W^{a\mu\nu} W_{\mu\nu}^a$$

$$\text{D7b :} \quad \frac{1}{\Lambda^3} \bar{\chi} i \gamma_5 \chi W^{a\mu\nu} W_{\mu\nu}^a$$

$$\text{D7c :} \quad \frac{1}{\Lambda^3} \bar{\chi} \chi W^{a\mu\nu} \widetilde{W}_{\mu\nu}^a$$

$$\text{D7d :} \quad \frac{1}{\Lambda^3} \bar{\chi} i \gamma_5 \chi W^{a\mu\nu} \widetilde{W}_{\mu\nu}^a.$$