

Dark Matter with t-channel mediator A simple step beyond contact interaction

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Based on the work arXiv:1308.0592[hep-ph] in collaboration with Haipeng An and Lian-Tao Wang







Dark Matter Searching

- Weakly Interacting Massive Particle (WIMP) is one of the most popular dark matter candidates.
- The interaction between WIMP and the SM particles makes it detectable by satellites (indirect detection), underground detectors (direct detection), and colliders.





• Effective operator method: model independent, less free parameters.

 $\frac{C_{AB}}{\Lambda n} \bar{\chi} \Gamma^A \chi \bar{q} \Gamma^B q$









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• The cutoff scale ~ hundreds GeV.

CMS PAS EXO-12-048

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- The typical energy of the jets pass cut ~ hundreds GeV.
- The mass of dark matter ~ hundreds GeV.
- Effective energy of c.m.s ~ TeV.



• Effective operator method: model independent, less free parameters.



We need to go beyond the EFT!



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- Beyond EFT? New physics models?
- Supersymmetry, and etc :)
- More than 100 parameters and tens more particles...
- It is too far than a simple improvement of EFT.





- A simple step beyond the EFT.
- Adding a mediator to reduce the dimension of the effective operators in the theory to be smaller than 5.



The effective Lagrangian of a simple t-channel mediator theory can be written as

$$\mathcal{L}_{\chi} = \lambda_q \bar{\chi} \phi^* q + \text{h.c.}$$

 The quarks in the effective Lagrangian could be right-handed or lefthanded. In the MFV scenario,

$$\begin{aligned} \mathcal{L}_{\chi} &= \lambda_{Q} \bar{\chi} \mathbb{P}_{L} Q \phi_{Q}^{*} + \lambda_{u} \bar{\chi} \mathbb{P}_{R} u \phi_{u}^{*} + \lambda_{d} \bar{\chi} \mathbb{P}_{R} d \phi_{d}^{*} \\ &+ \frac{\lambda_{Qu}^{(1)} \bar{\chi} H \phi_{Q}^{*} Y_{u} \mathbb{P}_{R} u}{\Lambda} + \frac{\lambda_{Qd}^{(1)} \bar{\chi} \tilde{H} \phi_{Q}^{*} Y_{d} \mathbb{P}_{R} d}{\Lambda} \\ &+ \frac{\lambda_{Qu}^{(2)} \bar{Q} H Y_{u} \phi_{u} \mathbb{P}_{R} \chi}{\Lambda} + \frac{\lambda_{Qd}^{(2)} \bar{Q} \tilde{H} Y_{d} \phi_{d} \mathbb{P}_{R} \chi}{\Lambda} \end{aligned}$$

$$+ \text{h.c.} ,$$



Friday, September 20, 2013

 \bigcirc

d1)

d4)

a3)

 χ

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Direct detection

• Direct detection

$$\mathcal{O}_1 = \frac{\lambda^2}{2M_\phi^2} \bar{\chi}_L \gamma_\mu \chi_L \bar{q}_R \gamma^\mu q_R$$

- The direct detection cross section from this effective operator has been well studied.
- There are other operators induced by the t-channel mediator!

$$\mathcal{O}_2 = \frac{\alpha_S}{4\pi} G^{a\mu\nu} G^a_{\mu\nu} \chi^2, \qquad \mathcal{O}_3 = m_q \bar{q} q \chi^2$$

- Those operators are dim-7 and loop-induced from the t-channel mediator.
- The chiral symmetry enforces the Wilson coefficients to be proportional to the mass of the WIMP.

$$C_2 \sim \frac{\lambda^2 M_{\chi}}{M_{\phi}^4}, \quad C_3 \sim \frac{\lambda^2 m_t^2 M_{\chi}}{32\pi^2 M_{\phi}^2 v_{\rm ew}^2 M_h^2}$$

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- Dark matter production process: monojet + missing ET.
- New contributions:

(c) ---- from dim-7 operator
$$\mathcal{O}_8 = -\frac{ig_S\lambda^2}{M_{\phi}^4}T^a_{ij}A^a_{\mu}\left(\bar{\chi}P_Rq_j\right)\overleftrightarrow{\partial}^{\mu}\left(\bar{q}_iP_L\chi\right)$$

Higher suppressed by the mediator mass

No logarithm enhancement as the initial state QCD jet

(d1-d4) ---- WIMP-mediator associated production

- Two-body phase space enhancement
- High pT jet from heavy mediator decay







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(d1-d4) ---- WIMP-mediator associated production

- Two-body phase space enhancement
- High pT jet from heavy mediator decay
- The most recent monojet+missing ET constraint from the LHC is from CMS collaboration with 19.5 fb⁻¹ dataset from 8 TeV proton-proton collision.





- We use MadGraph5/MadEvent generate parton level event, shower it using PYTHIA6.4 and simulate the detector effects using PGS4 with anti-KT jet algorithm with a distance parameter of 0.5.
- Cuts:

Only one central jet $p_T > 110 \text{ GeV}, |\eta| < 2.4$

At most two jets s.t. $p_T > 30 \text{ GeV}, |\eta| < 4.5$

No isolated electron with $p_T > 10 \text{ GeV}, |\eta| < 1.44 \text{ or } 1.56 < |\eta| < 2.5$

No isolated muon with $p_T > 10 \text{ GeV}, |\eta| < 2.1$

 $\not\!\!\!\!/ E_T > 120~{\rm GeV}$

For events with a second jet, $\ \Delta \phi_{j_1 j_2} < 2.5$

 Events which pass those cuts are separated in seven signal regions according to the missing ET in the event.





• Checking our simulation using vector current contact effective operator.



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• Fix WIMP mass (5 GeV)







LHC phenomenology II: dijet

• The mediator is colored particle which can be produced by purely QCD process.







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LHC phenomenology II: dijet

- The total cross section is calculated using MadGraph5/MadEvent
- A typical value of the K-factor is smaller than 1.05. We will neglect it in our calculation.
- We compare the parton level cross section with the unfolded result of squark search given by the CMS collaboration.









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- We compare the constraints to the t-channel mediator model from direct detection and 8 TeV LHC.
- The difference between the t-channel mediator model and the effective operator approximation is shown clearly in the figures.



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Conclusion

- We study a simplified t-channel UV completion model where the interaction between DM and SM particles are mediated by colored mediators couples to the DM particle and the right-handed quarks.
- In this scenario, if the DM particle is Dirac, the dominant direct detection signal is SI, and the monojet+missing ET search starts to be sensitive to the interesting parameter space in the light DM region.
- In the case that the DM particle is Majorana, the dominant direct detection signal is SD, and the monojet+missing ET signal is stronger in the region that DM mass is smaller than a hundred GeV, and dijet +missing ET is more significant for heavier dark matter.
- Additional annihilation processes are needed for Dirac DM to give correct relic abundance. For Majorana DM, there is still allowed parameter region which can gives correct relic abundance.





Thank you!