Gravitational waves from BH-NS binaries: Effective Fisher matrices and parameter estimation using higher-harmonics

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> September 28, 2012 Midwest Gravity Meeting

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Motivation The Fisher matrix Effective Fisher matrix

Motivation

- At leading order, the GW signal from a binary oscillates at twice the orbital frequency. Other harmonics of the orbital frequency enter at higher post-Newtonian (PN) order
- Are higher harmonics/amplitude corrections (ACs) important for parameter estimation, particularly for BH-NS systems?
- This can be addressed with Bayesian inference techniques like MCMC or nested sampling runs, but these are expensive
- We would like to get some insights with simpler, cheaper analytic methods
- The Fisher matrix is a standard approach, but can be perilous

Fisher matrix assumptions

- The Fisher matrix formalism is widely used because it is much cheaper than a full Bayesian analysis
- However, it makes several assumptions that are rather optimistic for ground-based GW data analysis:
 - High SNR
 - Stationary, Gaussian noise
 - Signal depends linearly on all of its parameters
- What could possibly go wrong?

Motivation The Fisher matrix Effective Fisher matrix

The Fisher matrix

• Those assumptions are meant to justify approximating the likelihood (of an observed signal given some proposed waveform parameters) as a multivariate Gaussian in some neighborhood of the true parameters

$$\mathcal{L} \simeq e^{-\Gamma_{ij}\,\delta\lambda^i\,\delta\lambda^j/2}$$
 where $\Gamma_{ij} = \left\langle \frac{\partial h}{\partial\lambda_i} \middle| \frac{\partial h}{\partial\lambda_j} \right\rangle$

is the Fisher matrix.

• The covariance matrix, the inverse of the Fisher matrix, will estimate the expected errors and correlations of the parameter measurements

$$\begin{split} \Sigma &= \Gamma^{-1} \\ \sigma_i &= \sqrt{\Sigma_{ii}} \\ c_{ij} &= \frac{\Sigma_{ij}}{\sqrt{\Sigma_{ii}\Sigma_{jj}}} \end{split}$$

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Log Likelihood

• The likelihood that $h(\lambda')$ will produce a signal $s = h(\lambda) + n$ is:

$$\mathcal{L}(s|\lambda') = \exp\left[-\frac{1}{2}\langle s - h(\lambda') | s - h(\lambda') \rangle\right]$$

• Expanding *s* and neglecting the noise contribution to the likelihood, we get:

$$\mathcal{L} = \exp\left[-\frac{1}{2}\left(\left\langle h(\lambda) \big| h(\lambda) \right\rangle + \left\langle h(\lambda') \big| h(\lambda') \right\rangle - 2\operatorname{Re}\left\langle h(\lambda) \big| h(\lambda') \right\rangle\right)\right]$$

• Assuming $\rho^2 = \langle h(\lambda) | h(\lambda) \rangle = \langle h(\lambda') | h(\lambda') \rangle$,

$$\ln \mathcal{L} = -\rho^2 \left(1 - \frac{\operatorname{Re}\langle h(\lambda) | h(\lambda') \rangle}{\sqrt{\langle h(\lambda) | h(\lambda) \rangle \langle h(\lambda') | h(\lambda') \rangle}} \right)$$

Motivation The Fisher matrix Effective Fisher matrix

Ambiguity function

• The ambiguity function, the normalized overlap between different points in the parameter space, is

$$P(\lambda,\lambda') = \frac{\operatorname{Re}\langle h(\lambda) | h(\lambda') \rangle}{\sqrt{\langle h(\lambda) | h(\lambda) \rangle \langle h(\lambda') | h(\lambda') \rangle}}$$

• Computing the ambiguity function allows us to map out the log likelihood, and it can be related to the Fisher matrix

$$\ln \mathcal{L} = -\rho^2 \left(1 - P(\lambda, \lambda') \right) \simeq -\Gamma_{ij} \, \delta \lambda^i \, \delta \lambda^j / 2$$

• The (normalized) *effective* Fisher matrix is obtained by fitting the ambiguity function with a multivariate quadratic

$$P_{\rm fit}(\delta\lambda) = 1 - \hat{\Gamma}_{ij}^{\rm eff} \delta\lambda^i \,\delta\lambda^j / 2$$

• The range over which you fit is determined by the expected SNR of your signal: $1 - P \le 1/\rho^2$

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Advantages of the Effective Fisher matrix



1.000 0.998 0.996 0.994 ٩, 0.992 - Ambiguity function 0.990 Fitting function at P>0.99 Fitting function at P>0.999 0.988 0.986 -0.003 -0.002 -0.001 0.000 0.001 0.002 0.003 δM_c

• Plotting the ambiguity function will tell you if the Fisher matrix approximation is at all reasonable • Fitting to a scale set by the SNR makes it robust against fine-scale structure that is unobservable at that SNR

Non-spinning results Spin-aligned results Precessing results

Fiducial Binary

- We study a $(10 + 1.4) M_{\odot}$ BH-NS binary ($M_c = 2.99, \eta = 0.1077$) in initial LIGO
- Use precessing PN waveforms (SpinTaylorT4) with 0PN or 1.5PN amplitude
- Non-spinning/Spin-aligned: *L* has an inclination of π/4 to line of sight *N*; BH has spin 0 or 1 along *L*; NS is non-spinning
- Precessing: *J* is inclined to line of sight by $\theta_{NJ} = \pi/4$; precession cone has an opening angle $\beta_{JL} = \pi/4$; *L* is initially either along *N* or perpendicular to it



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Non-spinning results Spin-aligned results Precessing results

Non-spinning results





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Spin-aligned results

- Non-spinning: ACs do not improve measurement of intrinsic parameters
- Spin-aligned: ACs give a modest improvement to intrinsic parameter errors
- In both cases, extrinsic parameters are unmeasurable without amp. cor., and are measureable (but with large errors) with ACs



amp. order	0PN	1.5PN
$\Delta M_c/M_c$	1.08%	.795%
$\Delta \eta / \eta$	18.7%	12.7%
$\Delta \chi$	0.149	0.108
$\Delta \iota_{LN}$		0.282
$\Delta \phi_{ m ref}$	—	0.821

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Precessing results

- The errors are smaller for precessing binaries than for spin-aligned
- The gains from ACs become less pronounced, however
- The divide between intrinsic and extrinsic parameters becomes blurred, as angles describing the precession cone correlate to both

amp. order	0PN	1.5PN
$\Delta M_c/M_c$.208%	.195%
$\Delta \eta / \eta$	6.14%	4.91%
$\Delta \chi$	0.0495	0.0421
$\Delta \beta_{JL}$	0.0241	0.0210
$\Delta heta_{JN}$	0.117	0.113
$\Delta \alpha_{JL}$	0.187	0.191
$\Delta \phi_{ m ref}$	—	_

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Caveats

- We see non-quadratic behavior in the precessing case for some orientations (but not others)
- We use a complex inner product s.t. Re $\langle h|h'\rangle =$ $(h_+|h'_+) + (h_\times|h_\times)$ corresponding to an idealized network.
- The results should lie between those of a single-IFO and a real multi-detector network
- We cautiously believe the trends seen here and are working to confirm with MCMC



Conclusions

- The effective Fisher matrix allows one to see whether the Fisher matrix formalism is likely to be valid (and for which parameters)
- It is also robust against small-scale structure that can be problematic for the standard Fisher matrix
- Higher harmonics will likely have little effect on the measurement of intrinsic parameters, but they can improve the measurement of extrinsic parameters
- For non-precessing systems, there is a clear separation of intrinsic and extrinsic parameters, but this becomes blurred for precessing
- For more details, see arXiv:1209.4494
- Full MCMC runs underway to get more precise results and confirm the trends seen here