

Exploring the low mass frontier in dark matter direct detection

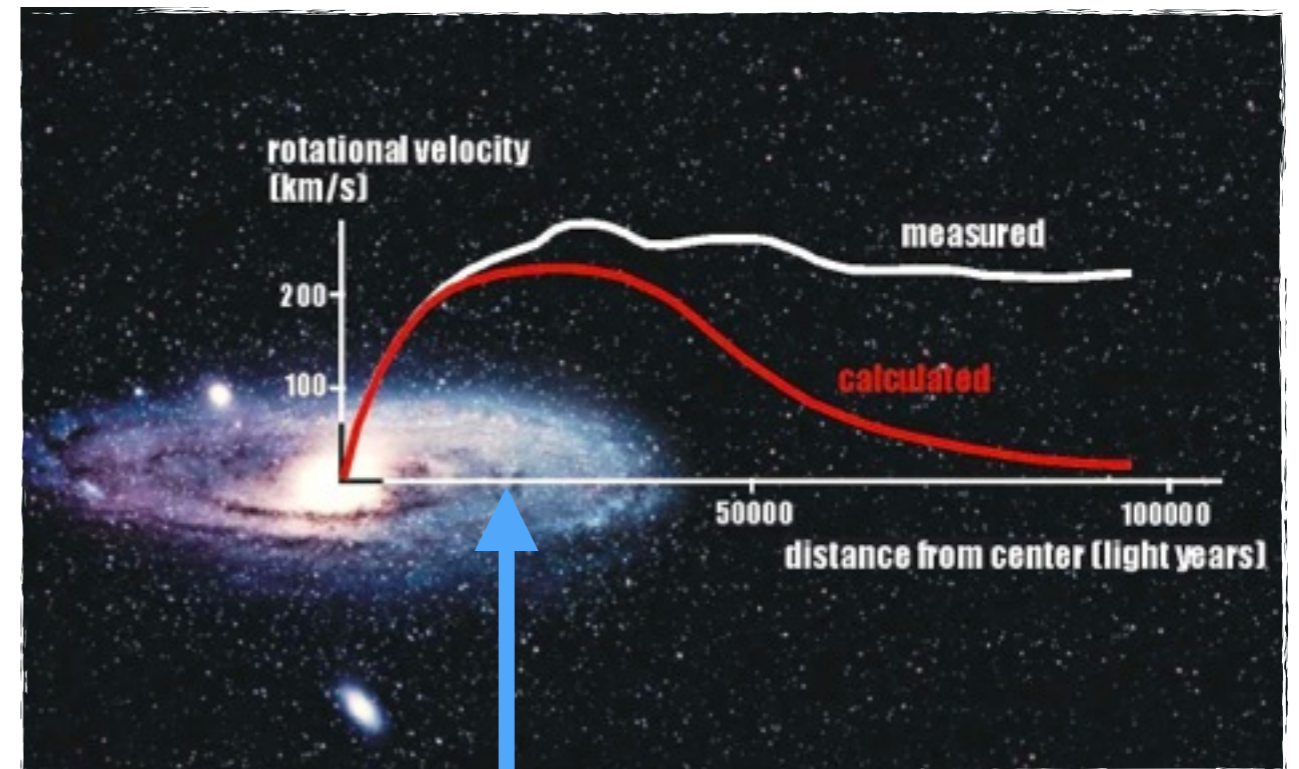
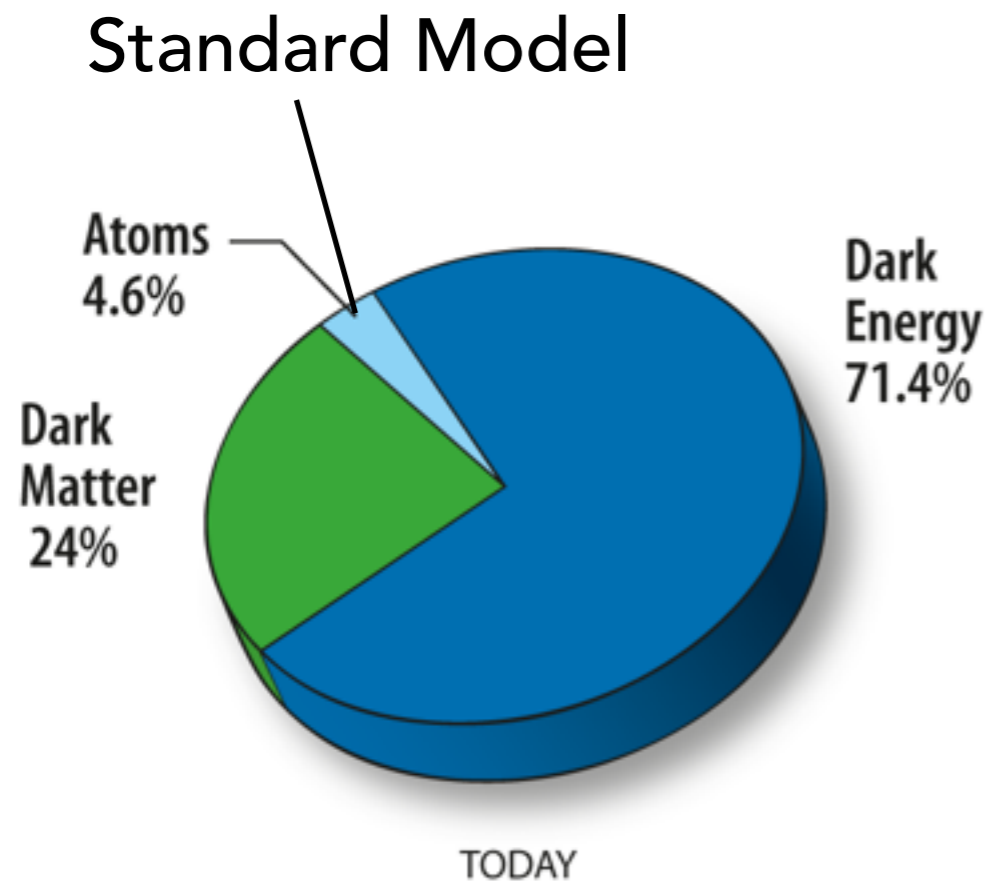
Tongyan Lin
UC Berkeley / LBNL

KICP Particle Cosmo Workshop
October 13, 2016

Outline

1. Why light dark matter?
2. Direct detection of light dark matter
3. Low-threshold targets for $M = \text{meV} - \text{MeV}$
 - a. absorption in semi/super-conductors
 - b. scattering in liquid helium

Dark matter puzzles



Solar neighborhood: $\rho_x \approx \frac{0.4 \text{ GeV}}{\text{cm}^3}$

Some new particle physics is needed

Directions



- Standard Model has other unresolved questions - use as a guide for dark matter.
- Look more broadly for a dark sector. Generically, this would be as rich as the Standard Model.

WIMP dark matter

WIMP (Weakly Interacting Massive Particle)

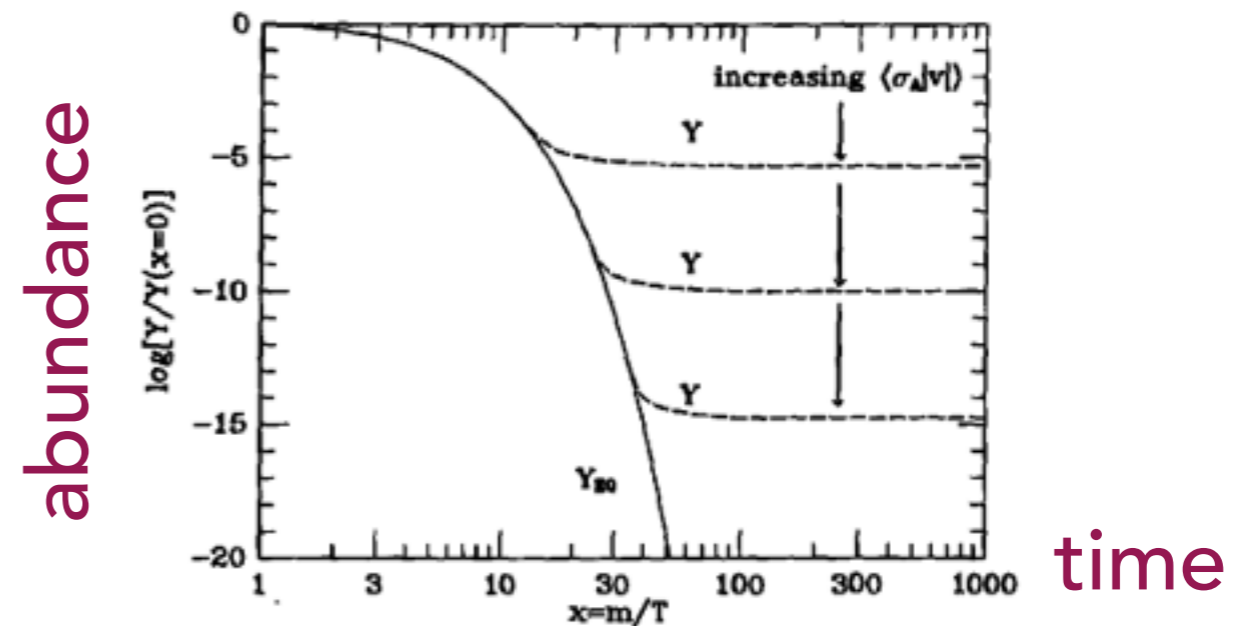
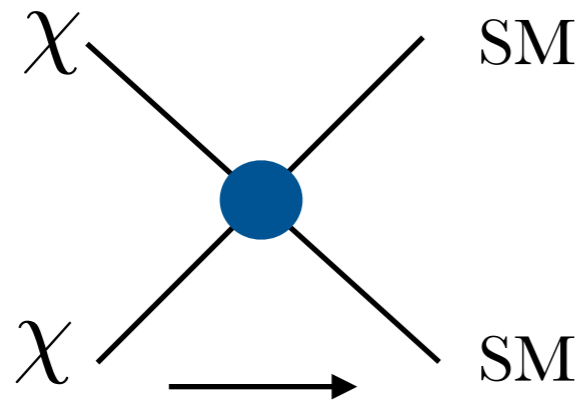


Fig. 5.1: The freeze out of a massive particle species. The dashed line is the actual abundance, and the solid line is the equilibrium abundance. Kolb & Turner

Relic abundance of dark matter today set by weak-scale interactions in the early universe

$$M = 10 \text{ GeV} - 10 \text{ TeV} ?$$

Mass scale of dark matter?



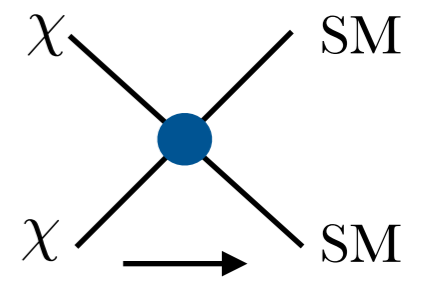
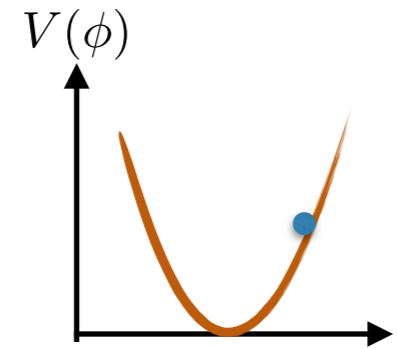
Light bosonic DM

WIMP

Many recent theory ideas

+

Many proposals for detection



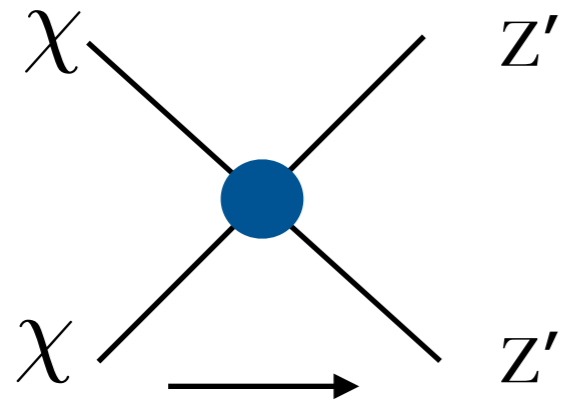
coherent light boson searches
(e.g. ADMX, CASPEr)



Nuclear recoil in direct detection

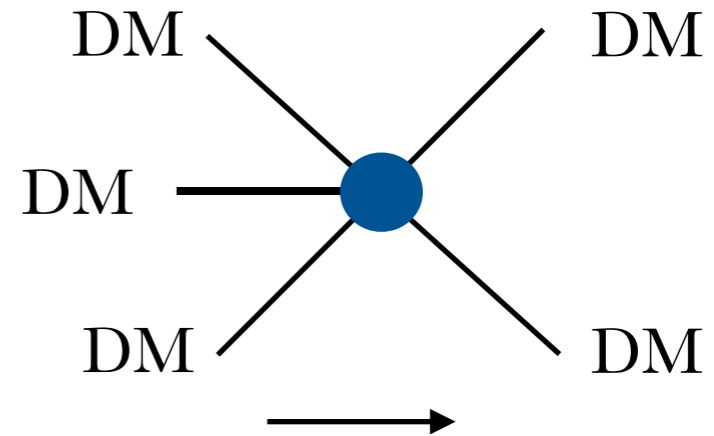
sub-GeV dark matter

light mediators



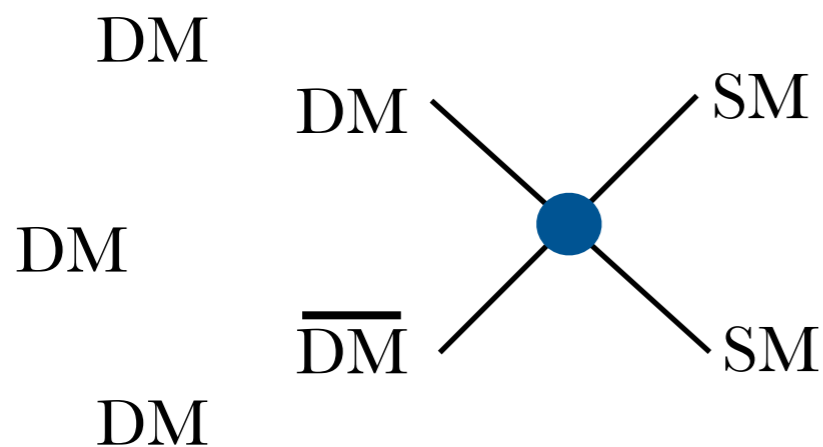
e.g. Fayet, Pospelov et al. 2007

3 to 2 annihilation, SIMPs



Hochberg et al. 2014, 2015

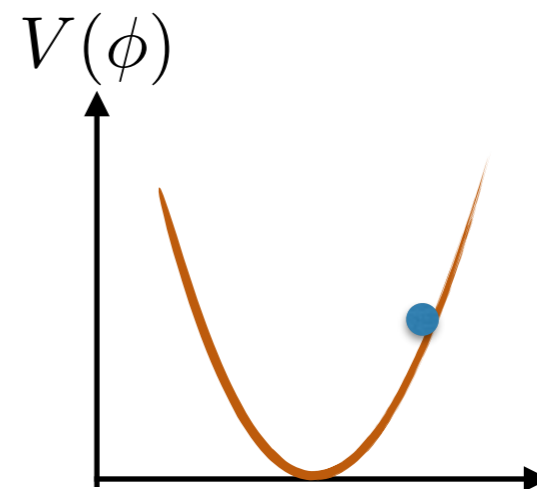
asymmetric



e.g. Nussinov, Barr, Kaplan et al.

7

sub-keV non-thermal



e.g. Pospelov et al. 2008, Arias et al. 2012

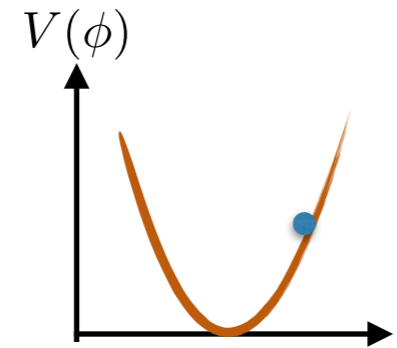
Mass scale of dark matter?



Light bosonic DM

Light thermal
Asymmetric
SIMP

WIMP

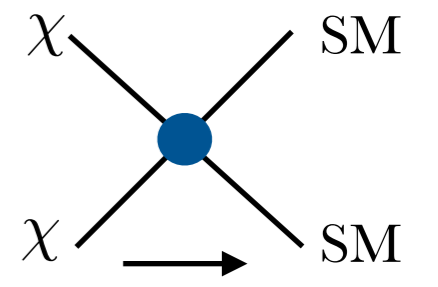


helium

semiconductors

superconductors

see also: graphene, scintillators,
chemical bonds, ...

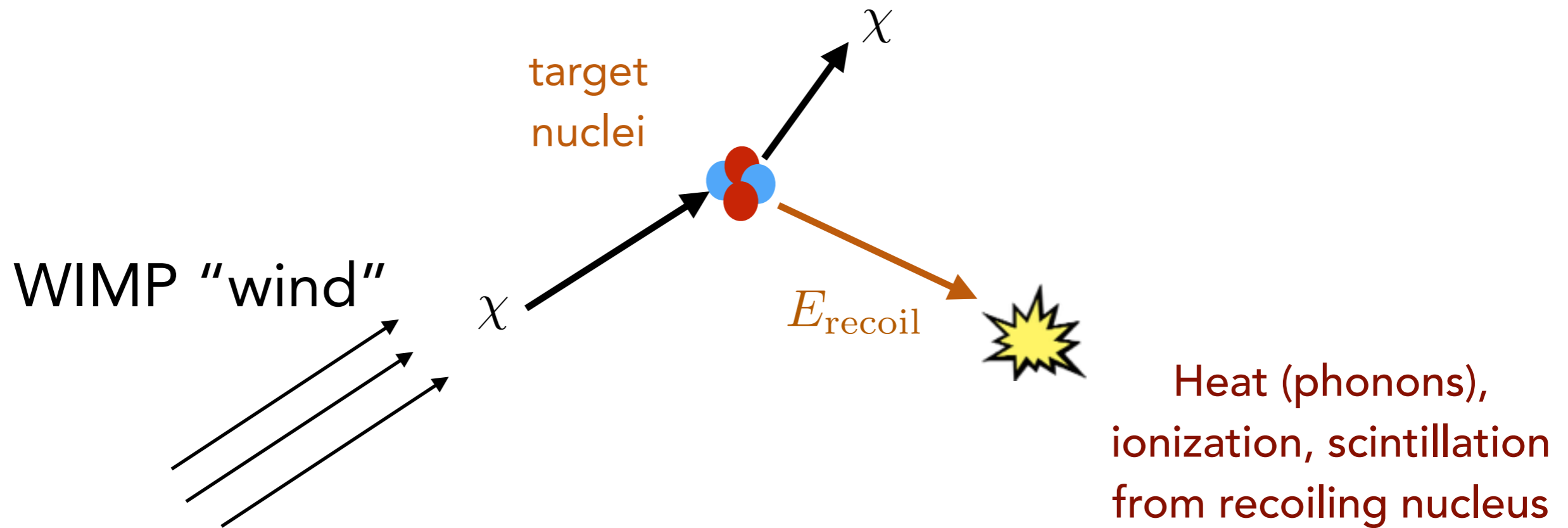


coherent light boson searches
(e.g. ADMX, CASPEr)

Nuclear recoil in
direct detection

2. Direct detection of light dark matter

Direct detection of WIMPs

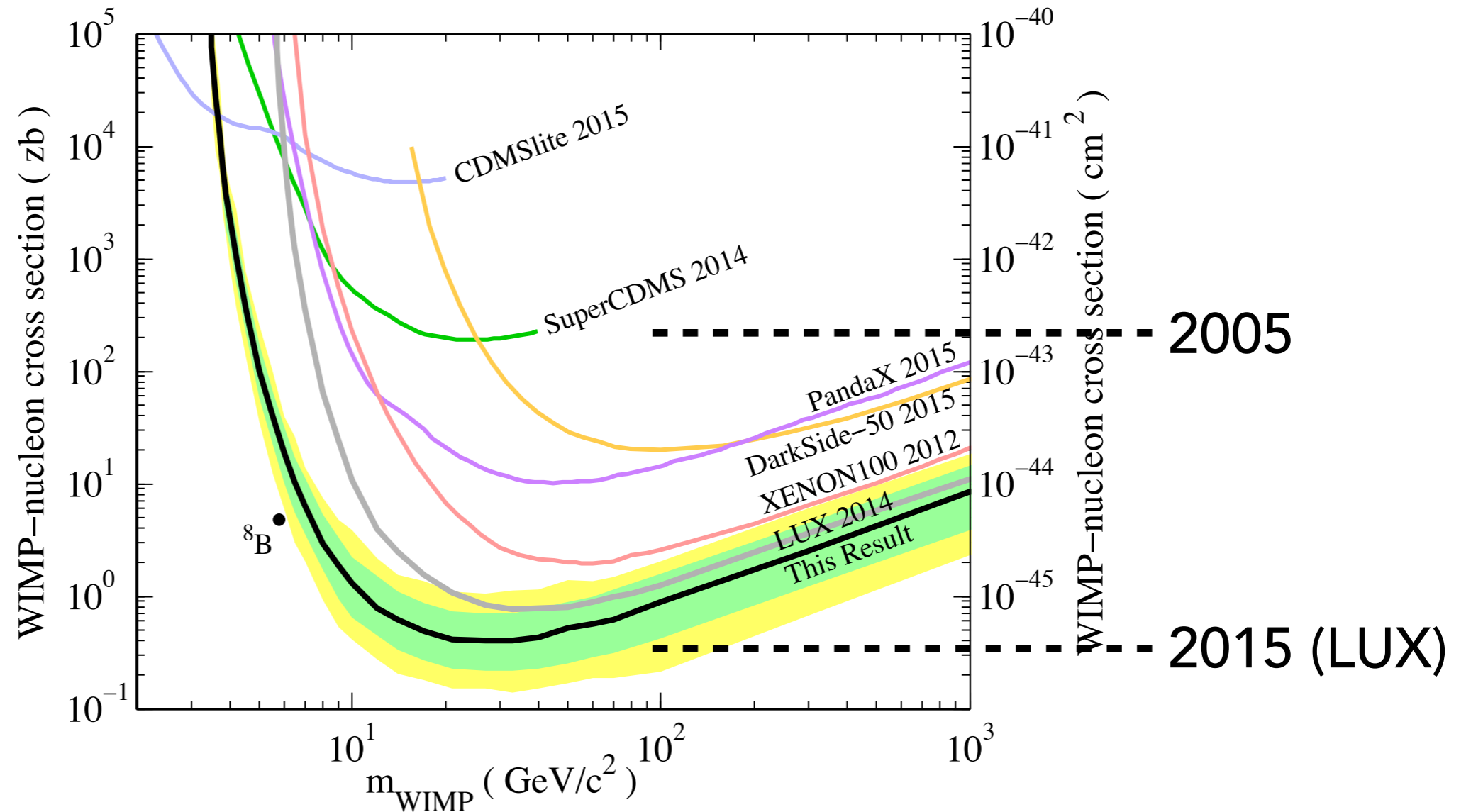


Energy deposited from WIMP in nuclear recoil:

$$E_R \sim \frac{\mu_{\chi N}^2 v^2}{m_N} \sim 1 - 100 \text{ keV}$$

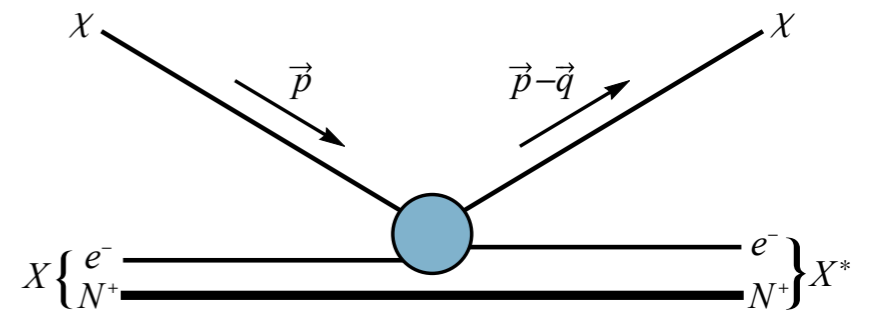
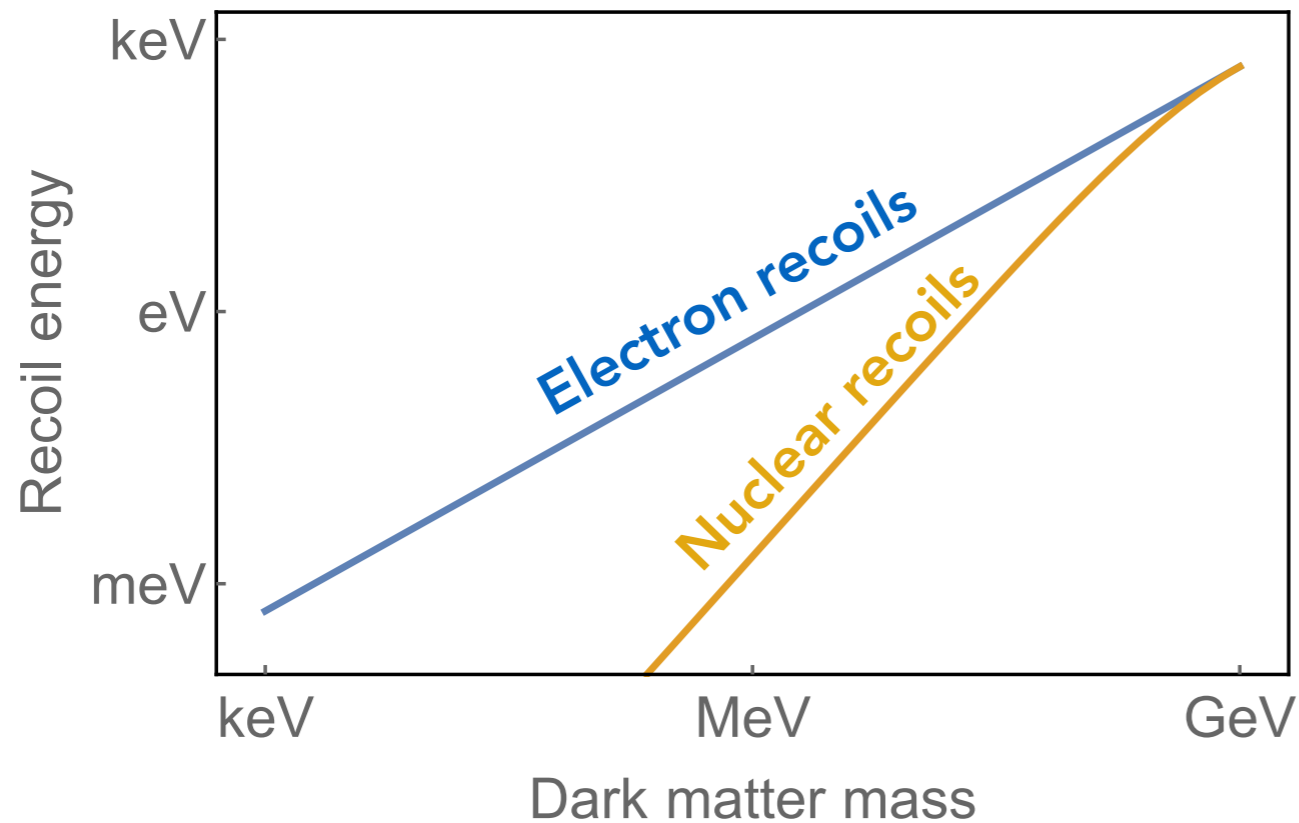
Typical dark matter velocity $v \sim 10^{-3}$

Direct detection of WIMPs



Typical threshold in experiment: > 1 keV recoil energy

Detecting light dark matter



Look for electron recoils

- Below ~ 1 GeV, inefficient energy transfer to nuclei.
- Idea applied in: Xenon, semiconductors...

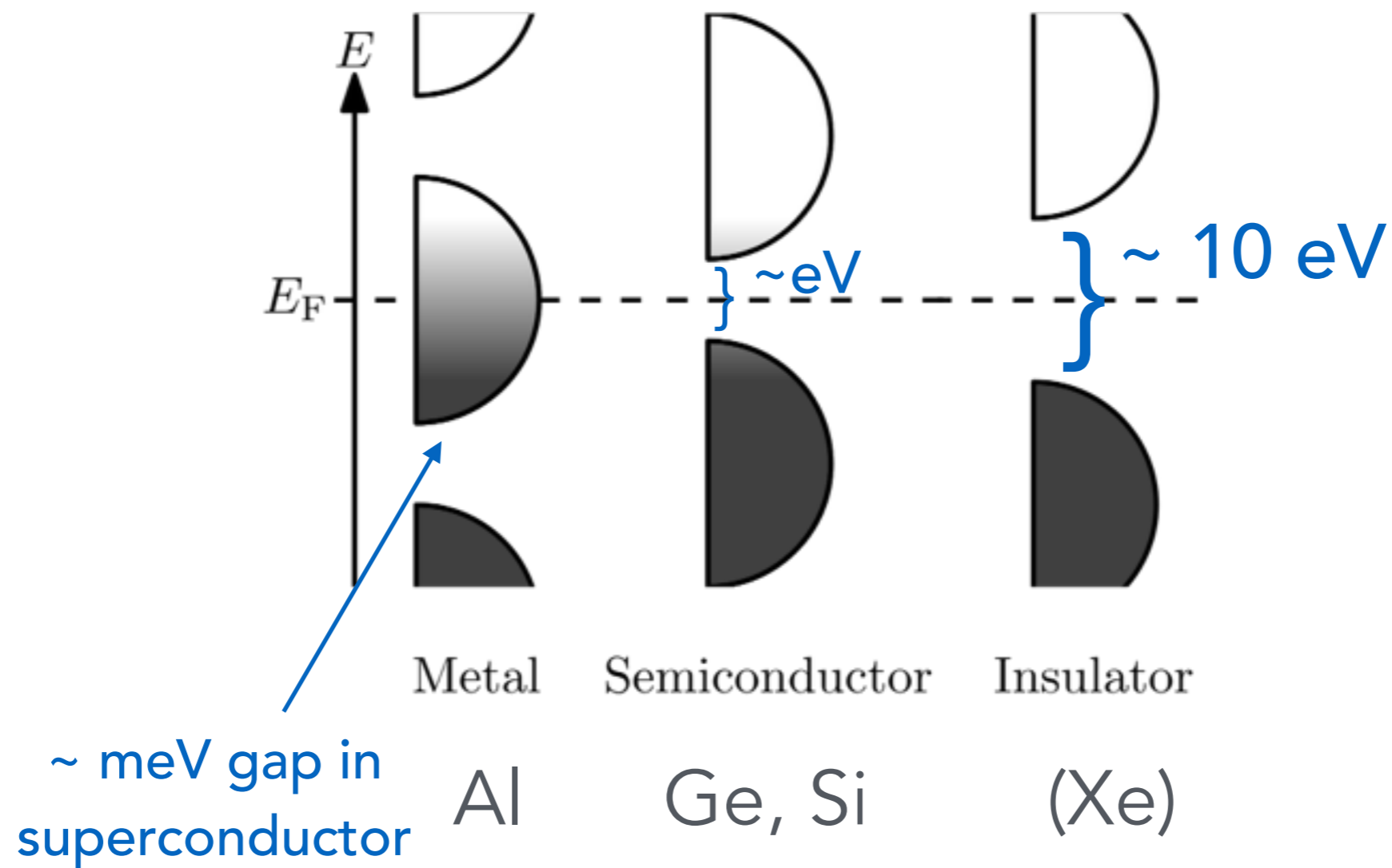
Electron recoils

- Electron interactions already constrained with Xenon10 low-threshold analysis, search for DM ionization signal

$$E_{th} \sim 12 \text{ eV}$$

Electron recoils

Electronic band structure



Gapless excitations

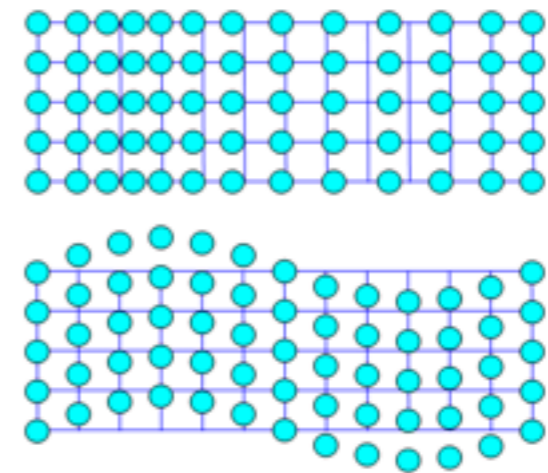
Long wavelength acoustic phonons have linear dispersion:

$$\Omega = c_s |\vec{Q}|$$

speed of sound

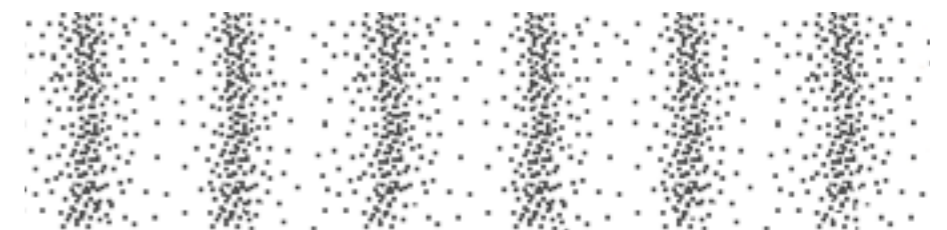
- Lattice vibrations in a solid

$$c_s \sim 10^{-5} \text{ in aluminum}$$



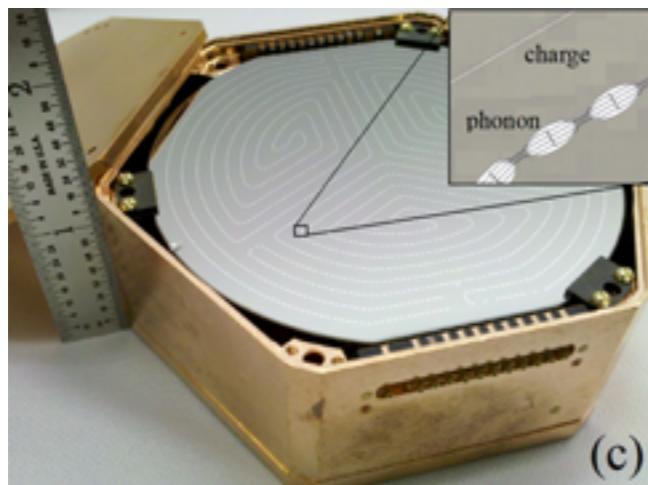
- Density perturbations in a liquid

$$c_s \sim 10^{-6} \text{ in helium}$$



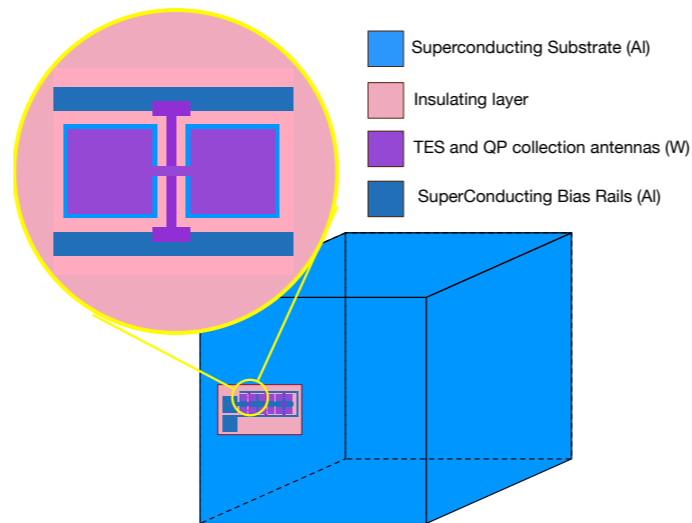
Low-threshold targets

Semiconductors



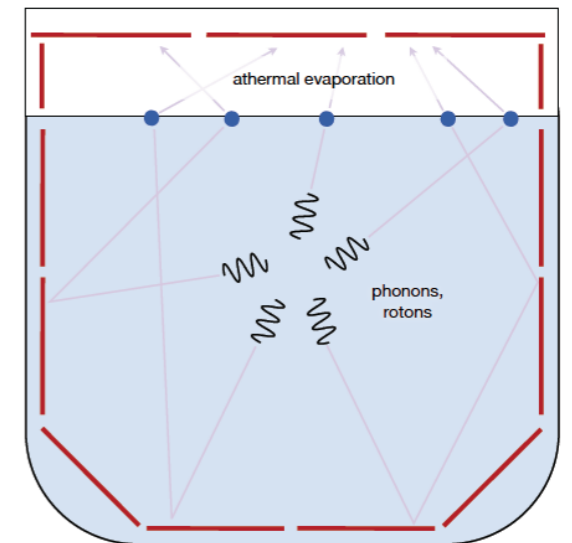
- $E_{th} \sim \text{eV}$, electron+phonon
- Reach MeV DM (scattering)
- Used in DAMIC, SuperCDMS, ...

Superconductors



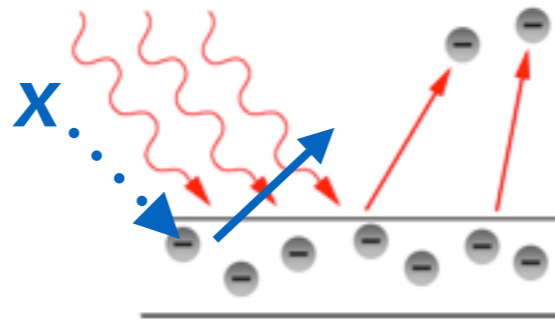
- $E_{th} \sim \text{meV}$, electron+phonon
- Reach keV DM (scattering)

Helium



- $E_{th} \sim \text{meV}$, phonons
- Reach keV DM (scattering)

3a. Absorption of light bosonic dark matter



with:

Yonit Hochberg

Kathryn Zurek

1604.06800, 1608.01994

sub-keV bosonic dark matter

- Candidates:
 - Hidden photon
 - Pseudoscalar (axion)
 - Scalar

- Coherent field below $m \sim \text{eV}$

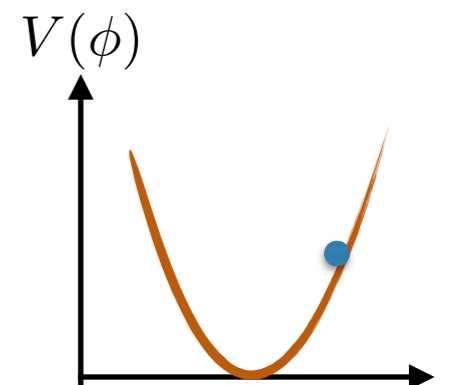
Local DM density: 0.4 GeV/cm^3 $\lambda_{\text{dB}} \sim \frac{2\pi}{m_{\text{DM}}v}$ $v \sim 10^{-3}$

Occupation number is high: $\frac{\rho_{\text{DM}}}{m_{\text{DM}}} \gg \lambda_{\text{dB}}^{-3}$

- Non-thermal relic abundance, e.g. “misalignment”

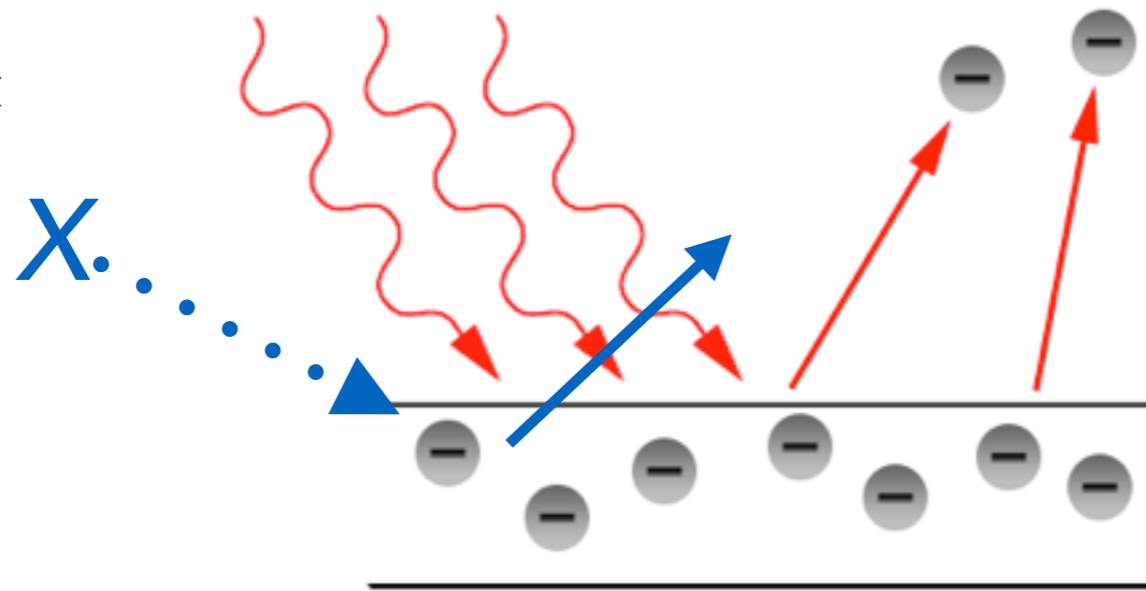
$$\rho_{\text{DM}} = \frac{1}{2} m_{\text{DM}}^2 \phi_0^2$$

ϕ_0 — field amplitude today



Absorption

Photoelectric effect:



absorb all of the energy of incoming dark matter

Absorption from halo

- mono-energetic
- doesn't require coherent field

Solar emission

- \sim keV energies
- "axio-electric" effect

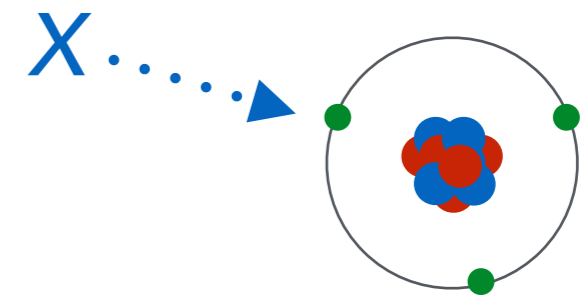
DM absorption in materials

Typical electron in material has $k \sim \text{keV}$.

Need momentum transfer $q \sim 100 m_X$ in absorption.

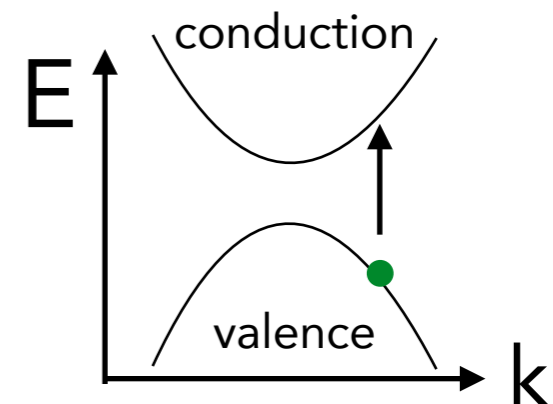
- **Atomic ionization**

- Ex: Xenon, ionization energy = 12 eV



- **Absorption on a crystal/solid**

- Rely on band structure
- Phonon emission

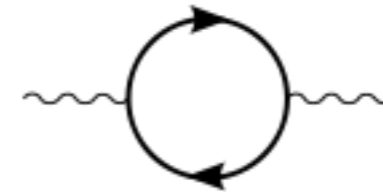


Relate the DM absorption rate to **photon** absorption rate:

$$\langle n_e \sigma_{\text{abs}} v \rangle_{\text{DM}} \propto \langle n_e \sigma_{\text{abs}} v \rangle_{\gamma} = \sigma_1 \quad \leftarrow \text{(photon absorption, conductivity } \sigma)$$

Conductivity

- Polarization tensor in medium:



$$\Pi^{\mu\nu} = e^2 \langle J_{\text{EM}}^{\mu\dagger}, J_{\text{EM}}^{\nu} \rangle \quad \Pi^{\mu\nu}(\vec{q}, \omega) = \Pi(\omega) \sum_{i=1,2} \epsilon_i^{T\mu} \epsilon_i^{T*\nu} + \Pi(\omega) \epsilon^{L\mu} \epsilon^{L\nu}$$

- Related to optical conductivity

$$\vec{J} = \hat{\sigma} \vec{E}$$

↑
Conductivity

$$\Pi(\omega) \approx -i\hat{\sigma}\omega$$

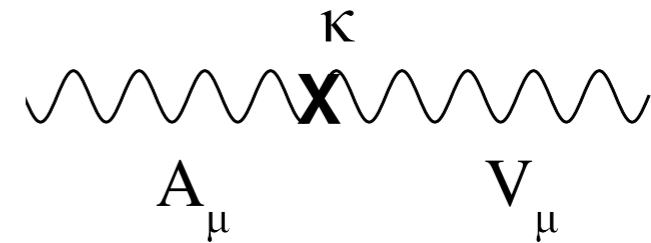
Real part gives effective mass , imaginary part gives absorption:

$$\text{Re } \Pi(\omega) \approx \omega_p^2 = \sigma_2 \omega \quad -\frac{\text{Im } \Pi(\omega)}{\omega} = \sigma_1 = \langle n_e \sigma_{\text{abs}} v \rangle_{\gamma}$$

Hidden photon dark matter

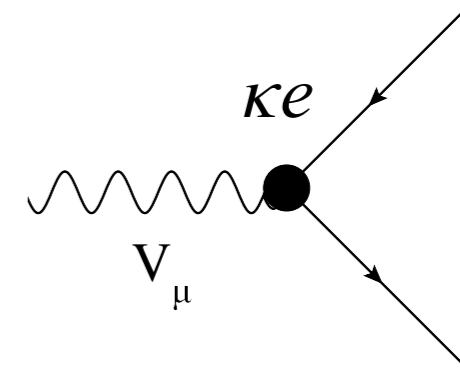
Kinetic mixing in vacuum:

$$\mathcal{L} \supset -\frac{\kappa}{2} F_{\mu\nu} V^{\mu\nu}$$



Matter coupling:

$$A_\mu \rightarrow A_\mu - \kappa V_\mu \longrightarrow \kappa e V_\mu J_{EM}^\mu$$



Absorption rate of halo DM

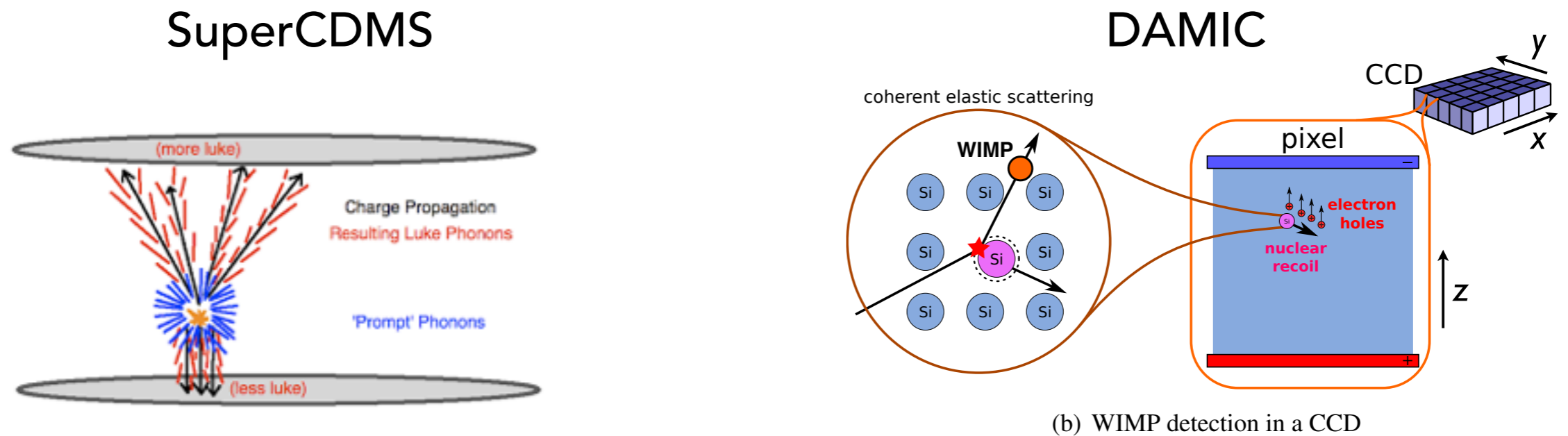
$$R = \frac{1}{\rho} \frac{\rho_{DM}}{m_{DM}} \kappa_{eff}^2 \sigma_1$$

photon
absorption
(conductivity σ)

effective kinetic mixing
in a material

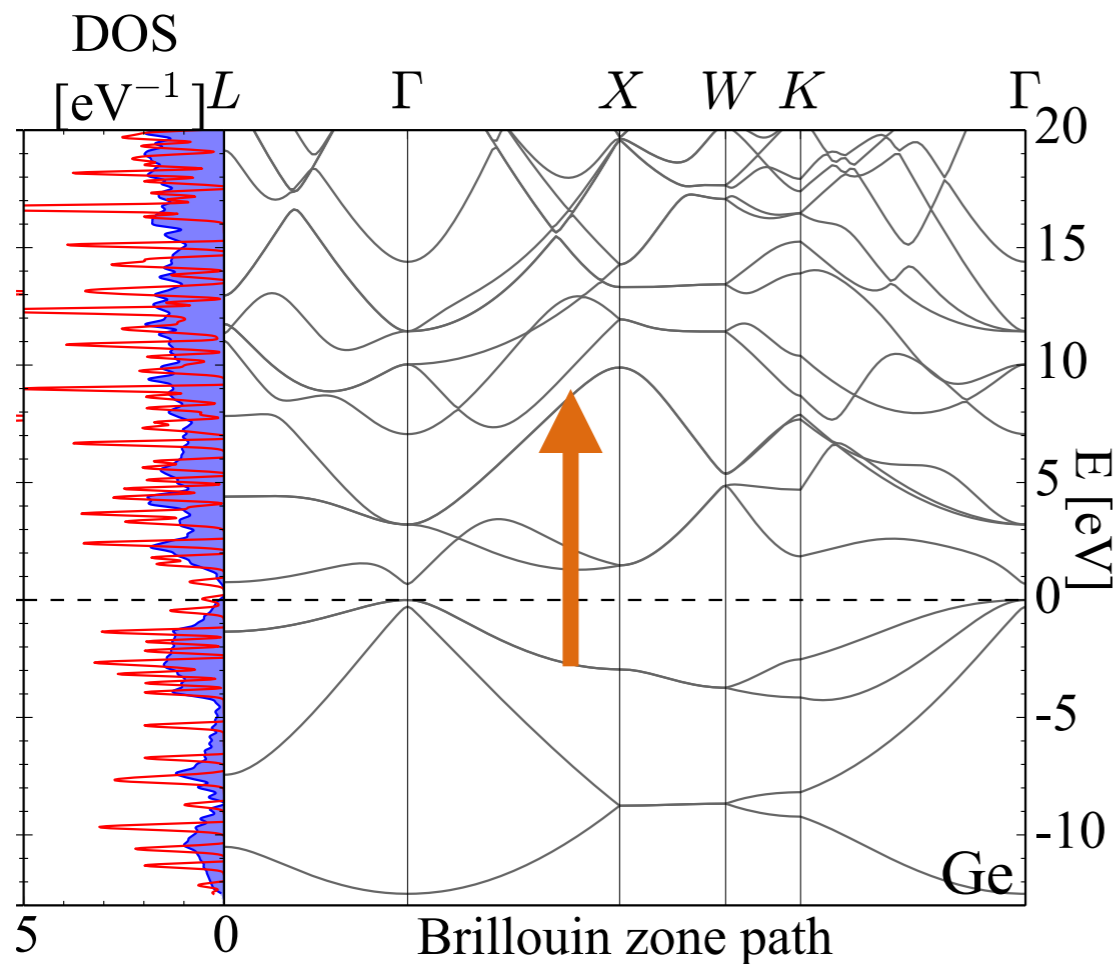
$$\kappa_{eff}^2 = \frac{\kappa^2 m_V^4}{[m_V^2 - \text{Re } \Pi(\omega)]^2 + [\text{Im } \Pi(\omega)]^2}$$

Semiconductor targets



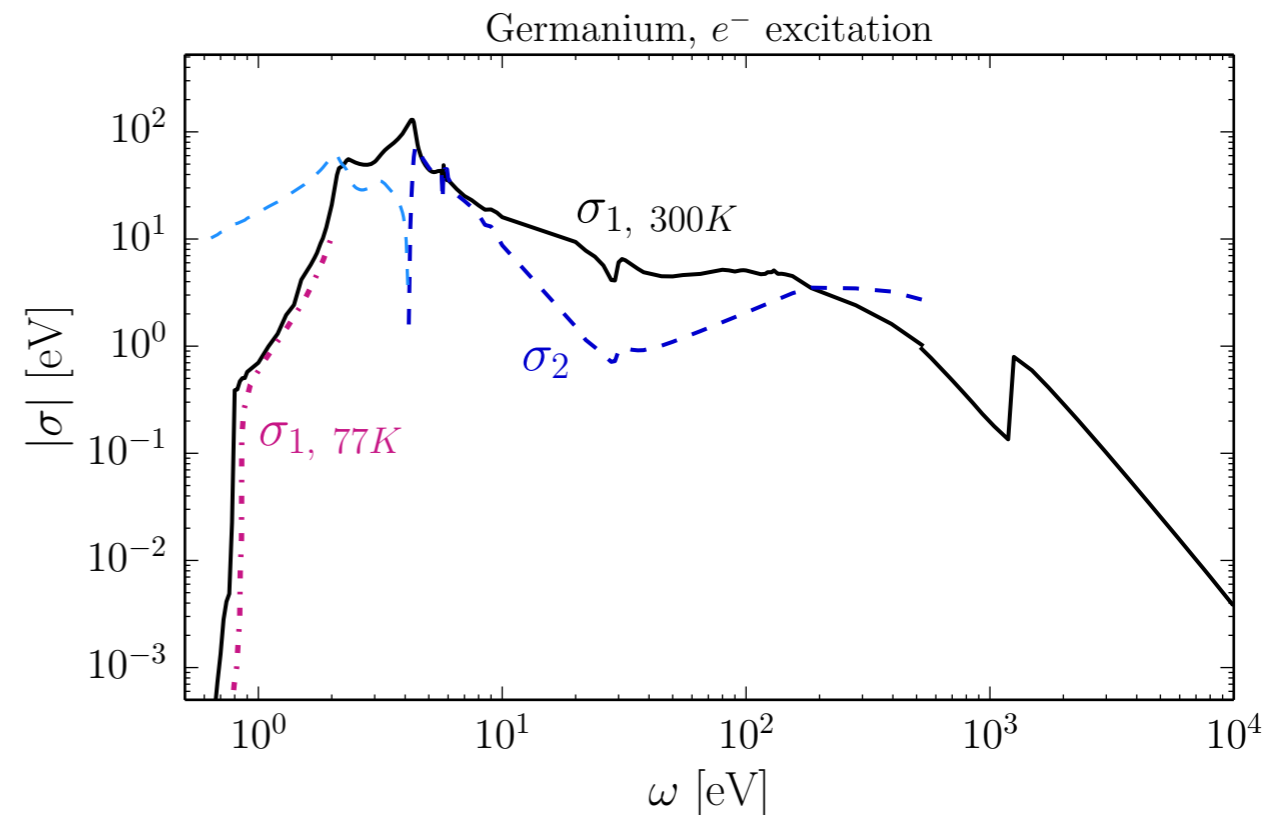
- Existing experiments already use Germanium, Silicon targets.
- Current $E_{th} \sim 50$ eV electron recoil.
- $E_{th} \sim \text{eV}$ could be reached in the near future.

Electron excitations



From Essig et al. 2015

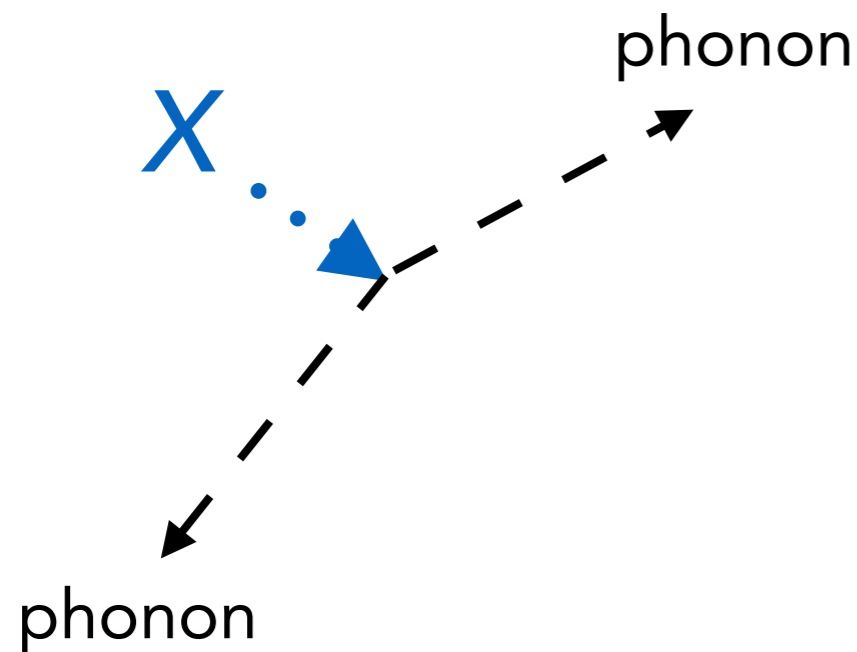
Photon absorption data



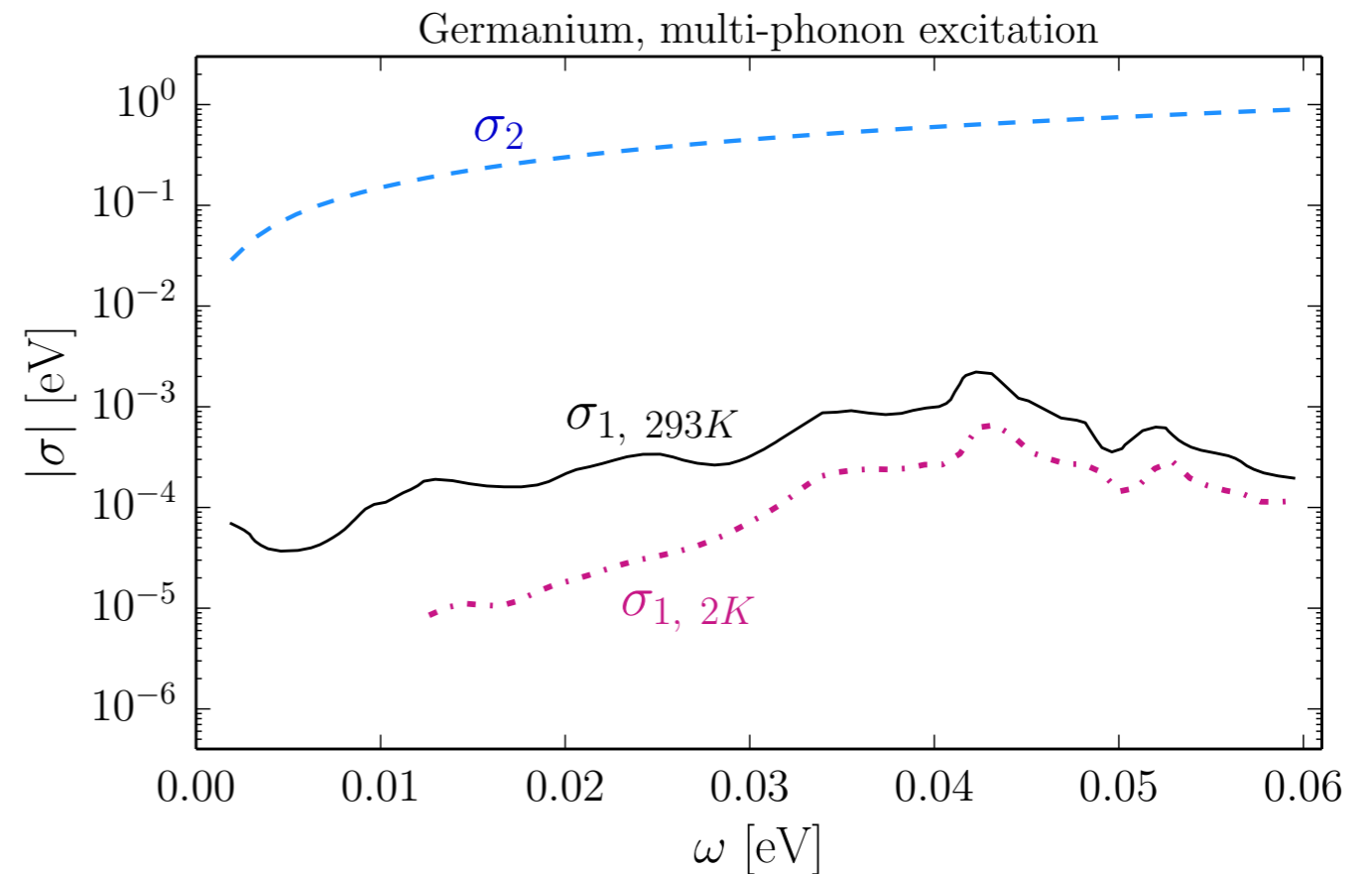
Optical absorption dominated by direct transitions

Band gap: 0.7 eV (germanium), 1.1 eV (silicon)

Multi-phonon excitation

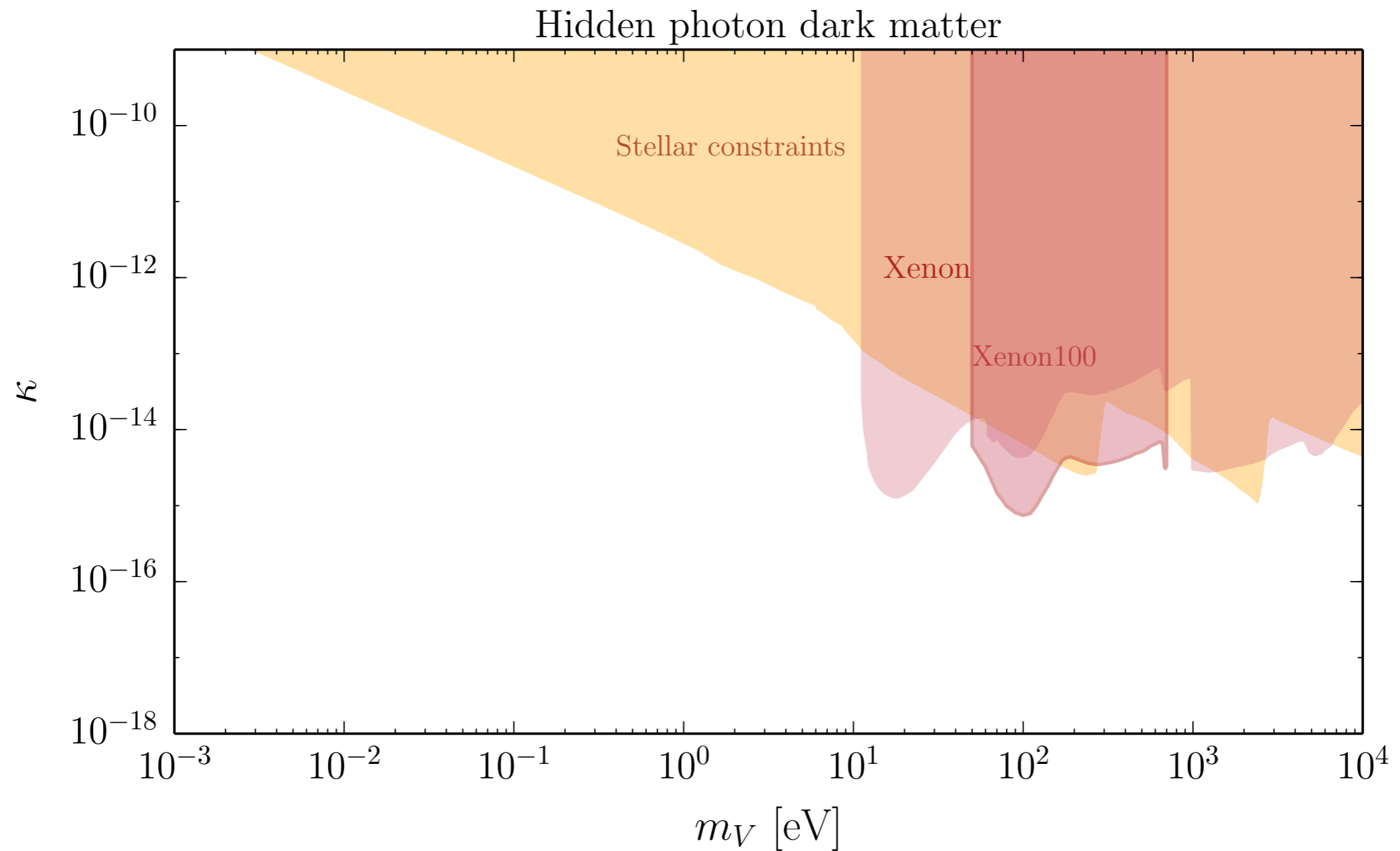


Photon absorption data



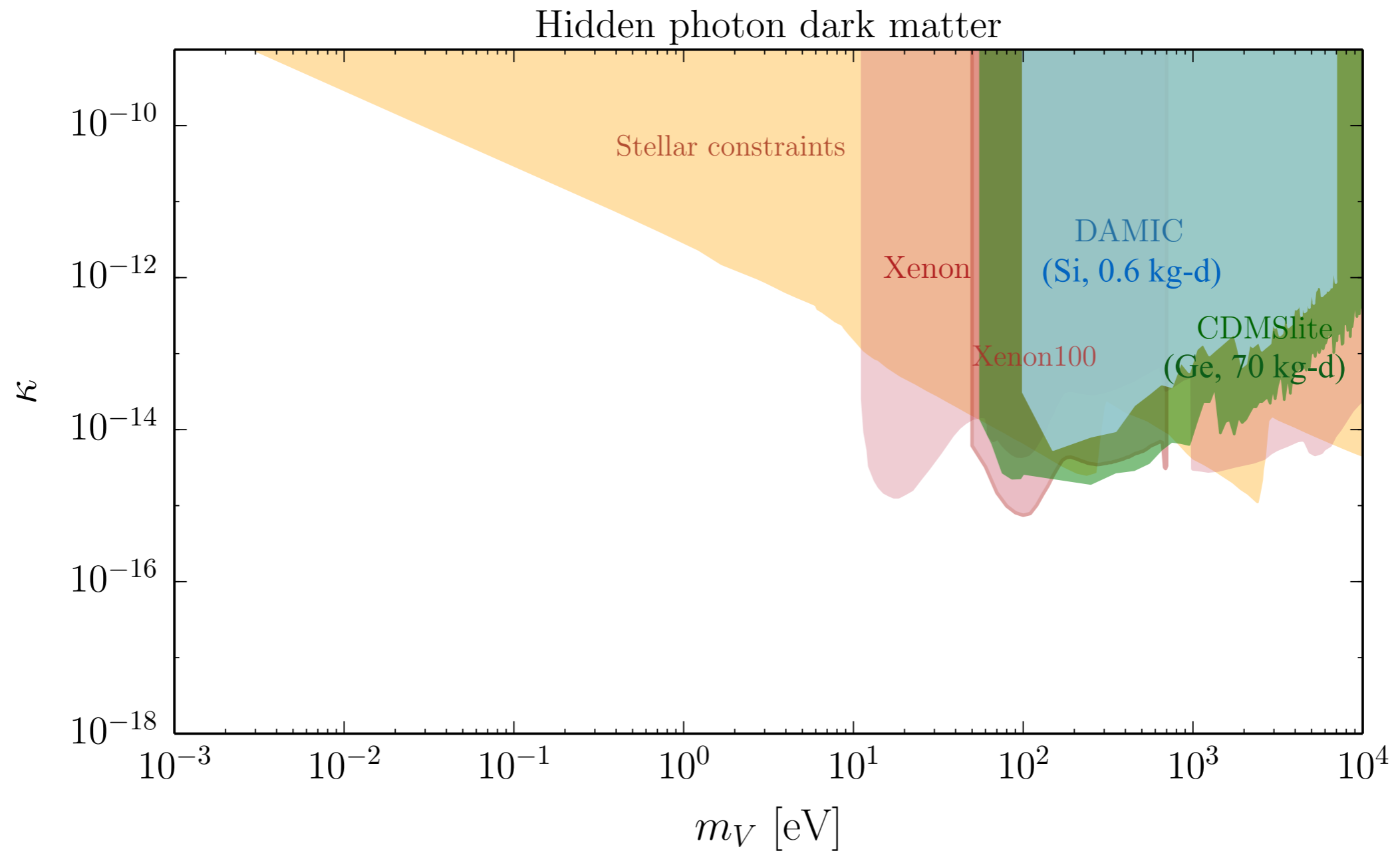
Optical absorption below the band gap is allowed if (multiple) phonons are excited instead!

Hidden photon DM



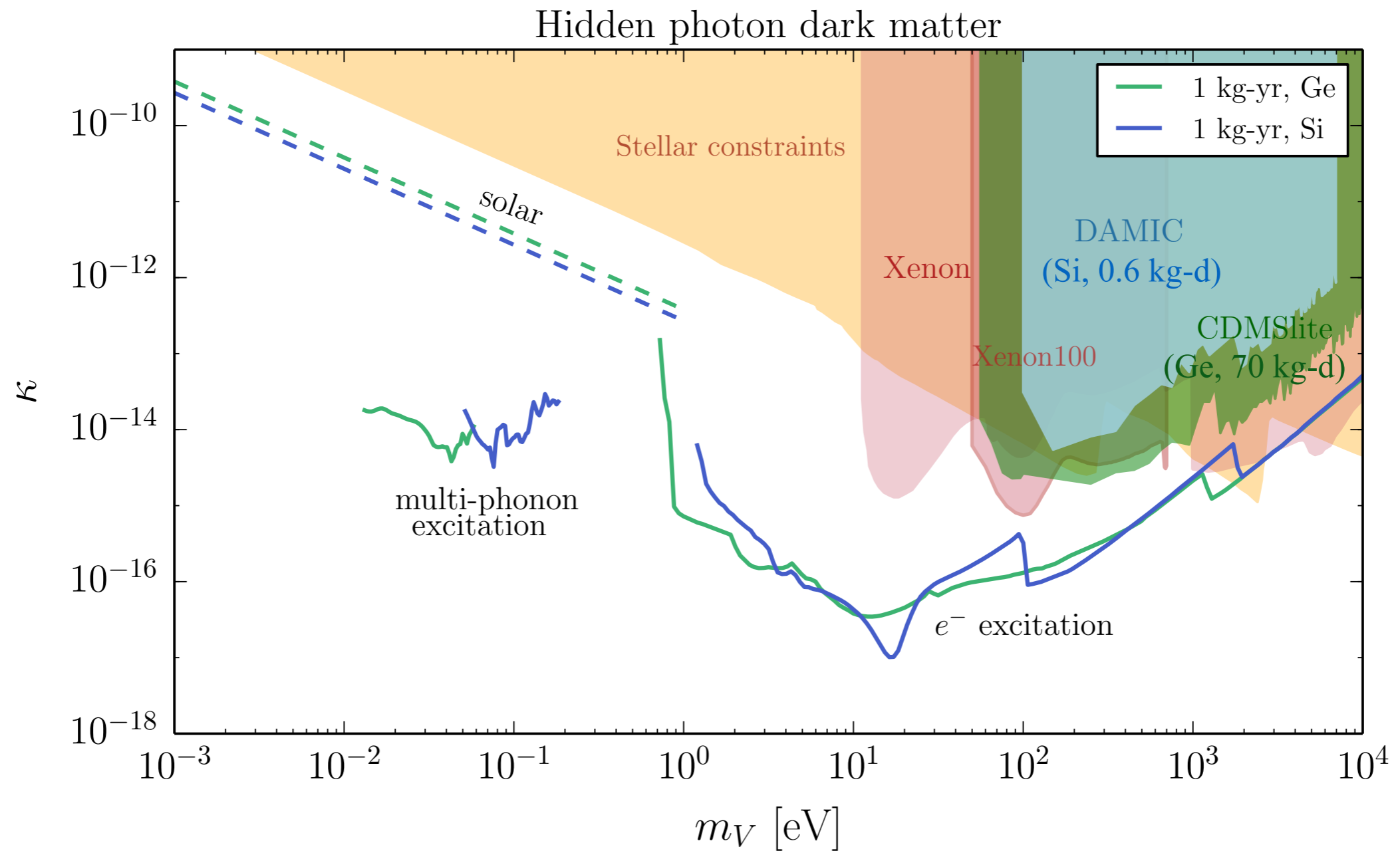
Stellar, Xenon10 constraints:
An, Pospelov, Pradler 2013, 2014
Redondo & Raffelt 2013

Hidden photon DM



Stellar, Xenon10 constraints:
An, Pospelov, Pradler 2013, 2014
Redondo & Raffelt 2013

Hidden photon DM



See also:

I. Bloch, Tien-Tien Yu, etc 2016
for similar study

Stellar, Xenon10 constraints:
An, Pospelov, Pradler 2013, 2014
Redondo & Raffelt 2013

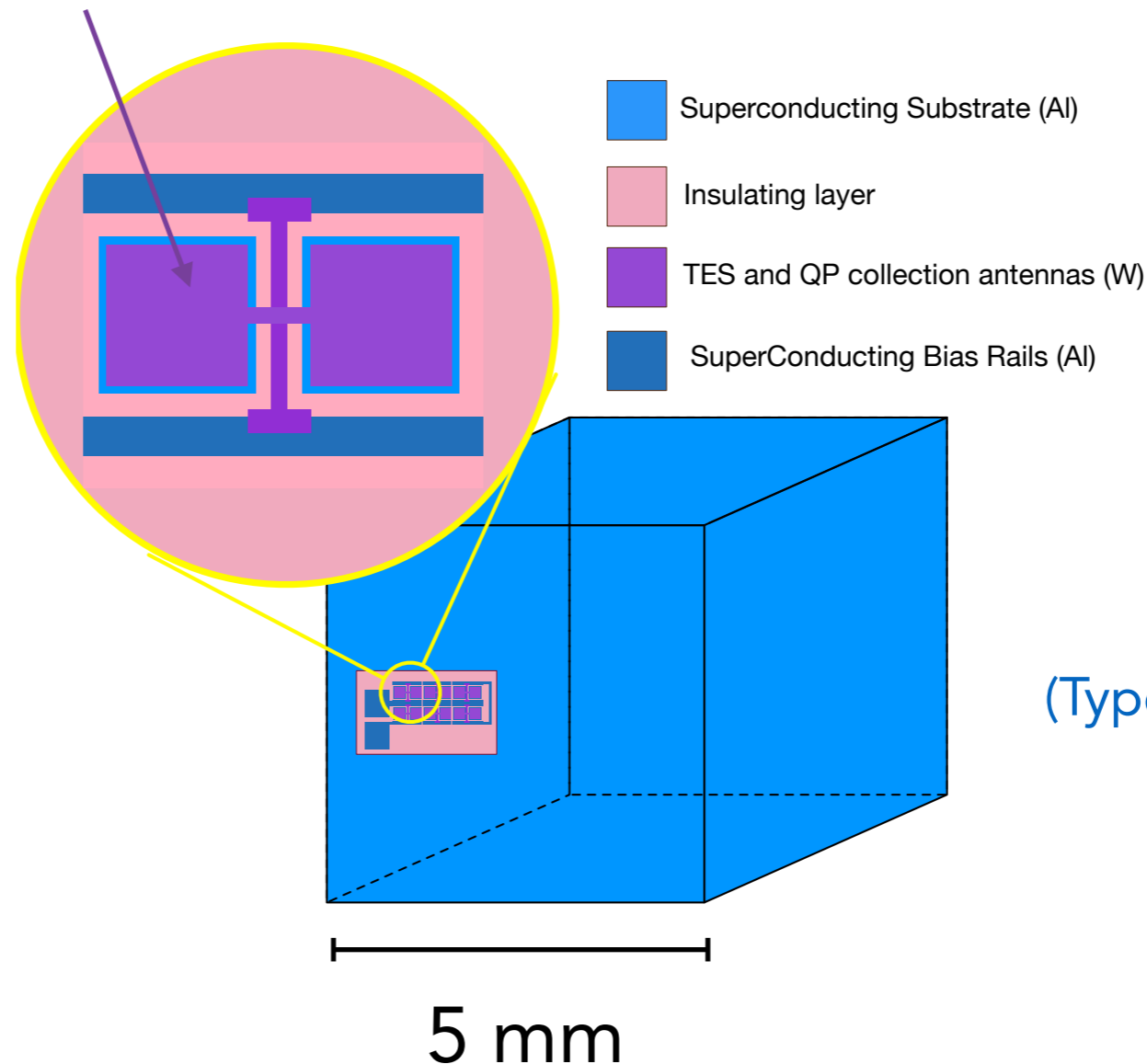
Superconductor target

Measurement
by sensitive
bolometer (TES)

$$T = 10 \text{ mK}$$

$E_{\text{th}} \sim \text{meV}$ goal

- Must improve substantially on current energy resolution $\sim 50\text{-}100 \text{ meV}$
- Noise is going to be an issue, backgrounds at low energies unknown.

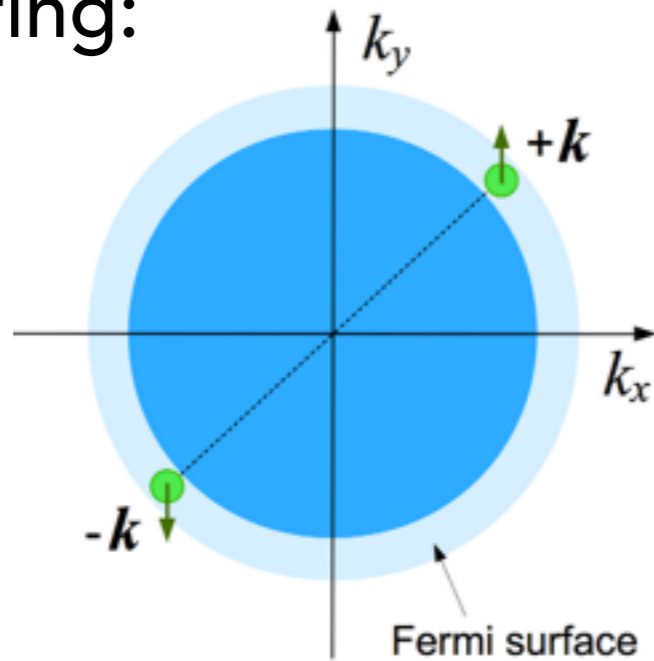


Aluminum
(Type I superconductor)

$$T_c = 1.2 \text{ K}$$

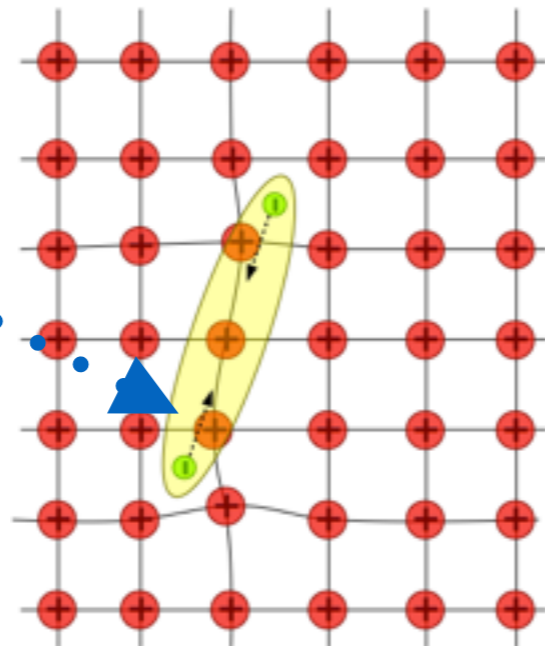
Why superconductors?

Cooper pairing:



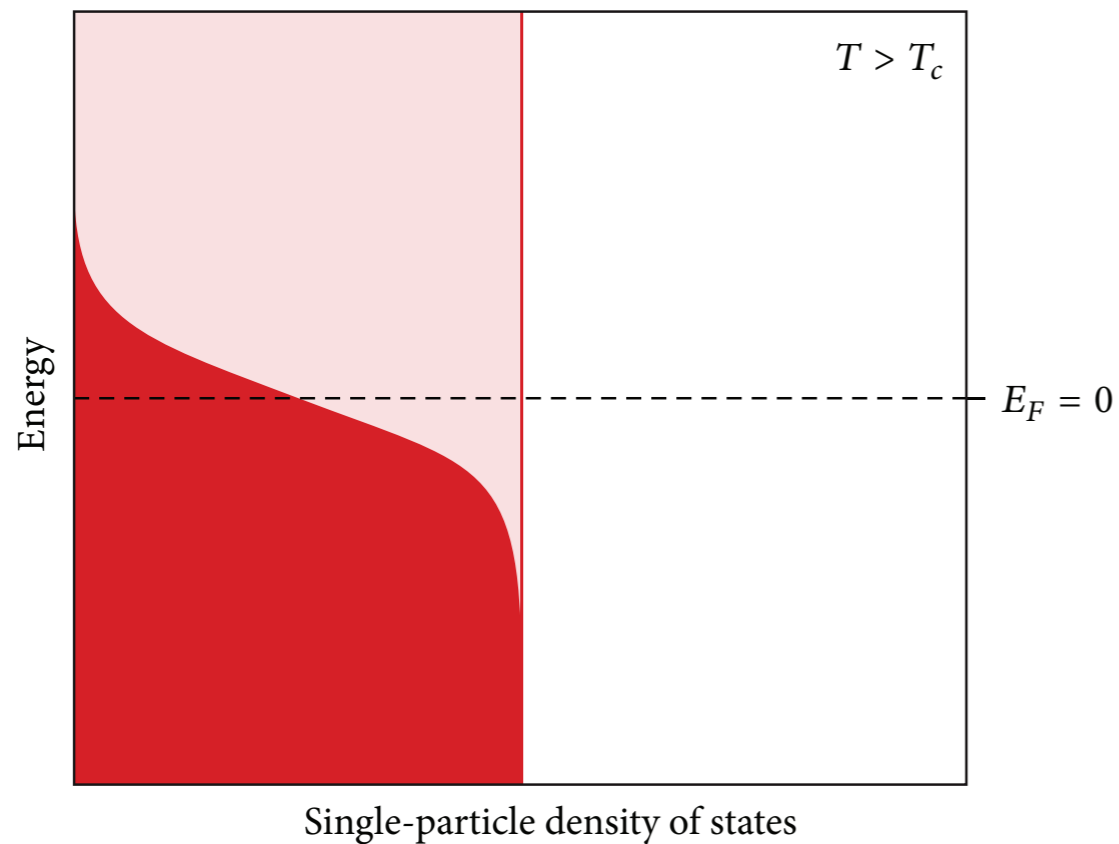
- Small band gap ($< \text{meV}$)
- Long-lived quasiparticle (electron) excitations
- Cooper pairs decoupled from thermal noise

$\Delta = \frac{3}{2}T_c \approx 0.3 \text{ meV}$

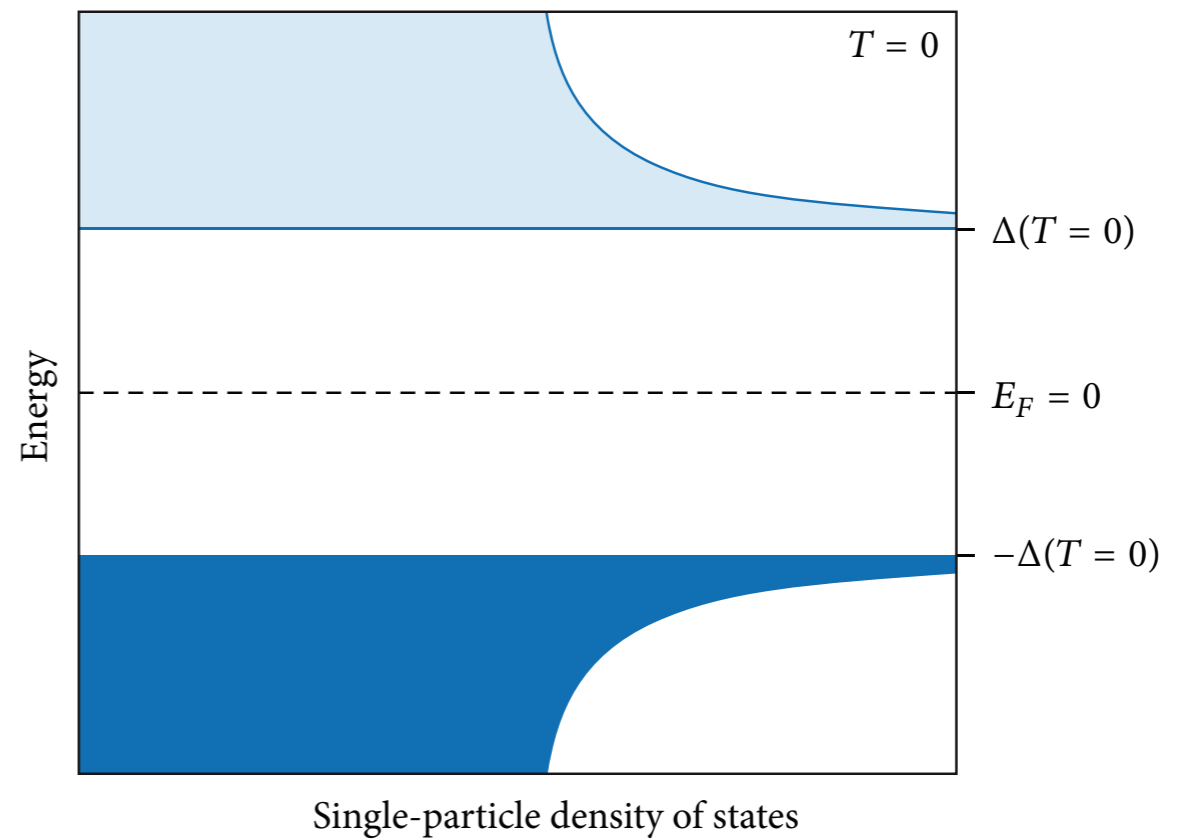


Superconductors

$$\Delta = \frac{3}{2}T_c \approx 0.3 \text{ meV} \quad \text{in aluminum}$$



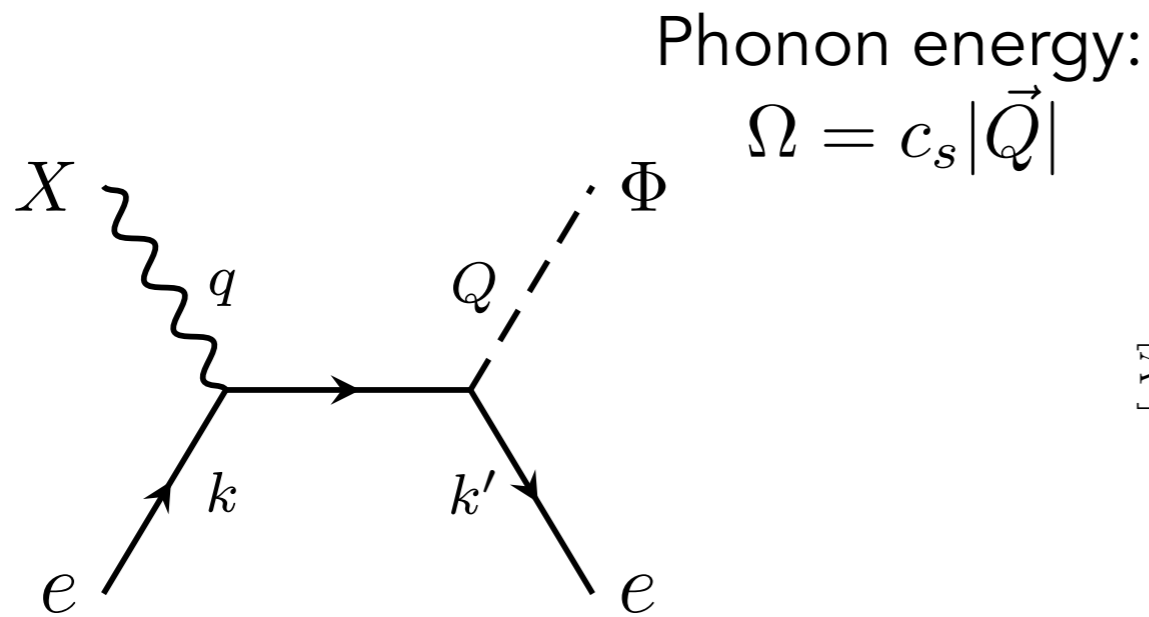
Metal



Superconductor

Because of gap, excitations (broken Cooper pairs) take a long time to recombine!

Absorption via phonon emission

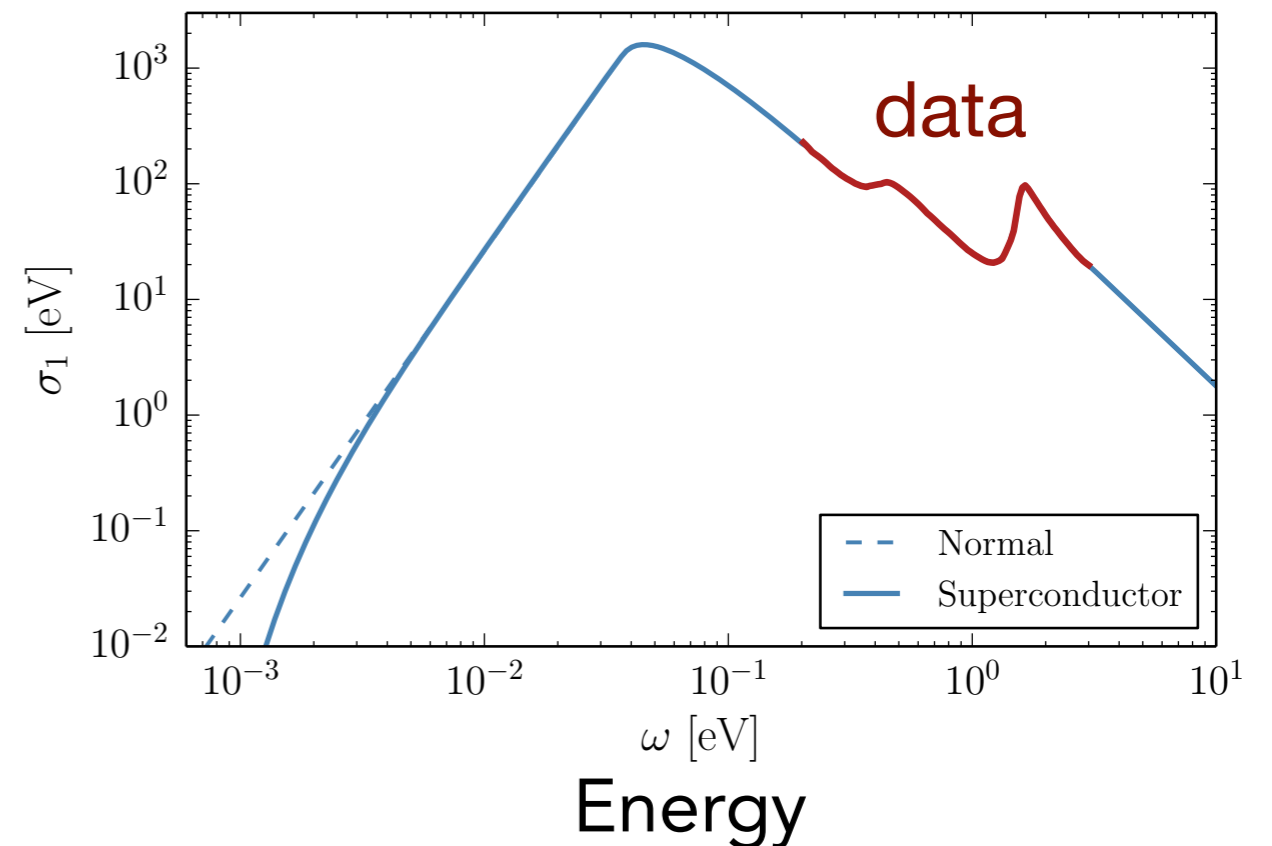


Speed of sound in aluminum:

$$c_s \simeq 6320 \text{ m/s} \sim 2 \times 10^{-5}$$

The phonon can carry large momentum, with small energy.

Photon absorption rate



Theory uncertainties:

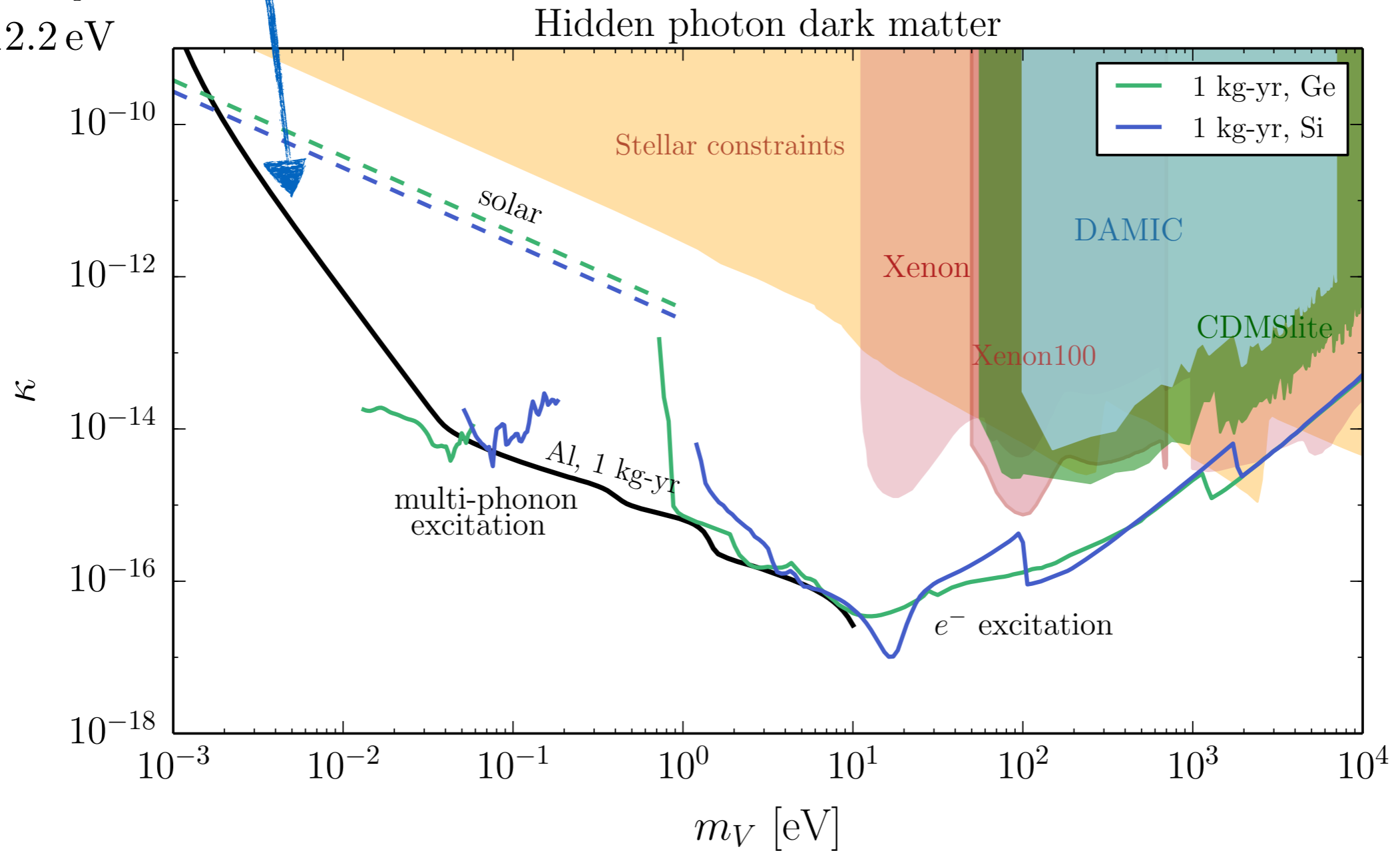
- phonon dispersion, coupling
- impurities

Suppression in absorption rate,
reduced effective kinetic mixing

Hidden photon DM

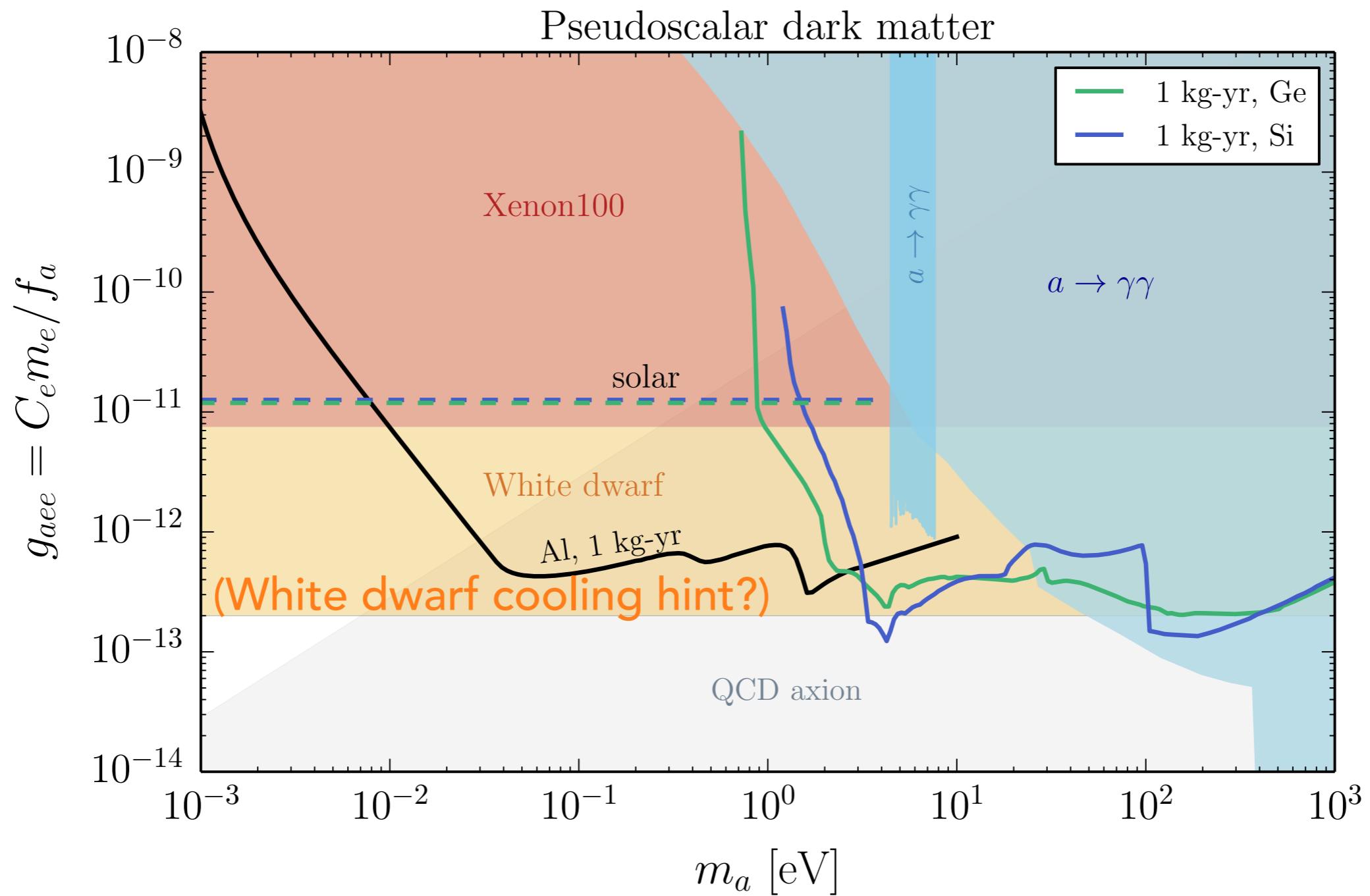
$$\kappa_{\text{eff}}^2 \simeq \frac{\kappa^2 m_V^4}{\omega_p^4}$$

$$\omega_p \approx 12.2 \text{ eV}$$

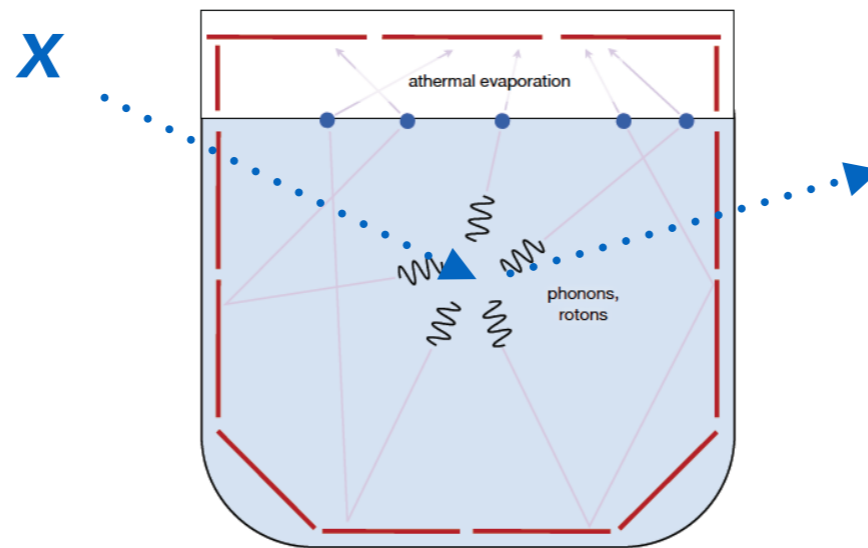


Stellar, Xenon10 constraints:
An, Pospelov, Pradler 2013, 2014
Redondo & Raffelt 2013

Pseudoscalar DM



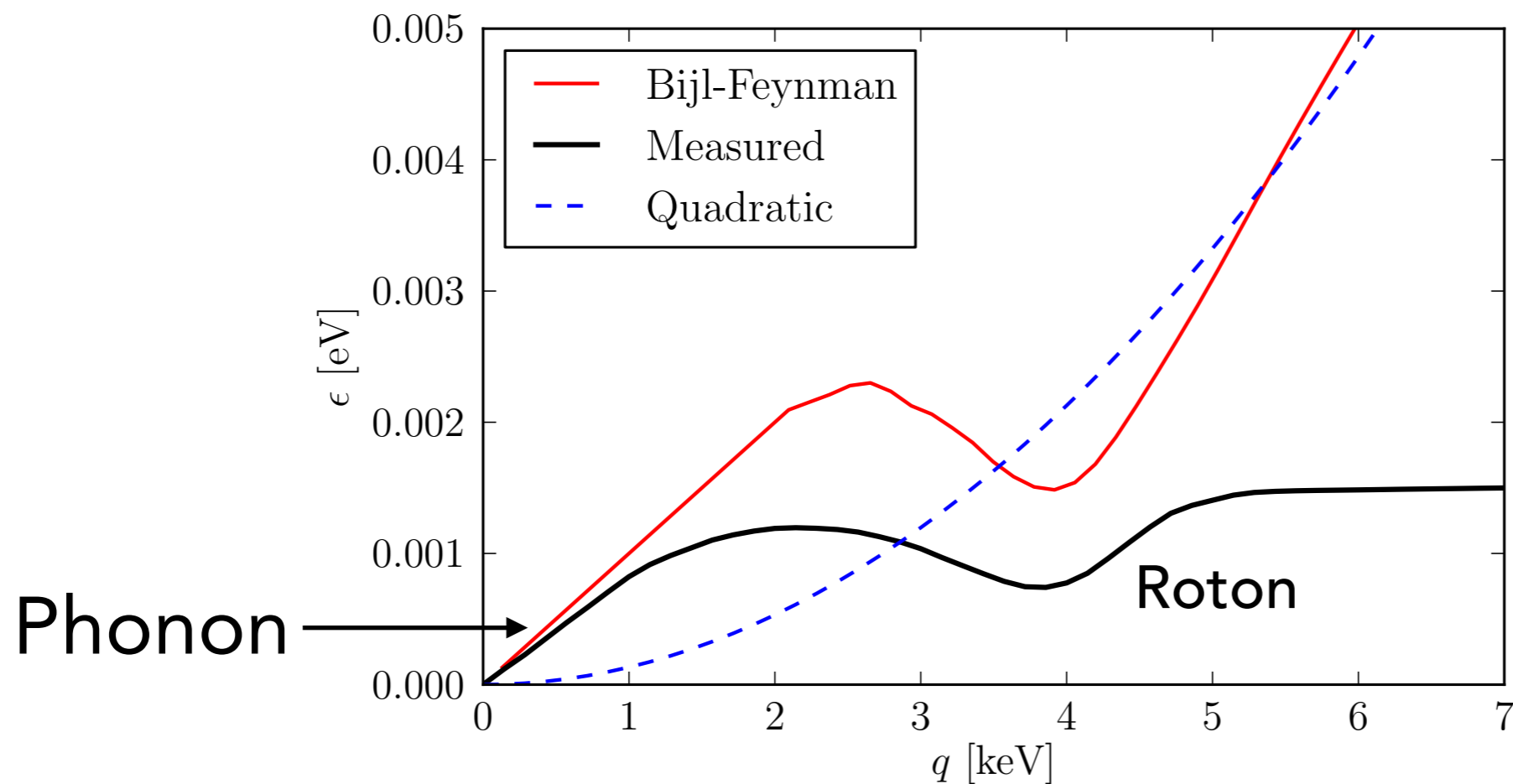
3b. DM scattering in superfluid He



in progress, with:
Simon Knapen
Eckhard Krotscheck
Kathryn Zurek

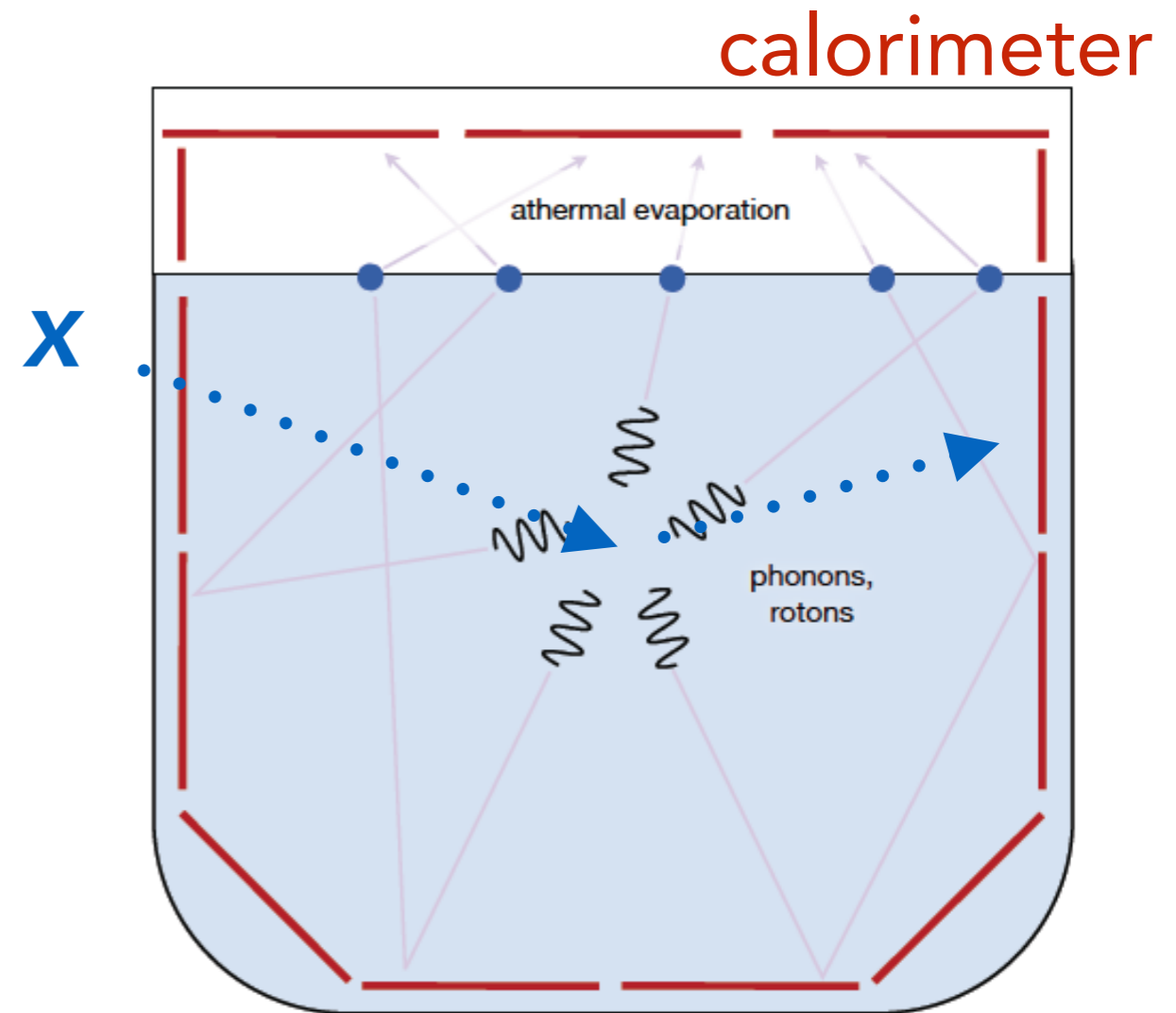
Why superfluid Helium?

- Long-lived quasiparticle excitations
- Potential detection of meV energy depositions



Detector concept

- Operation at 10-100 mK
- Evaporation and amplification
x10-40 down to 0.6 meV!
- Measure scintillation at
energies ≥ 100 eV
- Superfluid helium used for
ultracold neutrons, etc.
- Easy to obtain and purify



Scattering in Helium

- Neutron scattering:

$$\frac{d^2\sigma}{d\omega d\Omega} \equiv \frac{\sigma^n}{4\pi} \frac{k_f}{k_i} S(Q, \omega)$$

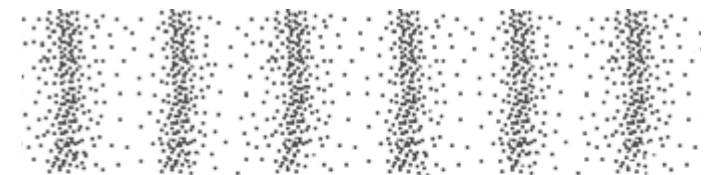
Dynamic structure factor:

$$S(\mathbf{Q}, \omega) = \frac{1}{n_0} \sum_F |\langle F | n_{\mathbf{Q}} | I \rangle|^2 \delta(E_F - E_I - \omega)$$

response of the fluid to a **density** perturbation with momentum transfer \mathbf{Q} , energy ω

Density operator:

$$\hat{n}_{\mathbf{k}} = \frac{1}{\sqrt{N}} \sum_i e^{i\mathbf{k} \cdot \mathbf{r}_i}$$



- Apply directly to dark matter scattering

$$\frac{d^2\sigma}{d\omega d\Omega} \equiv \frac{\sigma^n}{4\pi} \frac{k_f}{k_i} S(Q, \omega)$$

Multiphonon excitations

Dynamic structure factor $S(Q, \omega)$

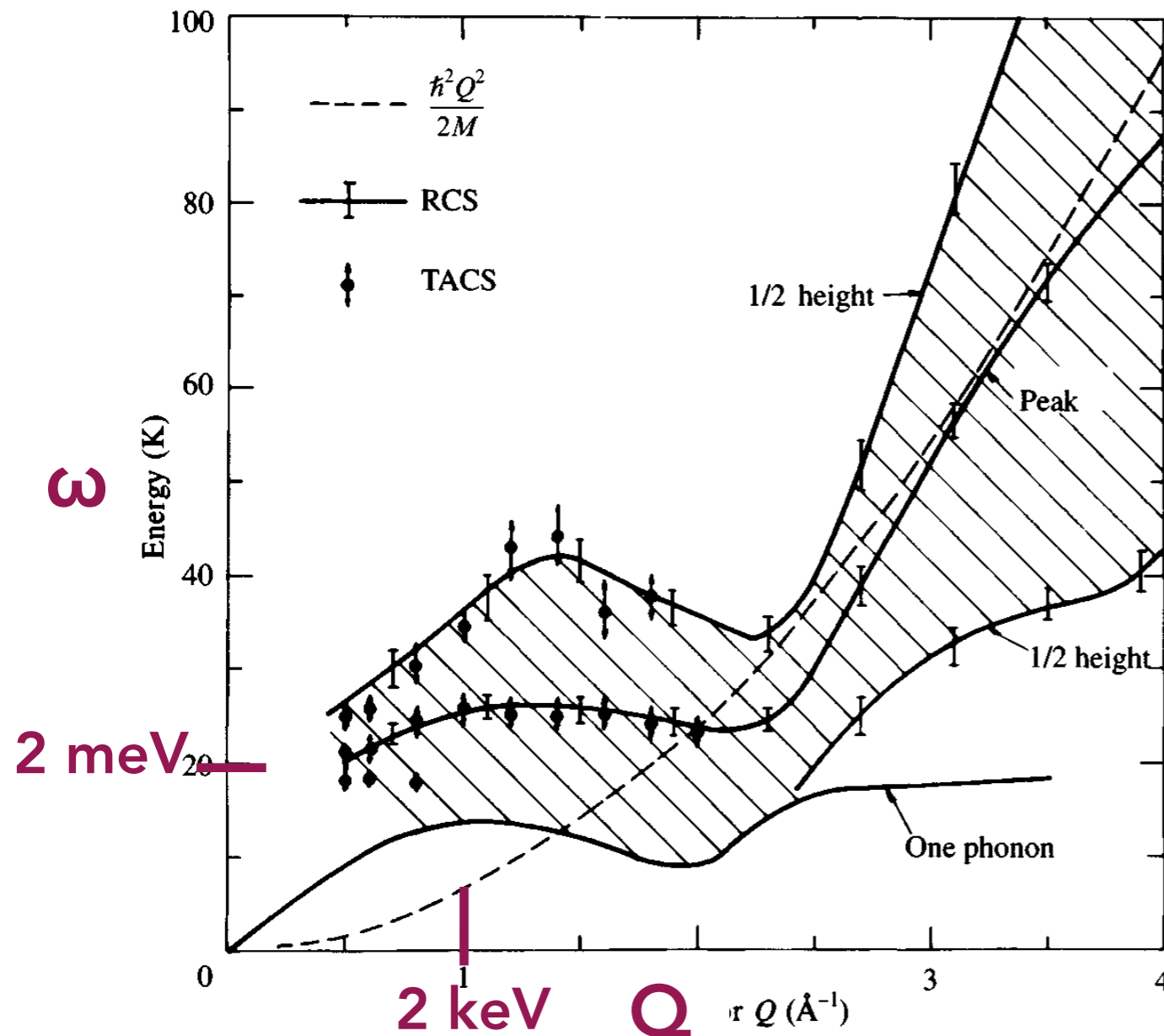
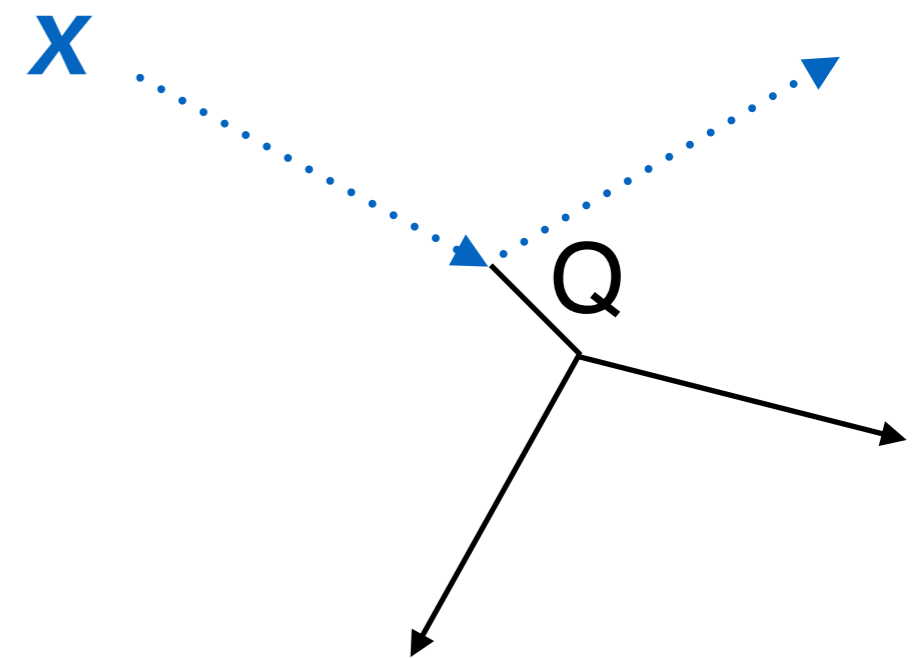


Fig. 1.6. The broad high-energyparticle (or multiphonon) structure in $S(Q, \omega)$ at intermediate Q at 1.1 K and SVP. The “peak” position and its width are indicated [Source: Cowley and Woods, 1971].

2-phonon excitation



Kinematics allows larger energy ω deposited, even for small Q

Theory vs. experiment

Dynamic structure factor $S(Q, \omega)$

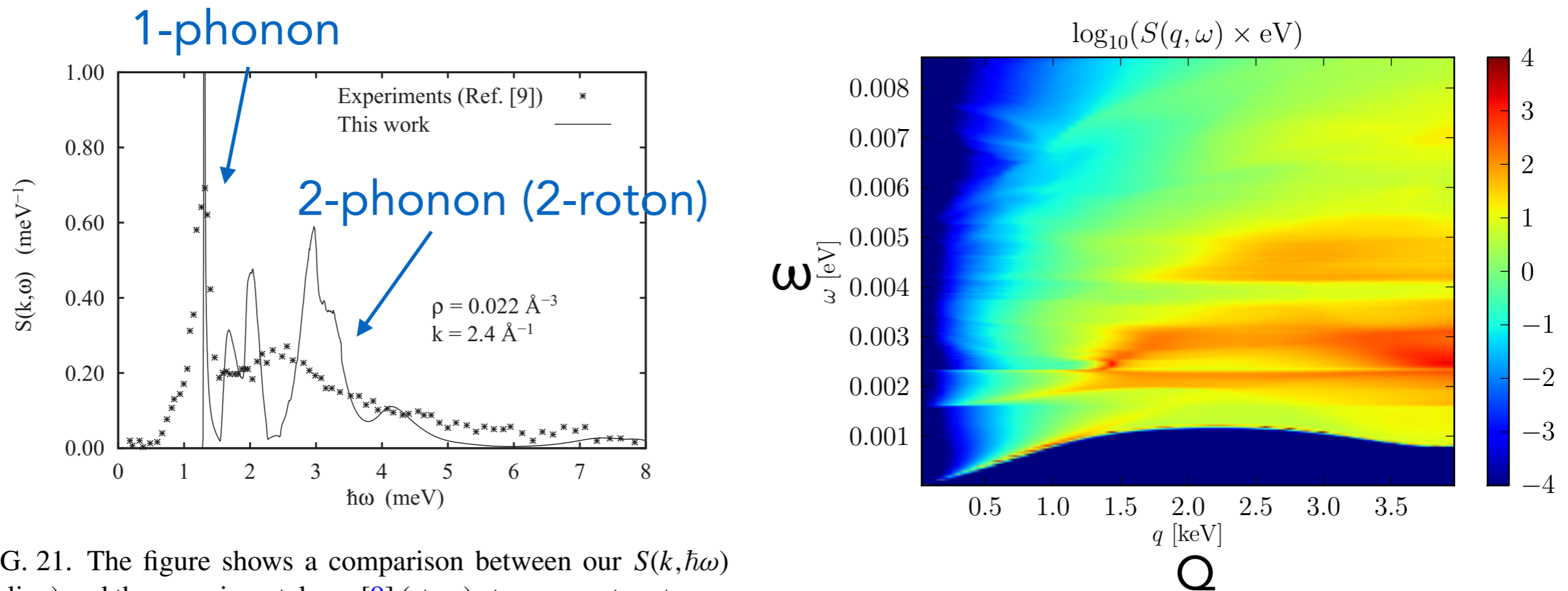
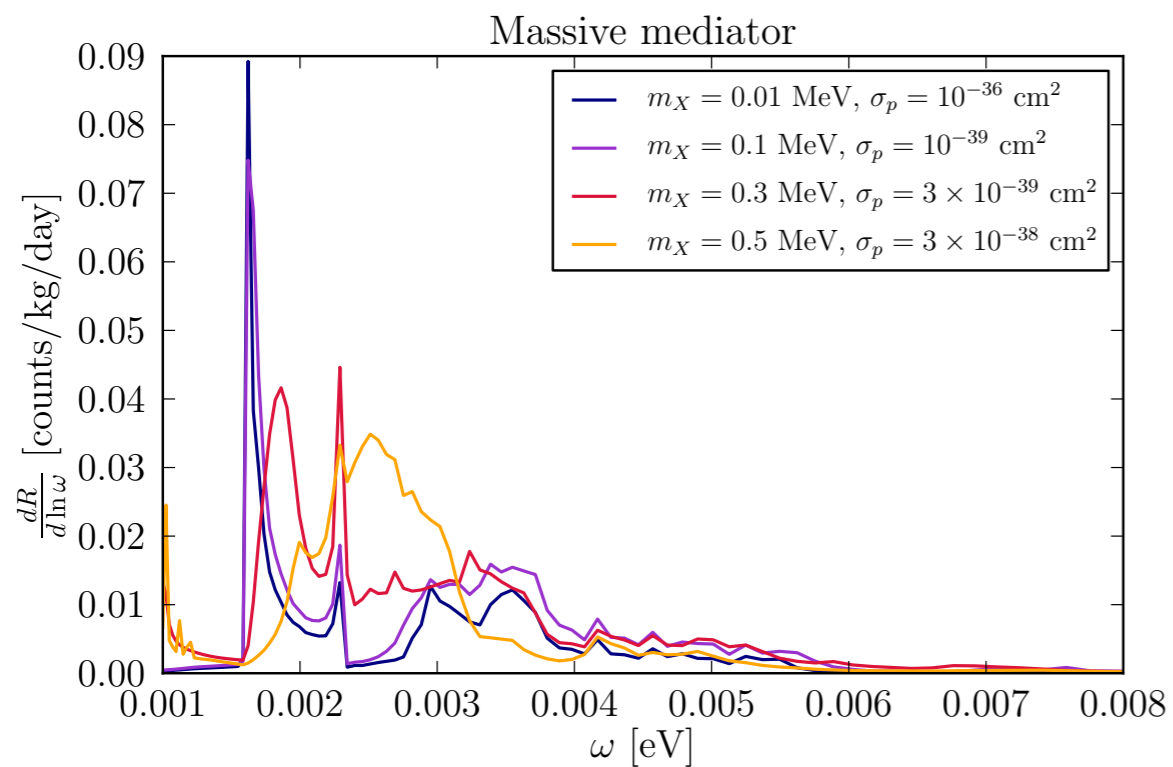


FIG. 21. The figure shows a comparison between our $S(k, \hbar\omega)$ (solid line) and the experimental one [9] (stars) at a momentum transfer of $k = 2.4 \text{ \AA}^{-1}$. The experimental values have been normalized to satisfy the m_0 sum rule.

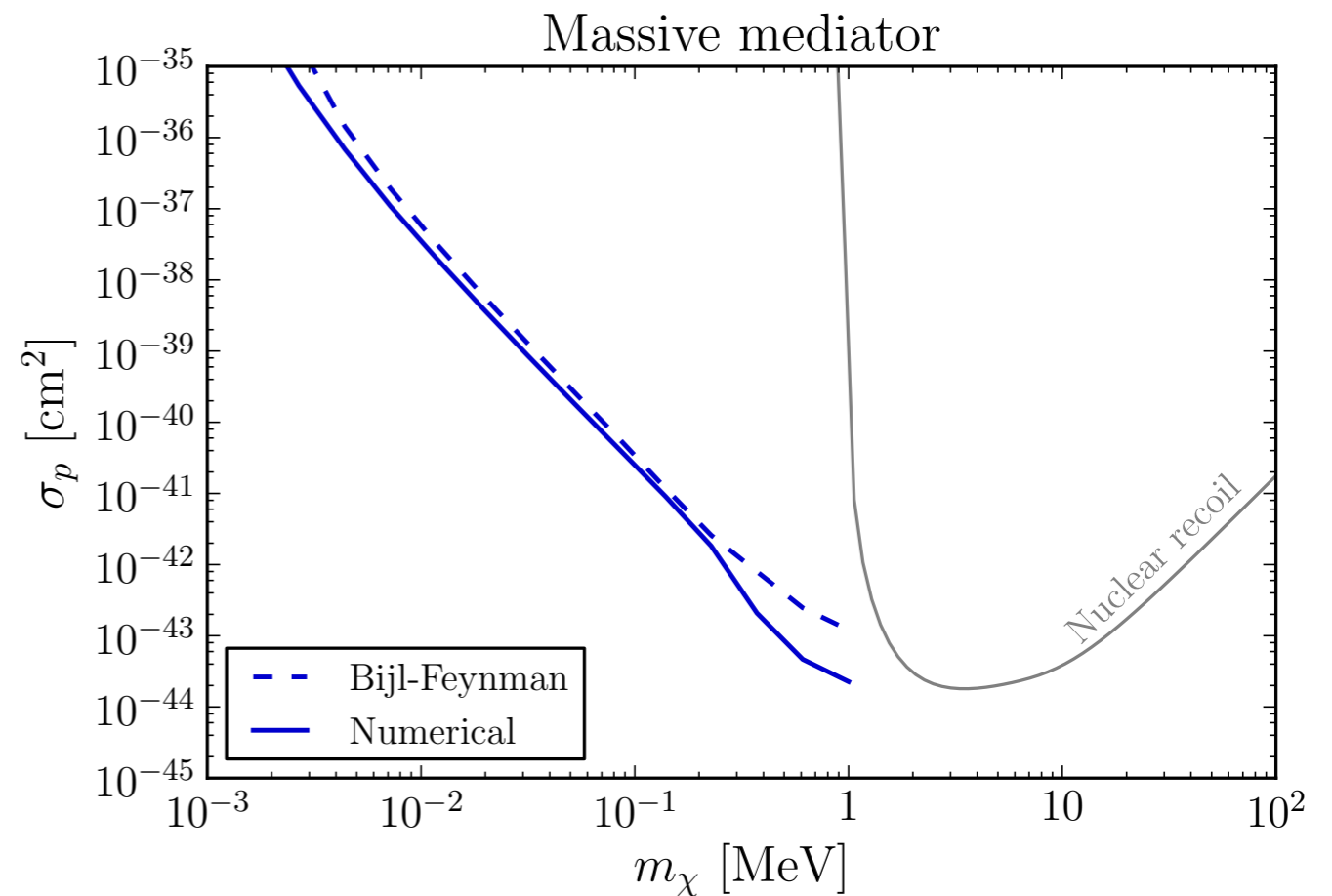
Theory

Sensitivity to light DM

Two-phonon scattering rate



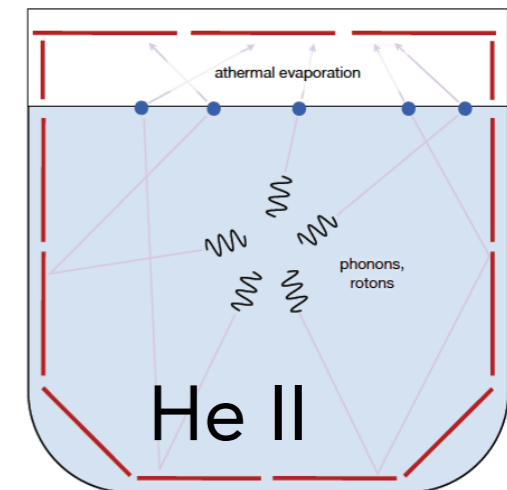
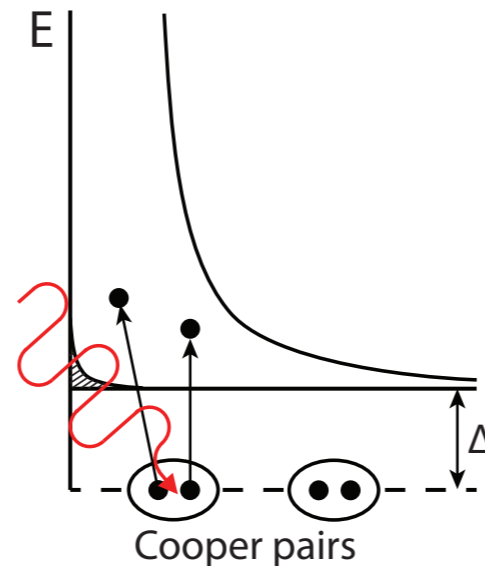
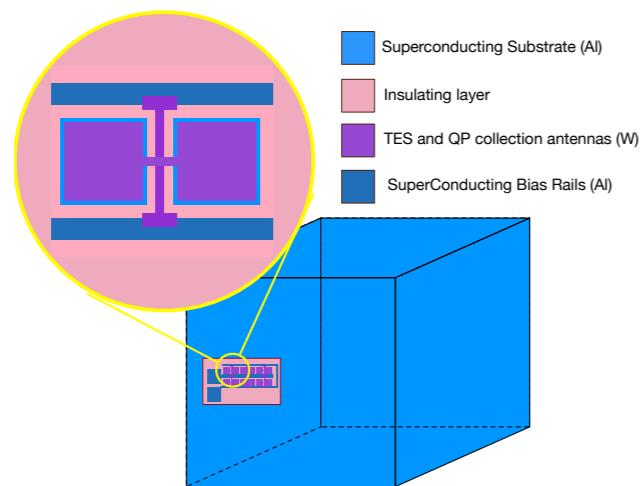
Dominant scattering into two rotons



Extends reach to lower mass DM!
(~10 meV resolution)

Conclusions

- Current direct detection experiments focus on WIMP dark matter at or above 1 GeV
- New ideas (for models and detection) explore many orders of magnitude in mass!



Absorption of bosonic DM down to meV

Scattering of $< \text{MeV}$ DM