

Partially Acoustic Dark Matter, Interacting Dark Radiation, & Large Scale Structure

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in collaboration with

Zackaria Chacko, Yanou Cui, Sungwoo Hong, and Takemichi Okui

KICP Seminar, Oct 14 2016

Explaining (σ_8, H_0) Discrepancies using Non-Minimal Dark Sector

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Non-minimal Dark Sector and LSS

I'll focus on the solution of LSS "puzzles" through DM-DR scattering

Z. Chacko, Y. Cui, S. Hong, T. Okui, YT (2016)

M. Buen-Abad, G. Marques-Tavares, and M. Schmaltz (2015)

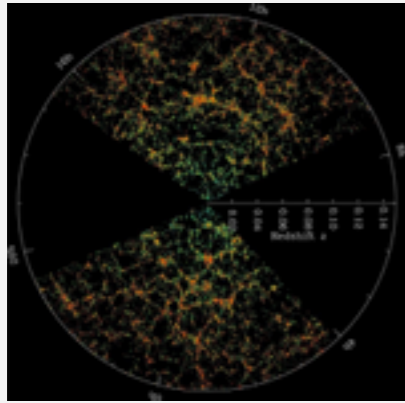
J. Lesgourgues, G. Marques-Tavares, and M. Schmaltz (2015)

P. Ko and Y. Tang (2016)

Other proposal

Decaying DM, Massive neutrinos, Dark energy models, ...

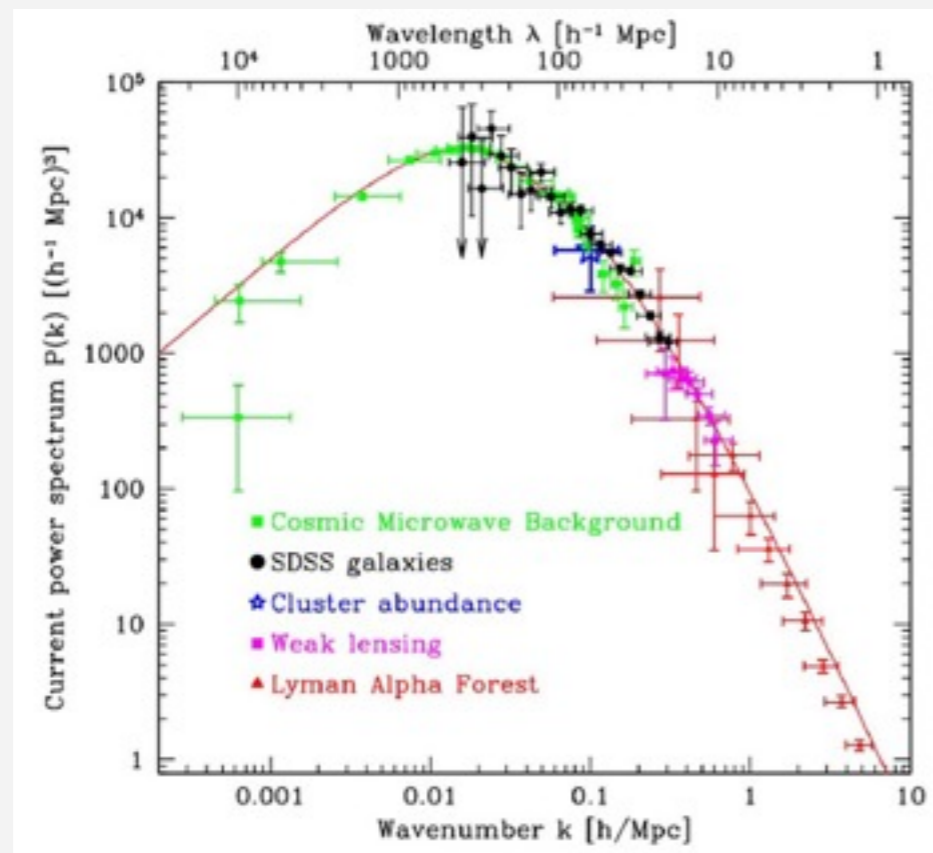
Precision measurement in cosmology



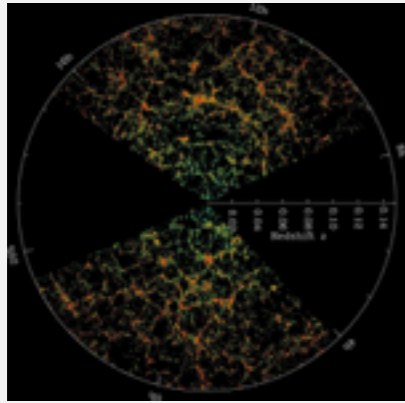
Matter Power Spectrum

Three ways to measure the spectrum

2004



Precision measurement in cosmology

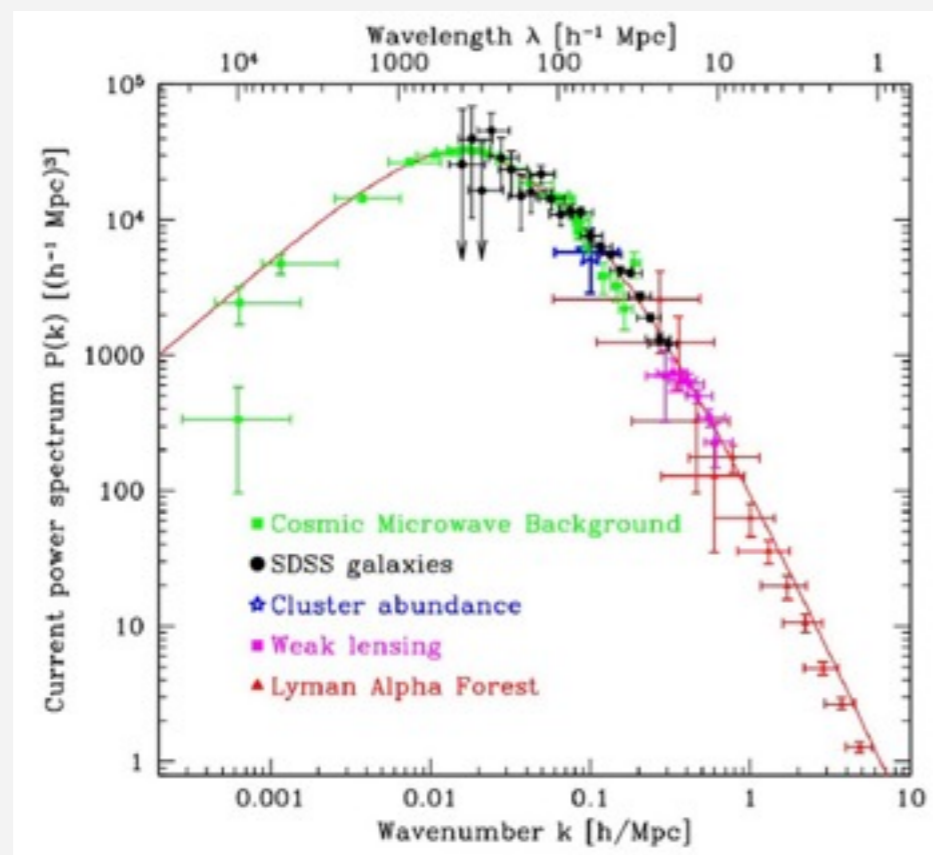


Matter Power Spectrum

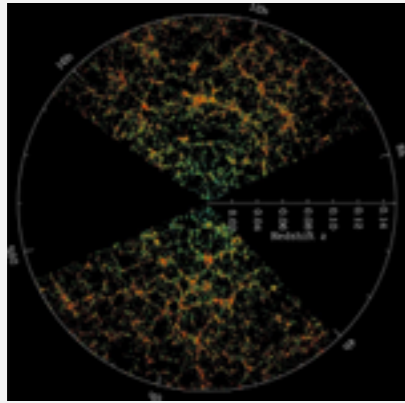
2004

Three ways to measure the spectrum

Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today



Precision measurement in cosmology



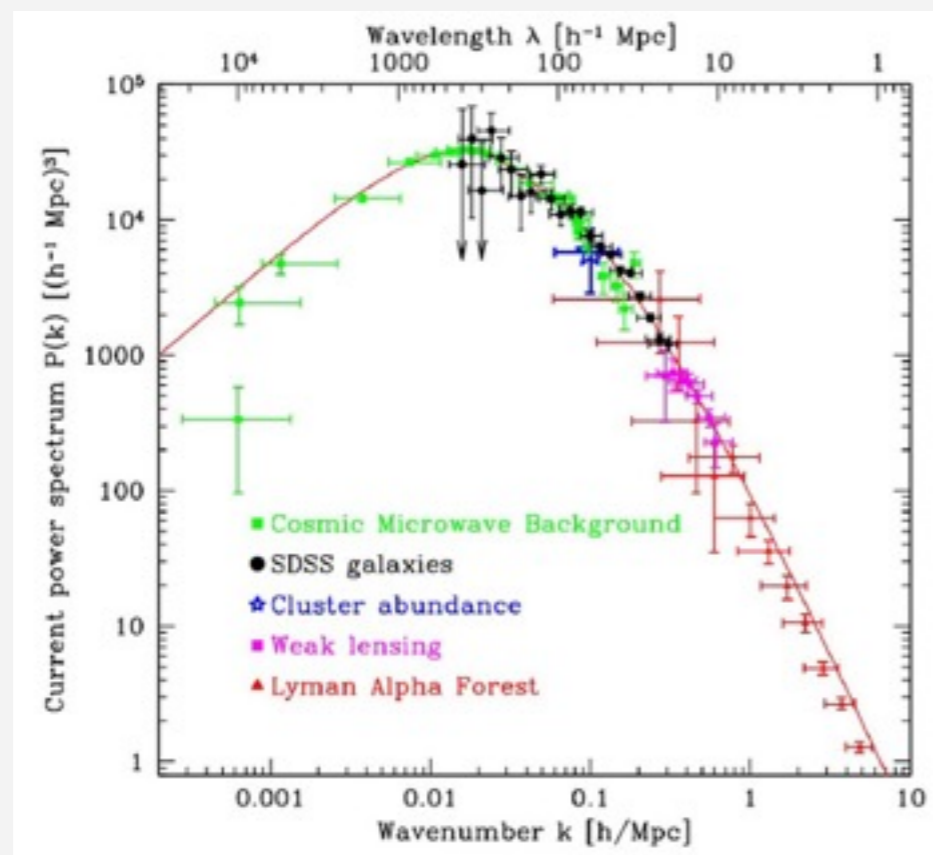
Matter Power Spectrum

2004

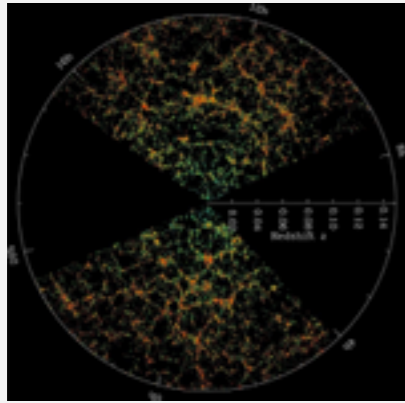
Three ways to measure the spectrum

Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today

Map the galaxy distribution, then fit the DM distribution



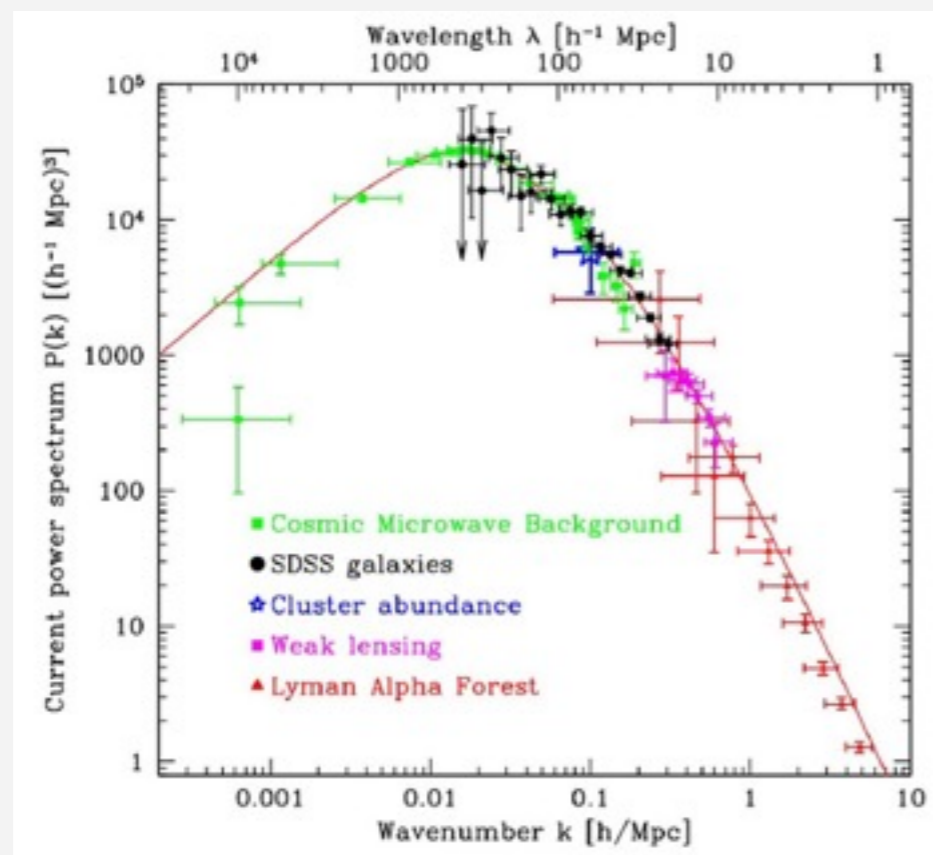
Precision measurement in cosmology



Matter Power Spectrum

2004

Three ways to measure the spectrum

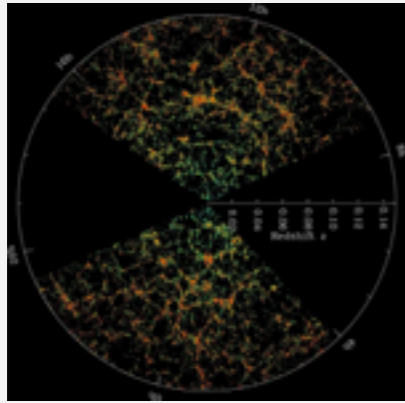


Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today

Map the galaxy distribution, then fit the DM distribution

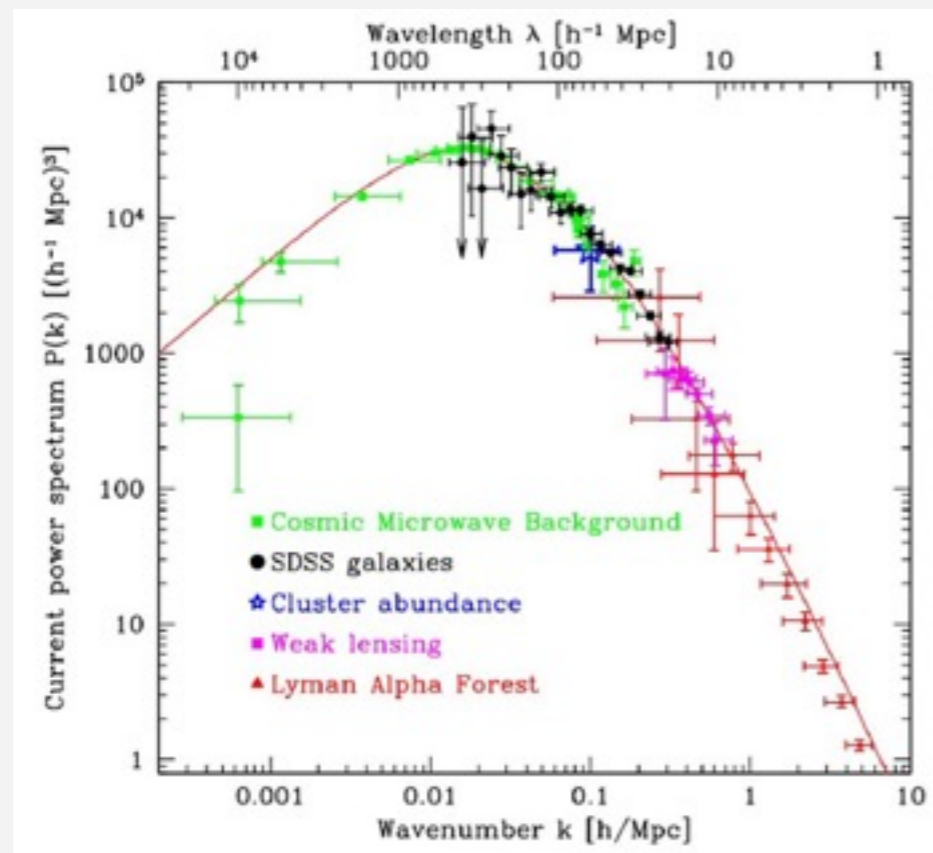
Map the DM distribution directly using weak lensing experiments

Precision measurement in cosmology

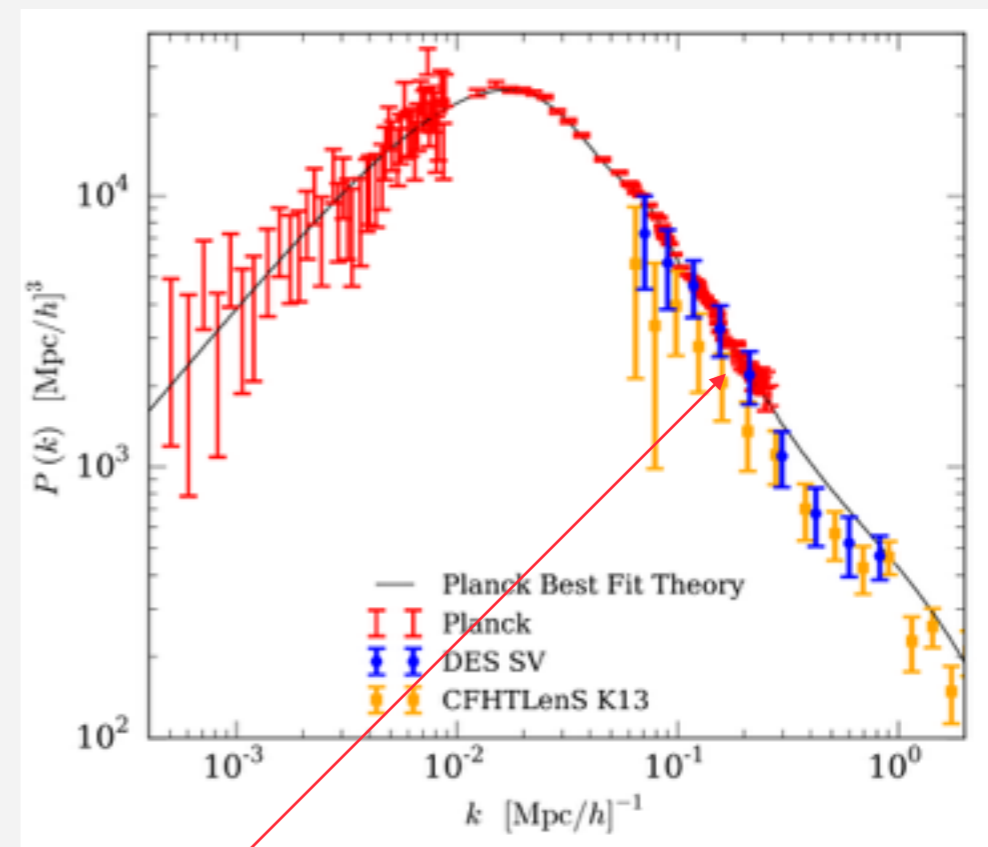


Matter Power Spectrum

2004

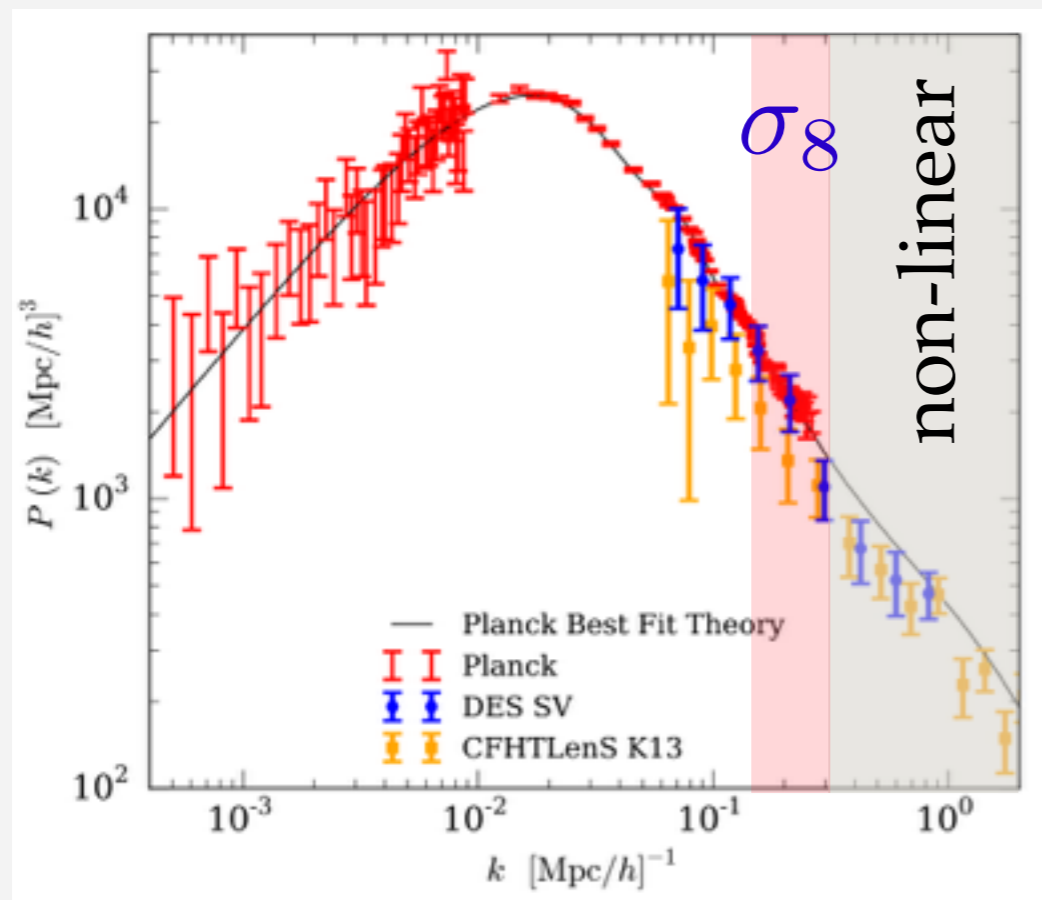


DES: 1507.05552



Can start to cross check the CMB & WL results

The σ_8 problem



DES: 1507.05552

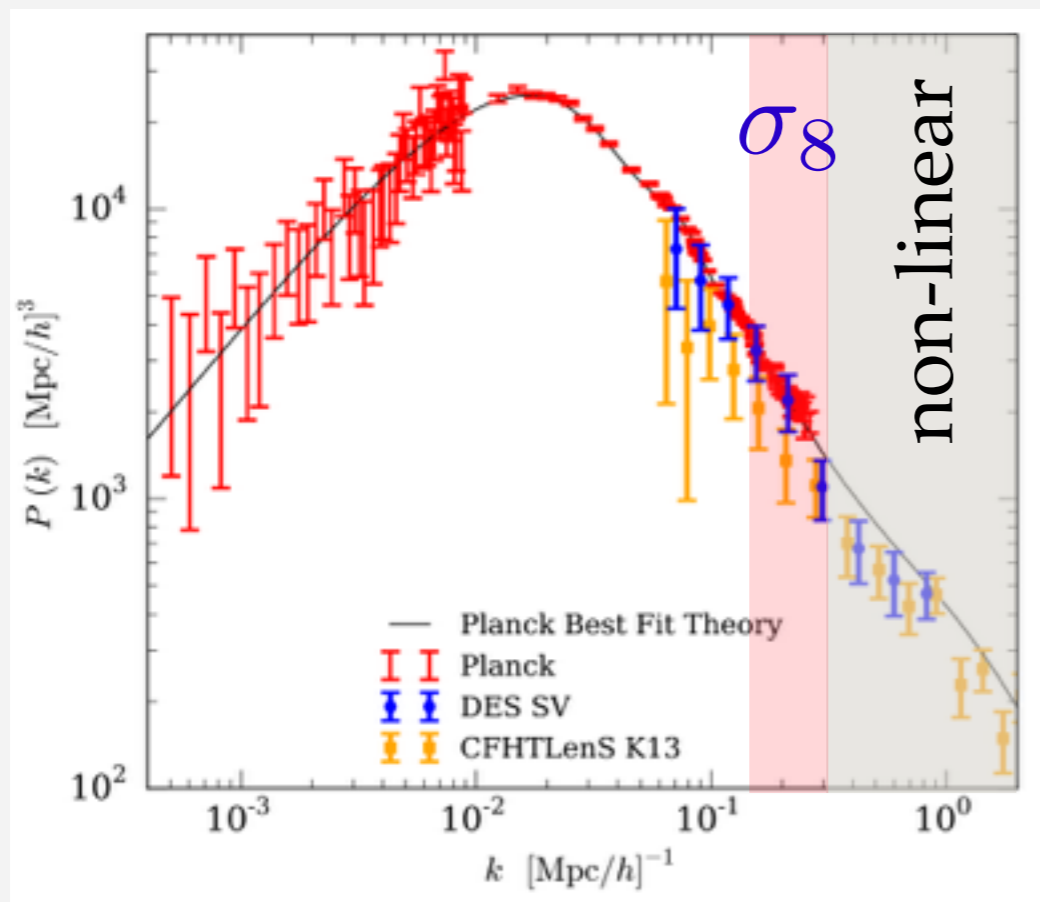
σ_8

~ amplitude of matter fluctuation
on the scale of $8 h^{-1} \text{Mpc}$.

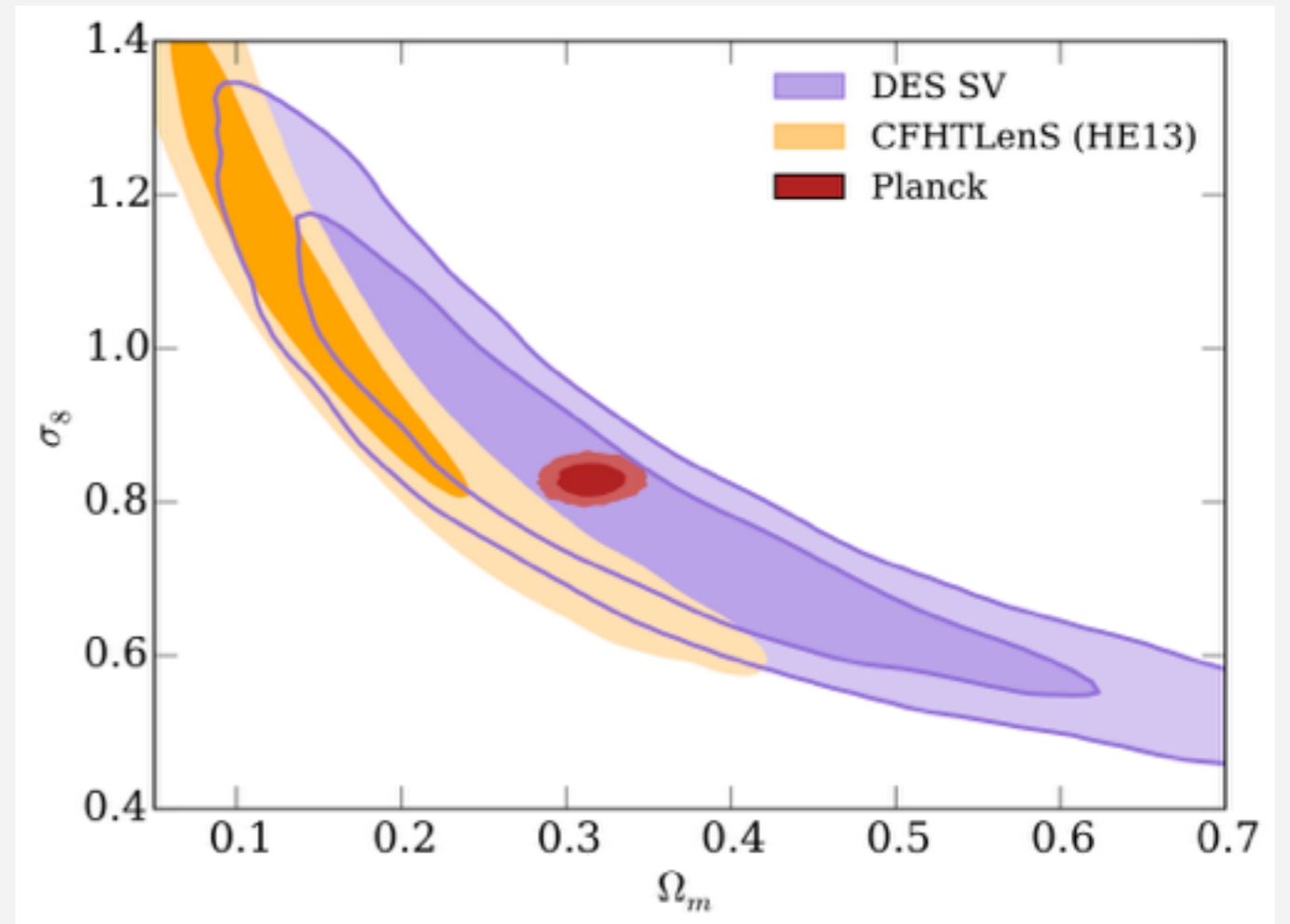
The smallest structure to study without
significant non-linearity effects

The σ_8 problem

Two σ_8 measurements: **CMB + Λ CDM** vs. **Weak Lensing**



DES: 1507.05552



The CFHTLenS & CMB results deviate by $\sim 2 - 3\sigma$.

H₀ problem

Two H₀ measurements

CMB + Λ CDM .

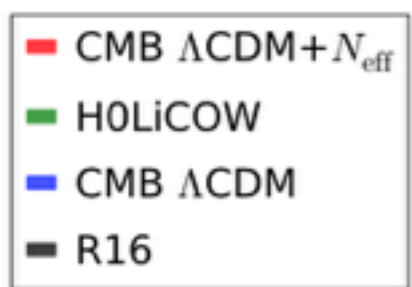
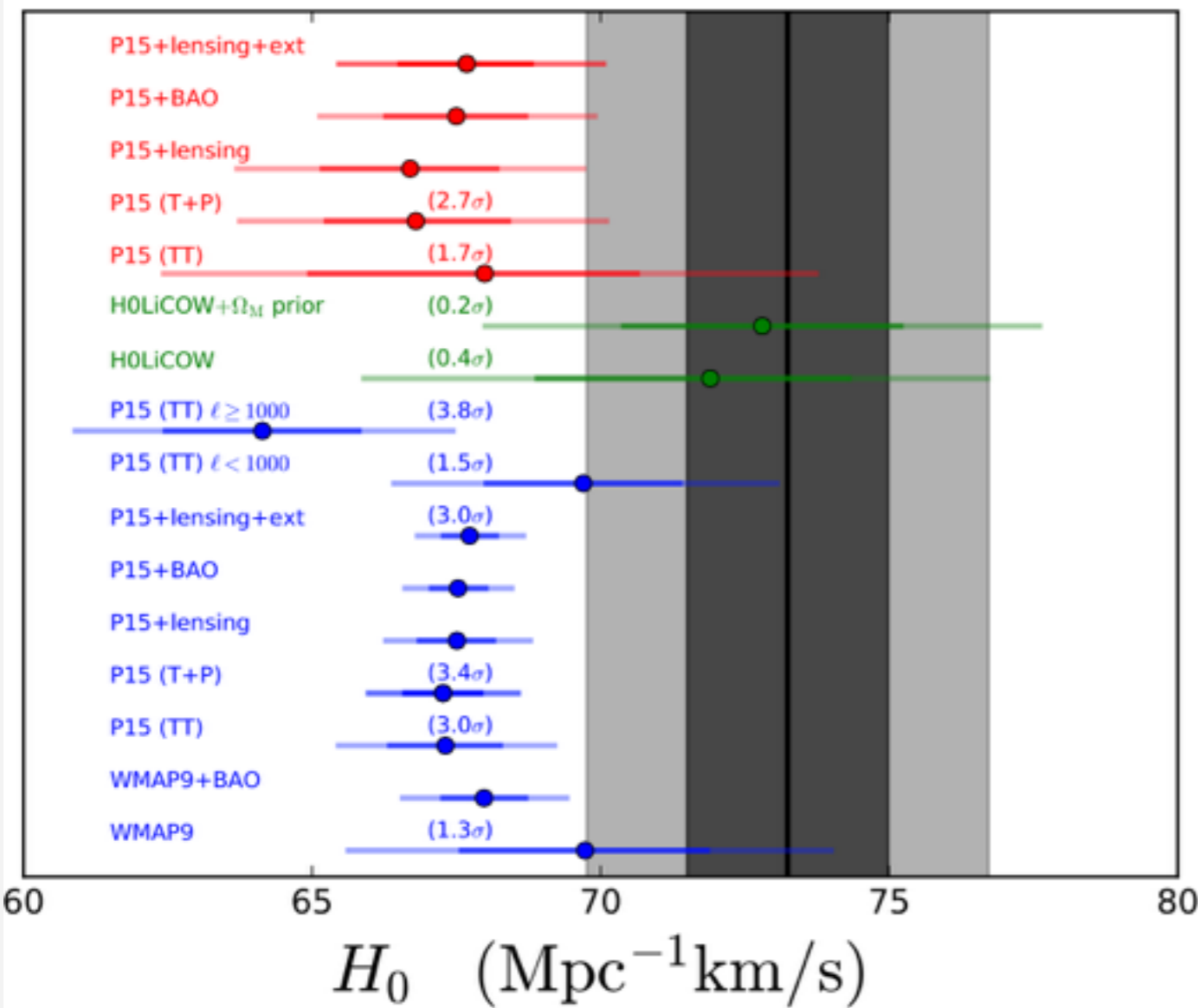
vs.

Local Measurements

$$H_0^{\text{Planck}} = 67.3 \pm 0.7 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$H_0^{\text{HST}} = 73.02 \pm 1.79 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

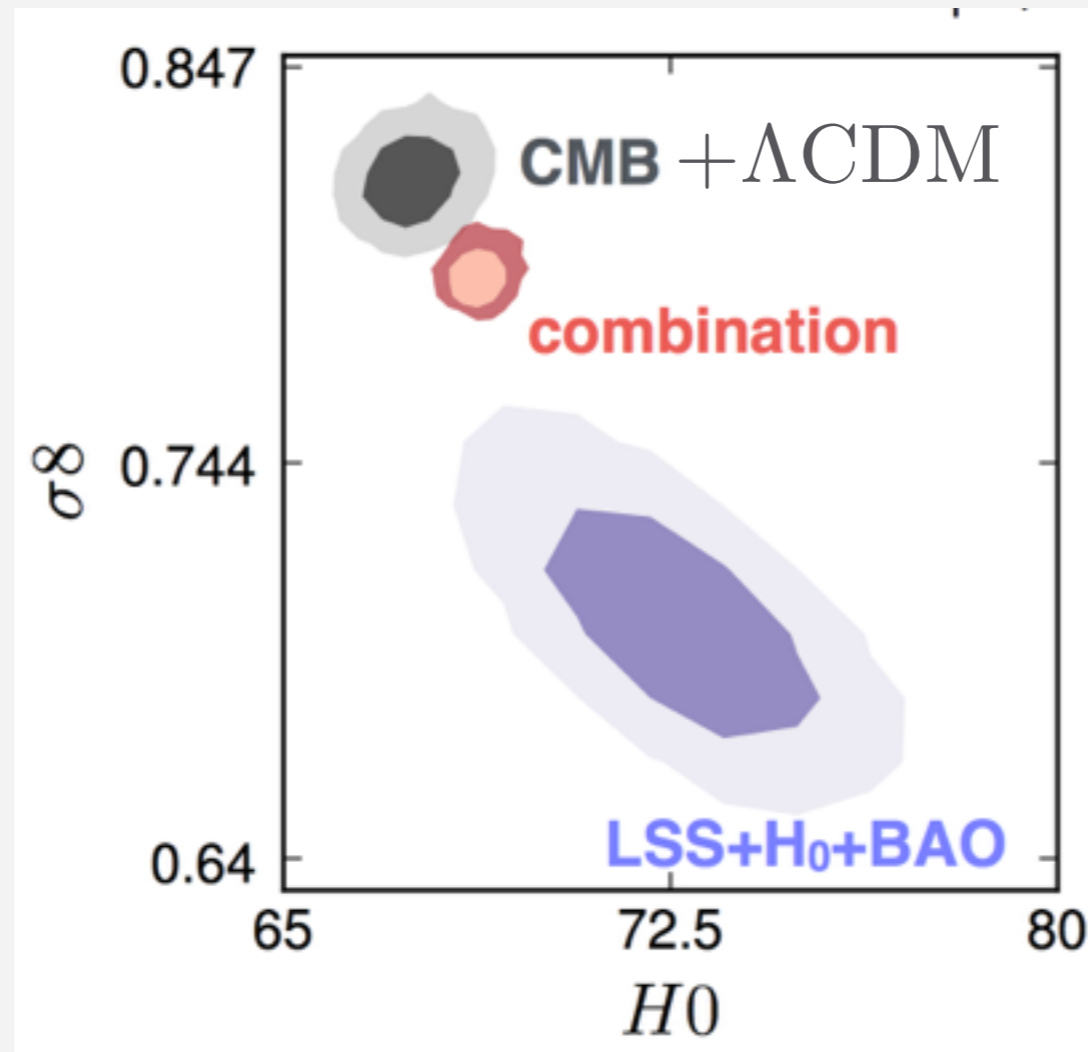
> 3 σ Discrepancy



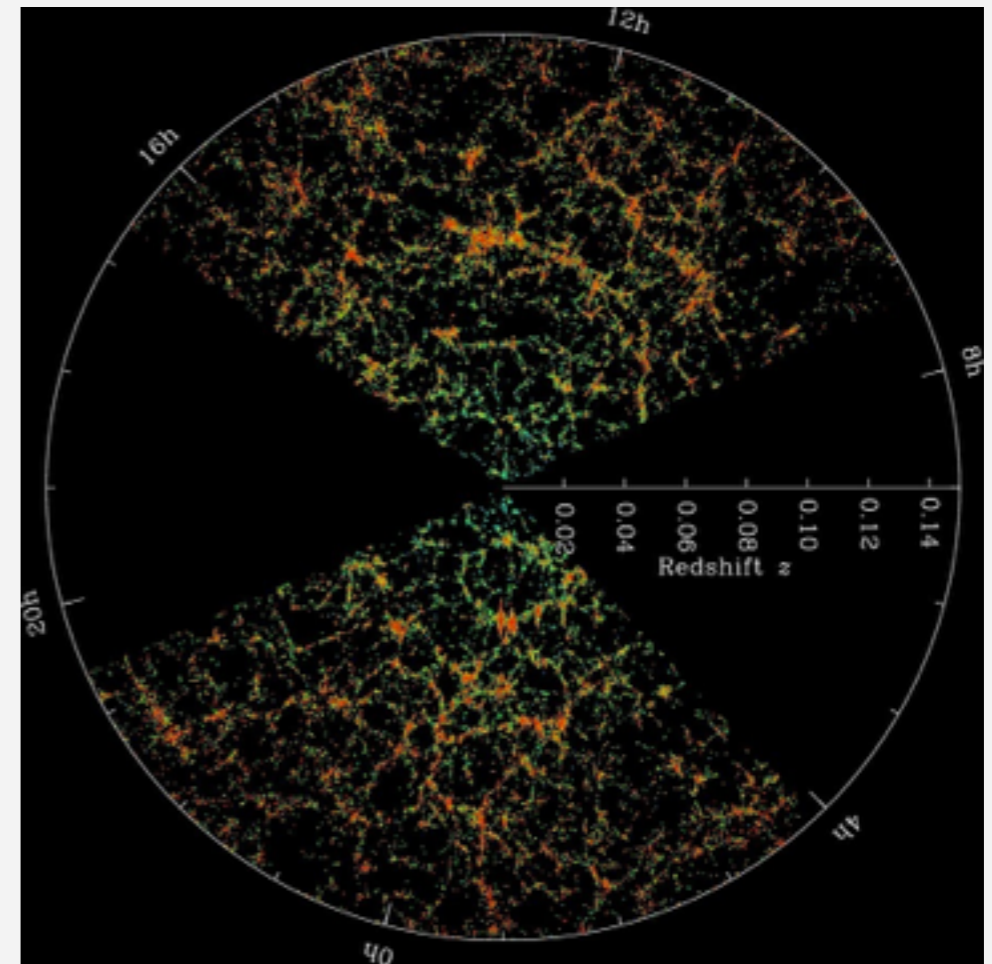
Bernal et. al. 1607.05617

Large Scale Structure problem

Poulin et. al. 1606.02073 : an illustration

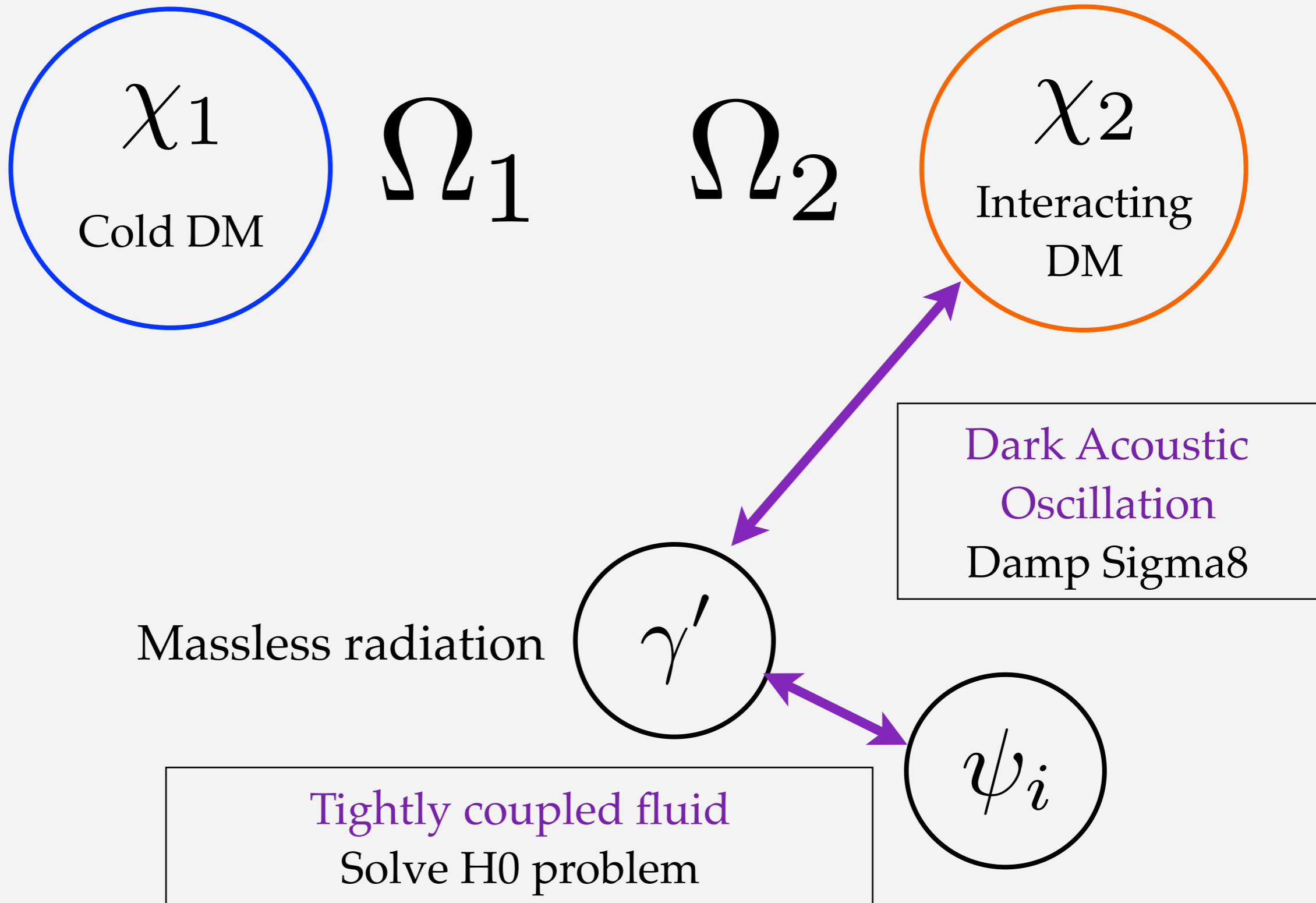


One solution: Partially Acoustic DM

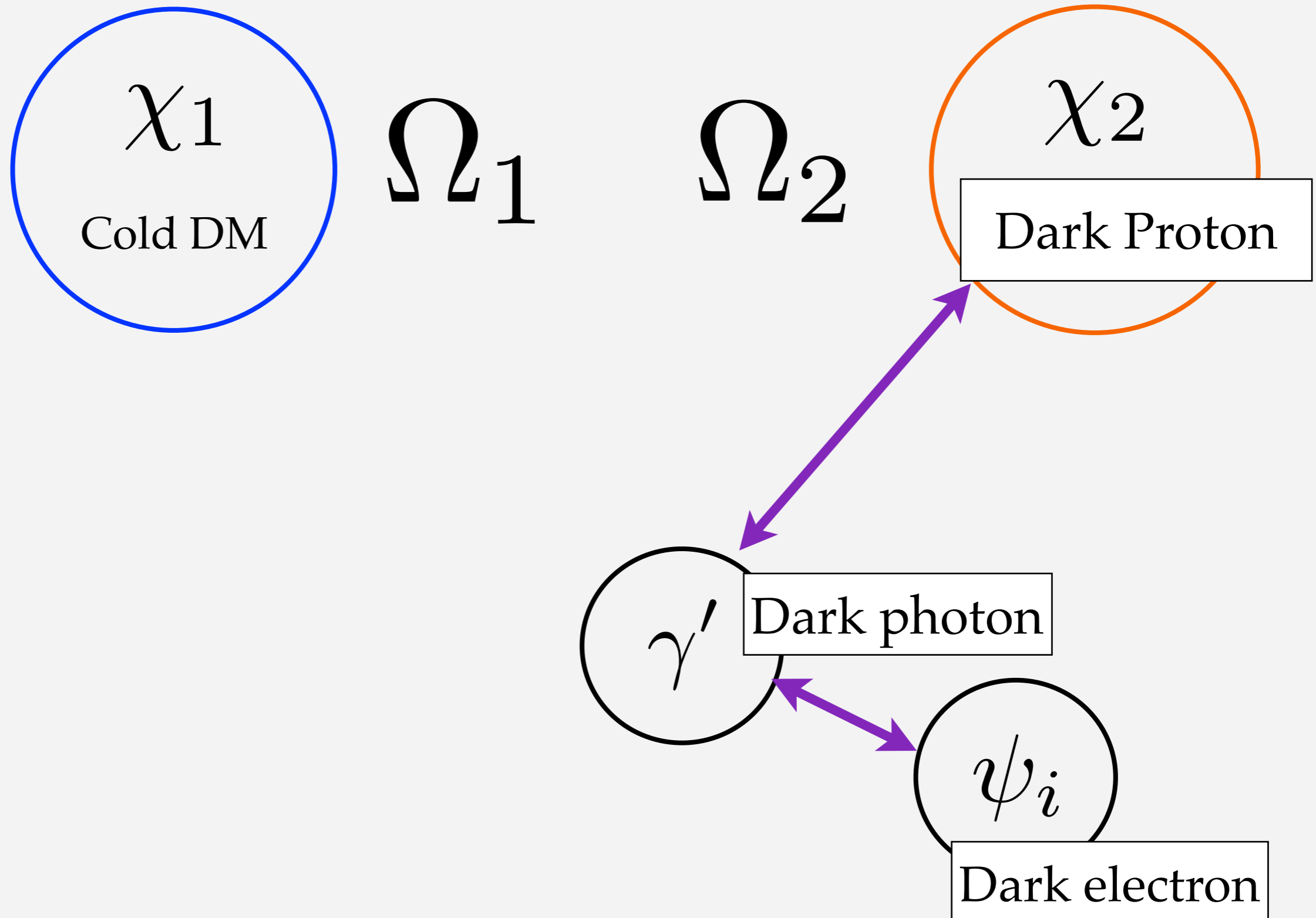


DISTRIBUTION OF GALAXIES IN OUR UNIVERSE. CREDIT: SDSS

Consider a non-minimal dark sector

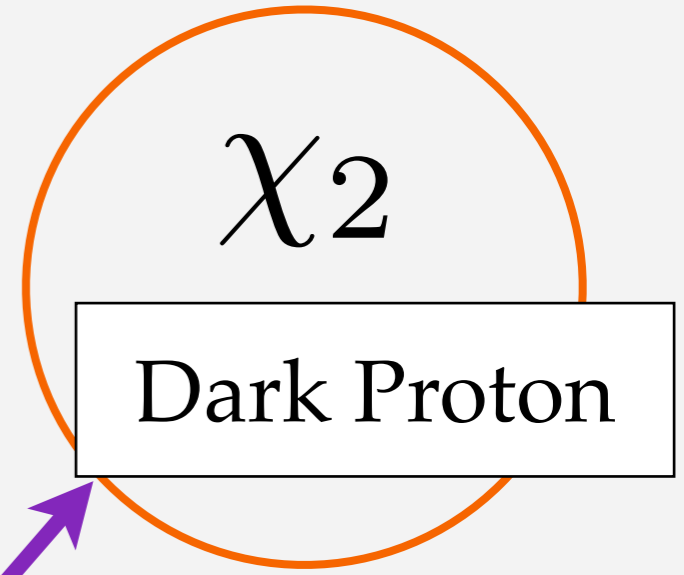


Consider a non-minimal dark sector



Consider a non-minimal dark sector

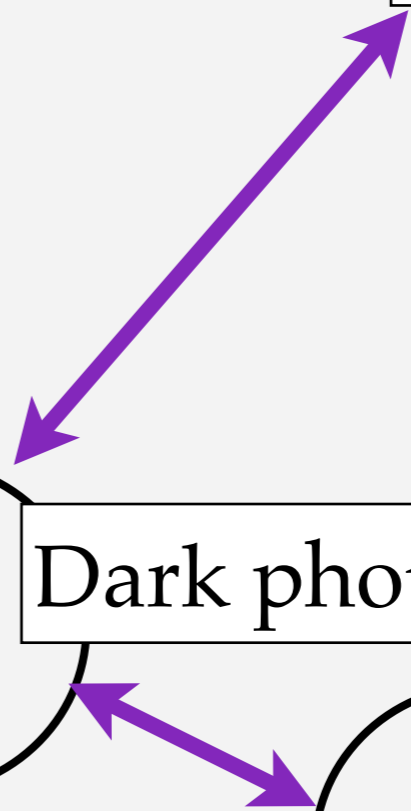
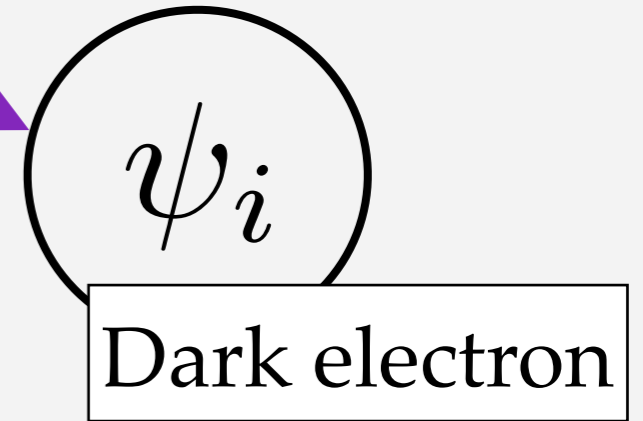
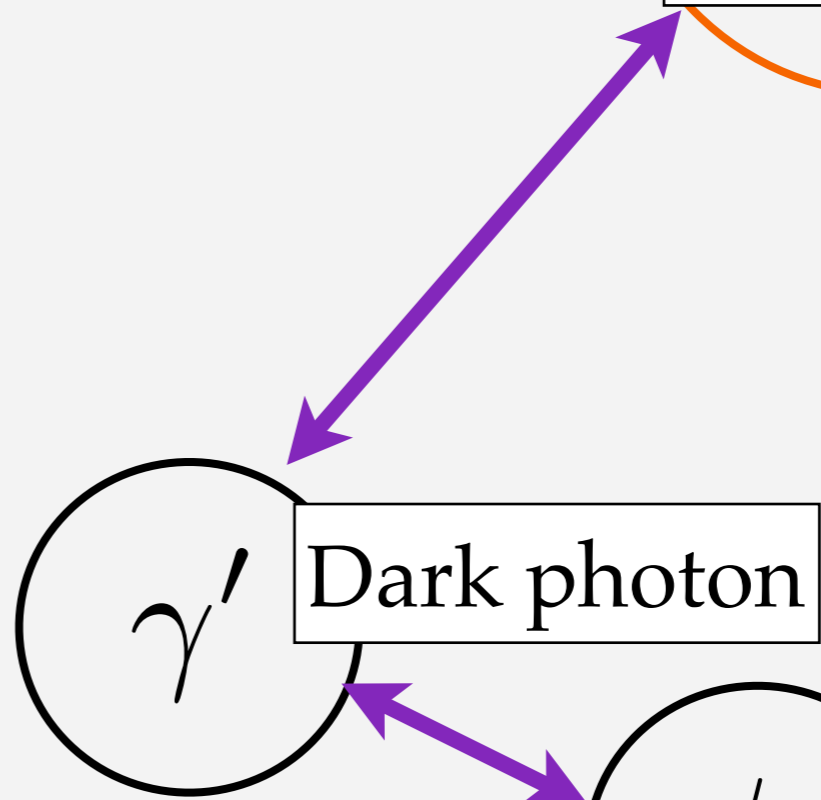
Ω_2



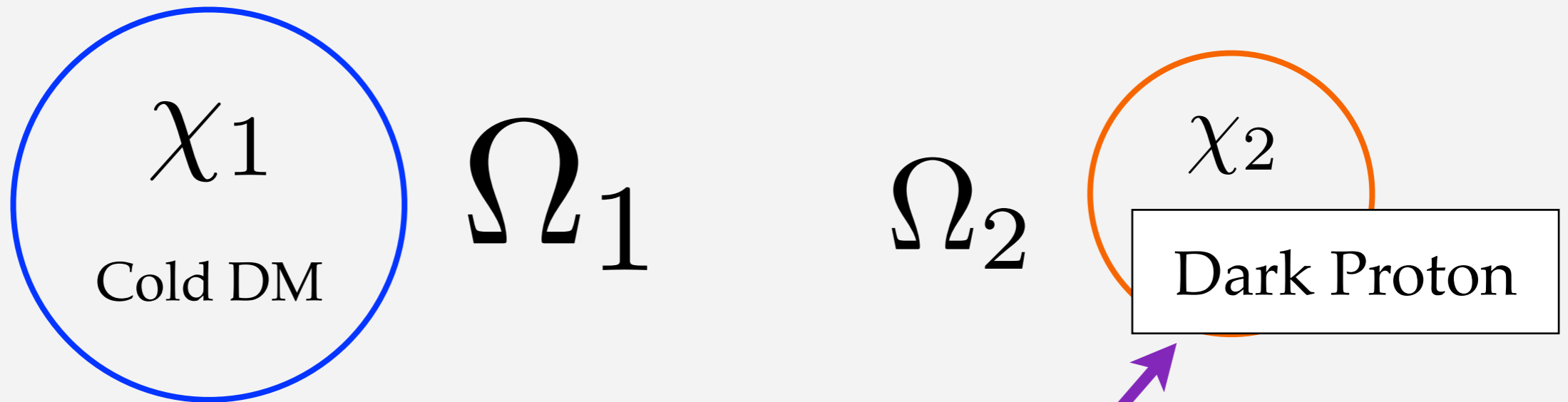
$$r \equiv \Omega_2 / \Omega_{\text{DM}}$$

$$r = 1$$

Fully Acoustic Oscillation

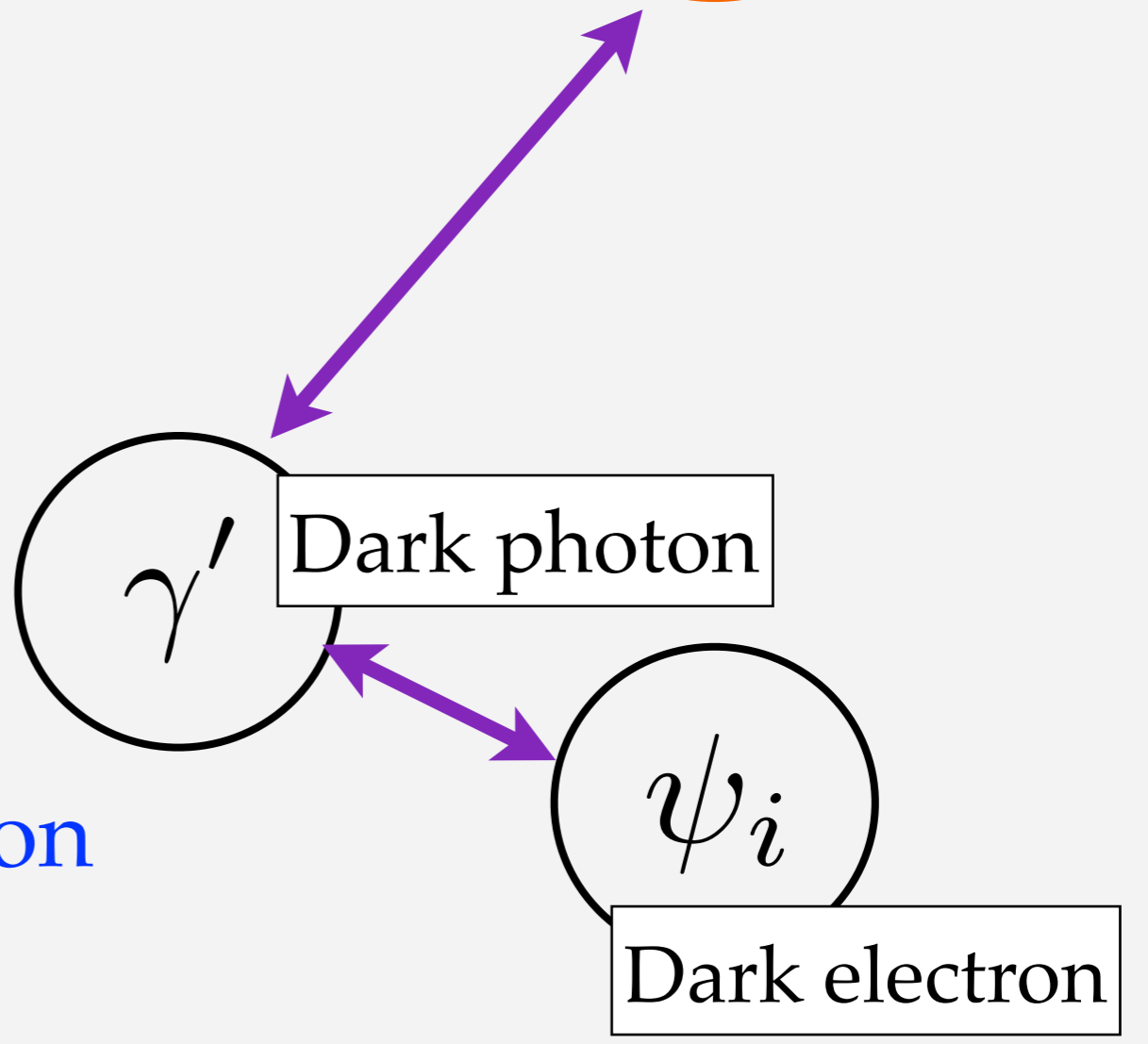


Consider a non-minimal dark sector



$$r \equiv \Omega_2 / \Omega_{\text{DM}}$$

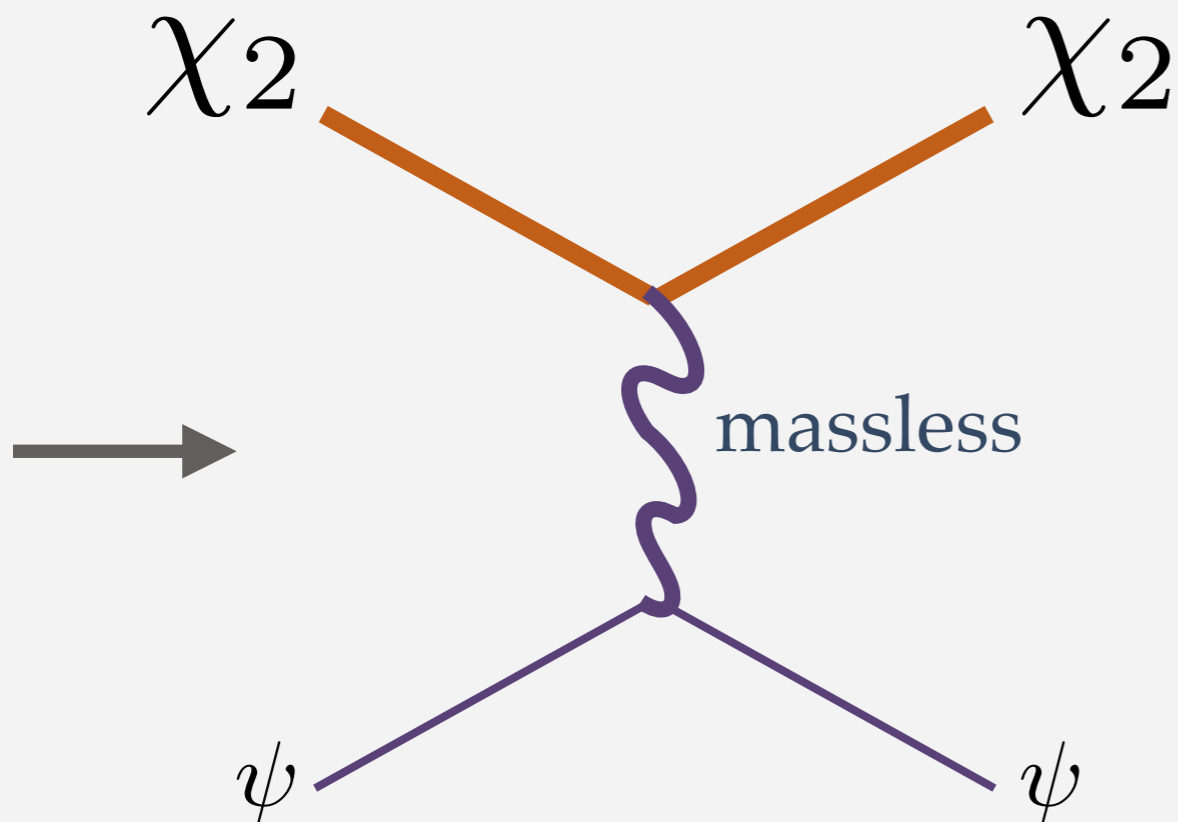
$$r < 1$$



Partially Acoustic Oscillation

For the acoustic oscillation to exist

We need the DM-DR scattering to remain non-decoupled



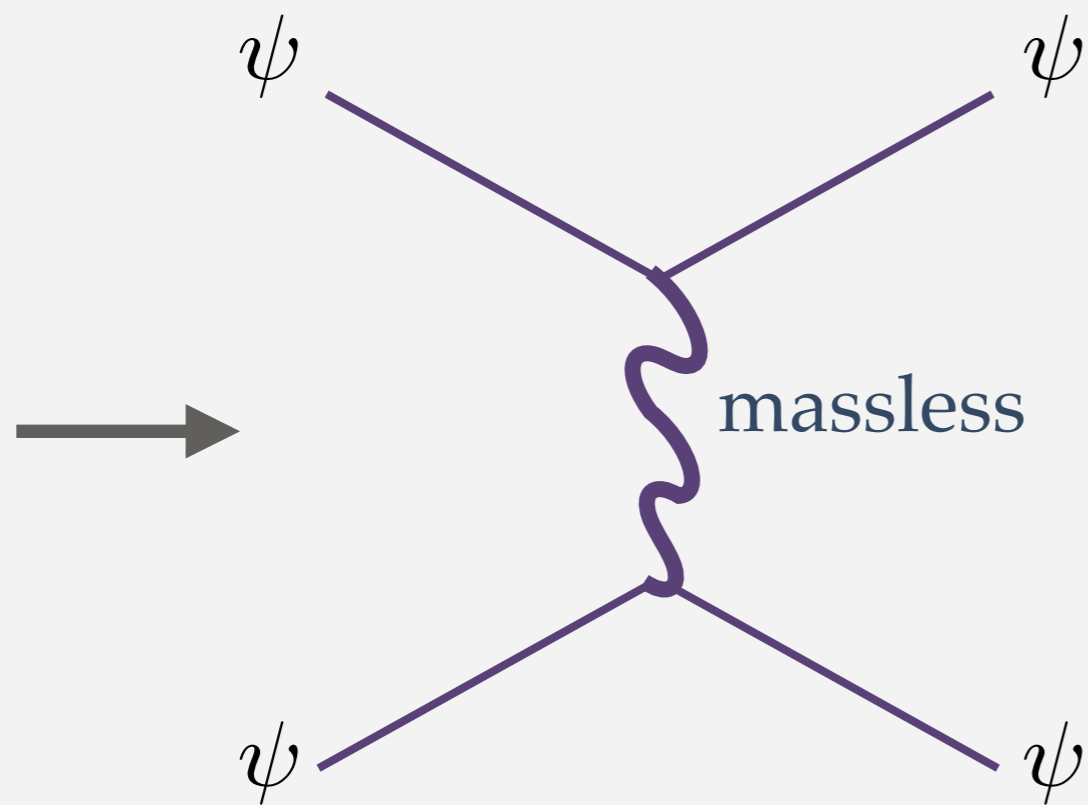
$$\Gamma \simeq \hat{\alpha}^2 \ln(\hat{\alpha}^{-1}) \frac{T_D^2}{m_{\text{DM}}}$$

Same temp-dependence as Hubble
in the radiation-dominant era

Easy to keep $\Gamma \gg H$ all the time, if $\hat{\alpha} \gg 10^{-8} \sqrt{\frac{m_{\chi_2}}{10 \text{ GeV}}}$

Tightly coupled dark radiation

We need the DM-DR scattering to remain non-decoupled

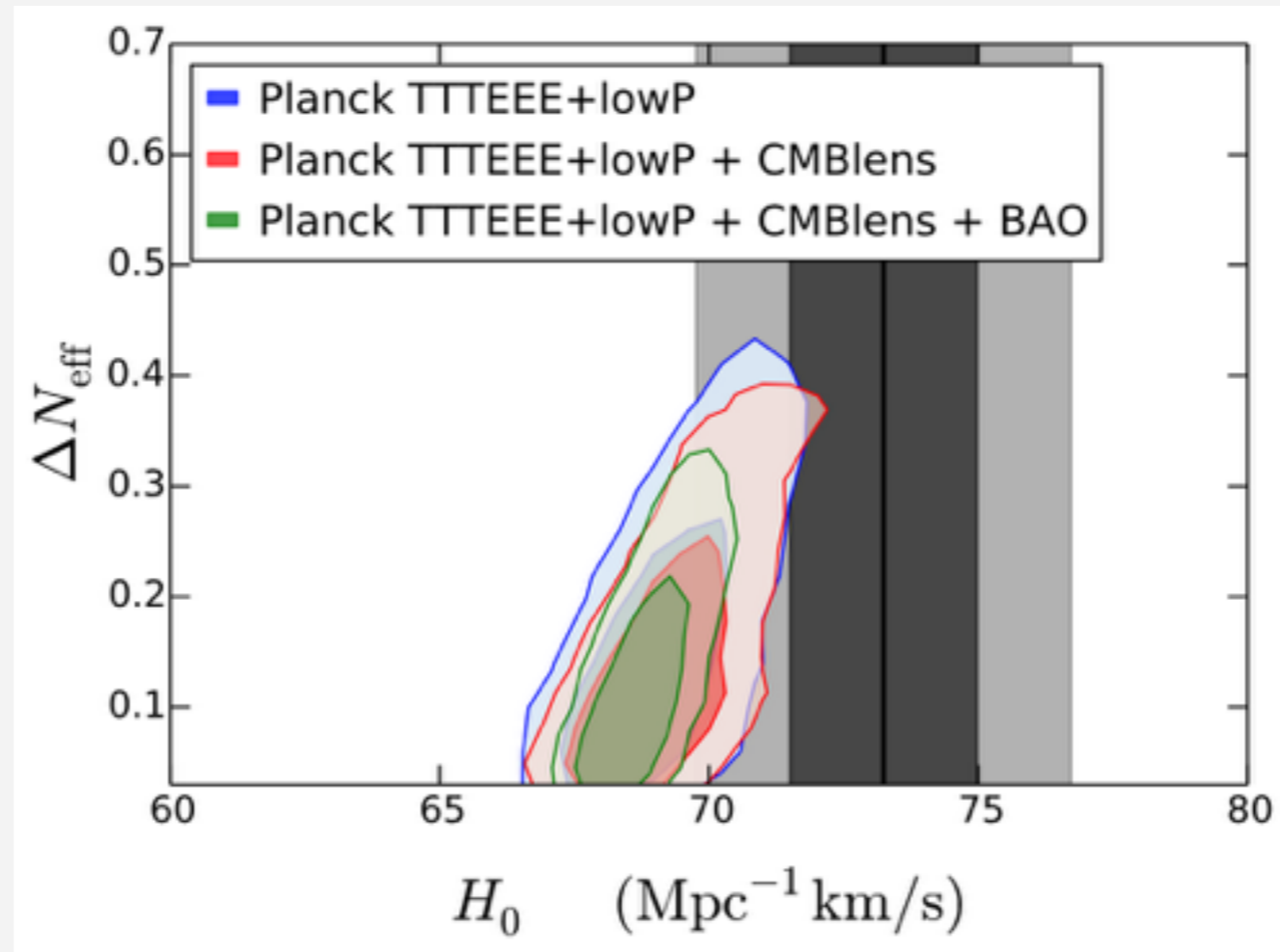


$$\Gamma \simeq \hat{\alpha}^2 \ln(\hat{\alpha}^{-1}) T_D$$

The same coupling keeps dark fermions / photon a tightly coupled fluid

Solving H_0 problem with extra dark radiation

Bernal et. al. 1607.05617

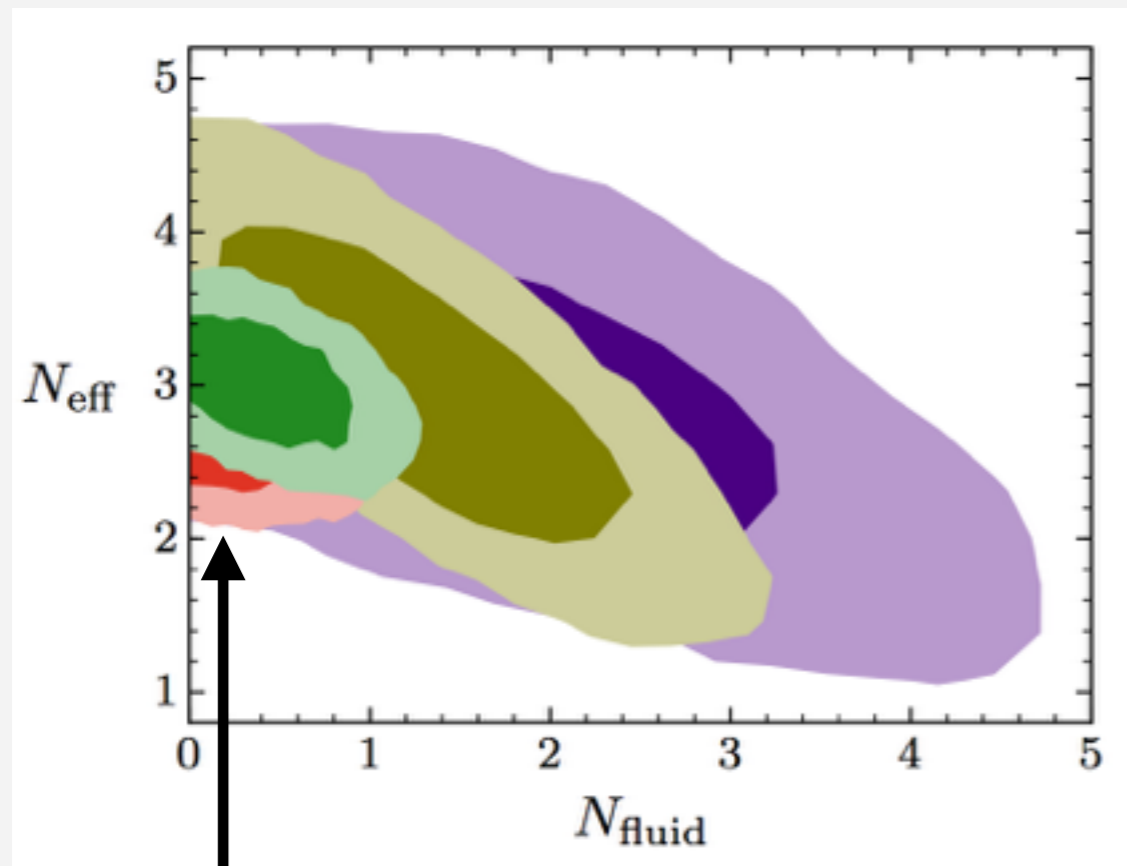


Can explain the larger H_0 by including $\Delta N_{\text{eff}} > 0.4$ dark radiation

Adam Riess et.al. 1604.01424

Dark fluid is better than FS-radiation

Baumann et. al. 1508.06342



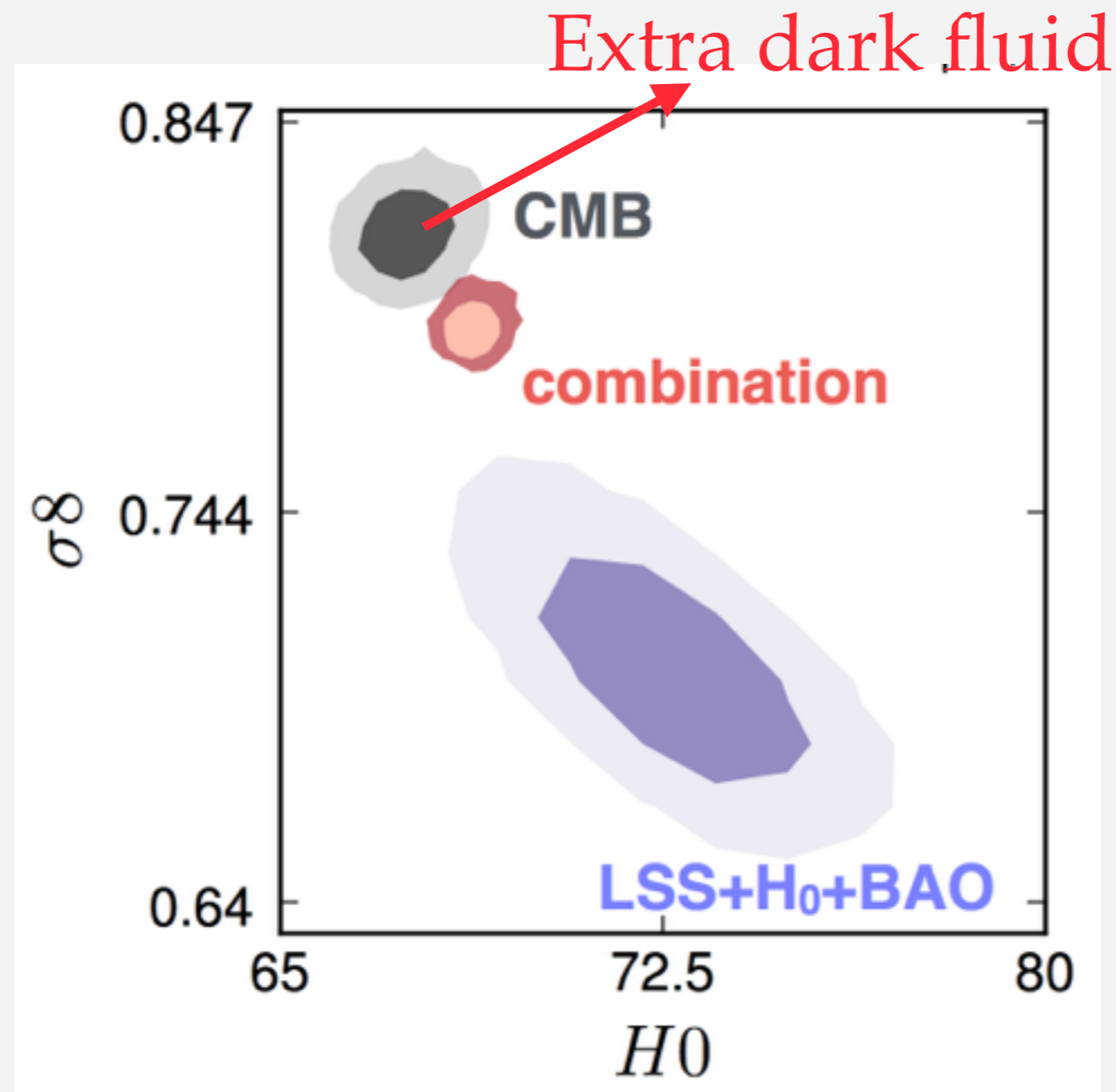
Planck TT, TE, and EE likelihoods

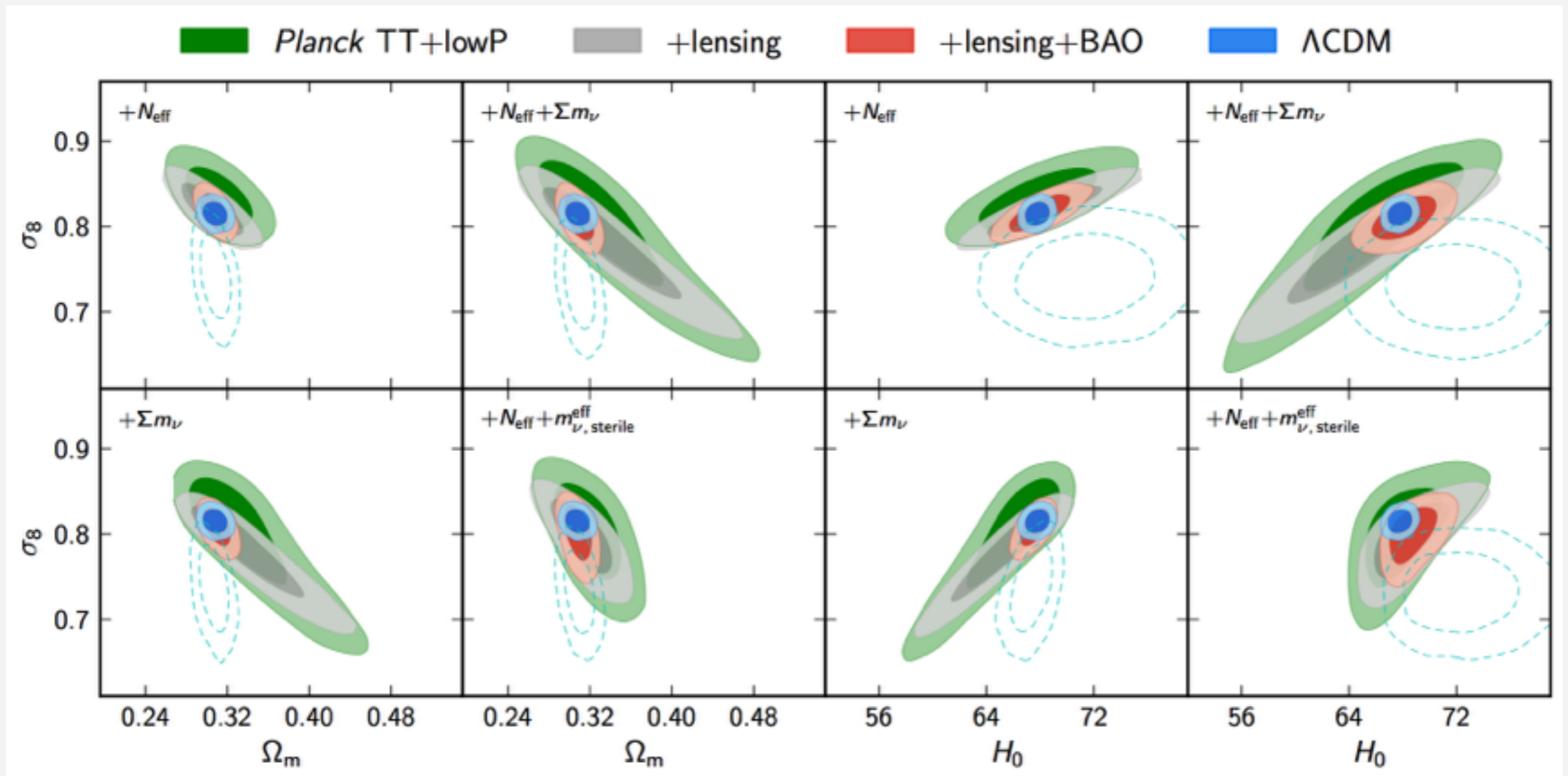
ΔN_{eff} bound on a tightly coupled fluid is weaker

	TT, TE, EE		TT-only	
	varying Y_p	fixed Y_p	varying Y_p	fixed Y_p
N_{eff}	$2.78^{+0.30}_{-0.35}$	$2.99^{+0.30}_{-0.29}$	$2.87^{+0.76}_{-0.74}$	$2.94^{+0.71}_{-0.69}$
N_{fluid}	< 0.88	< 1.06	< 3.93	< 2.65

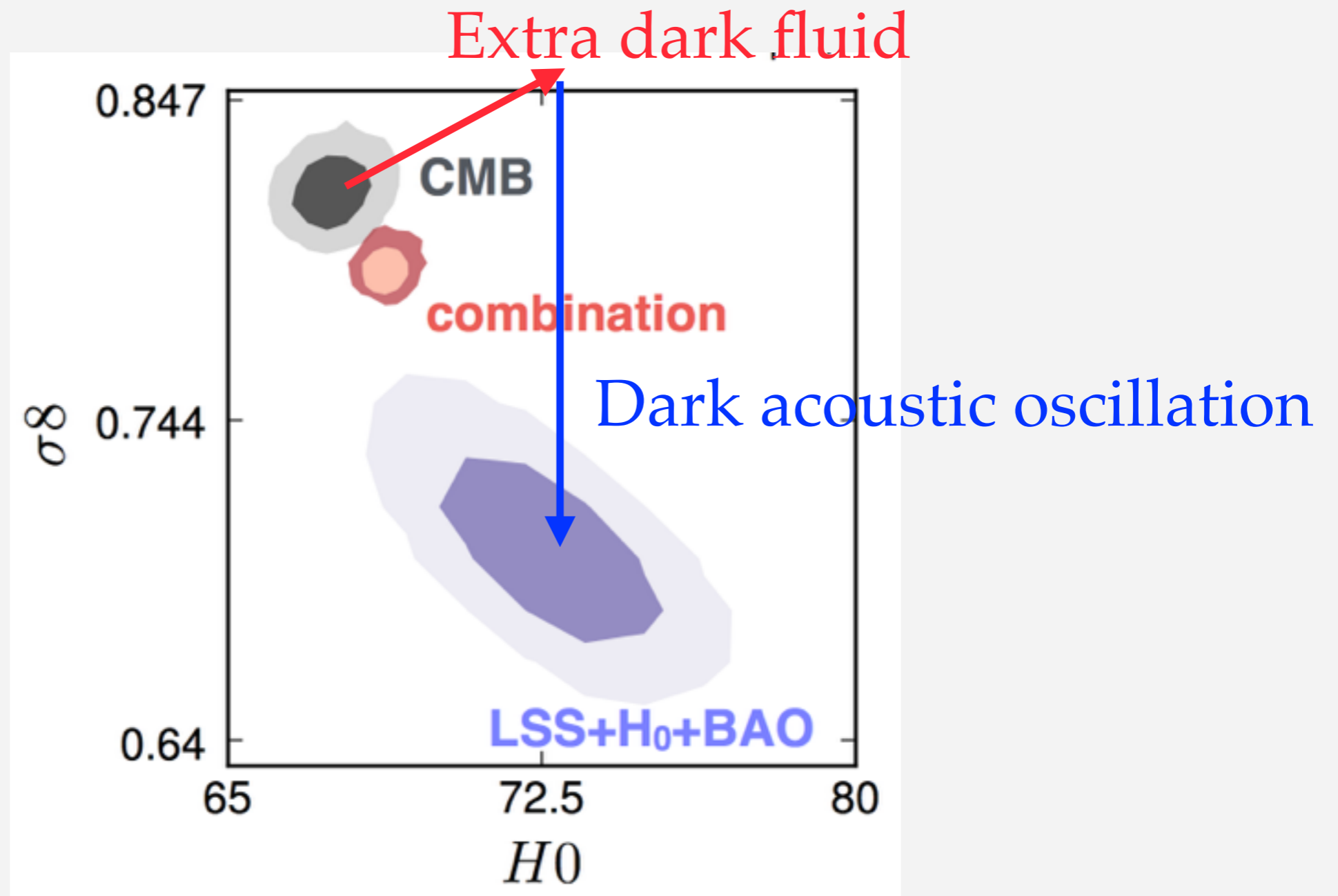
(2σ)

Reconcile H_0 , but makes σ_8 worse

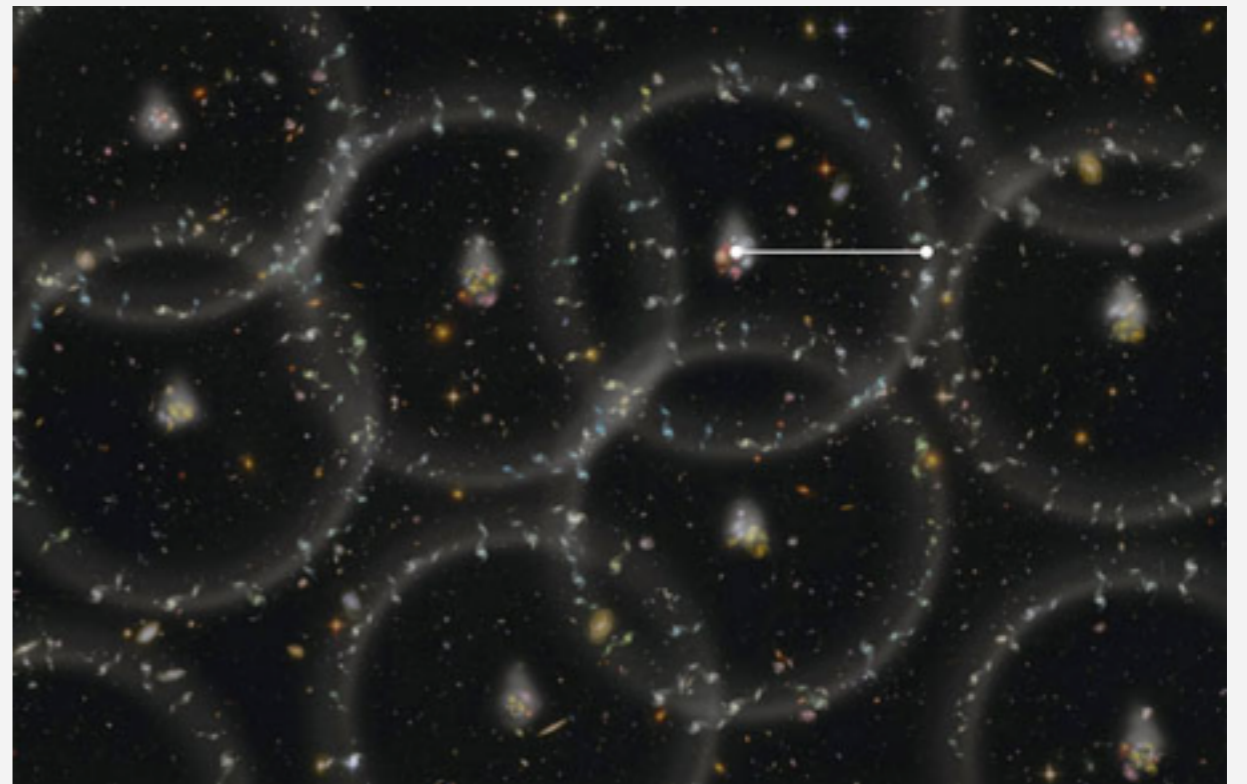




DM-DR scattering suppresses Σ_8

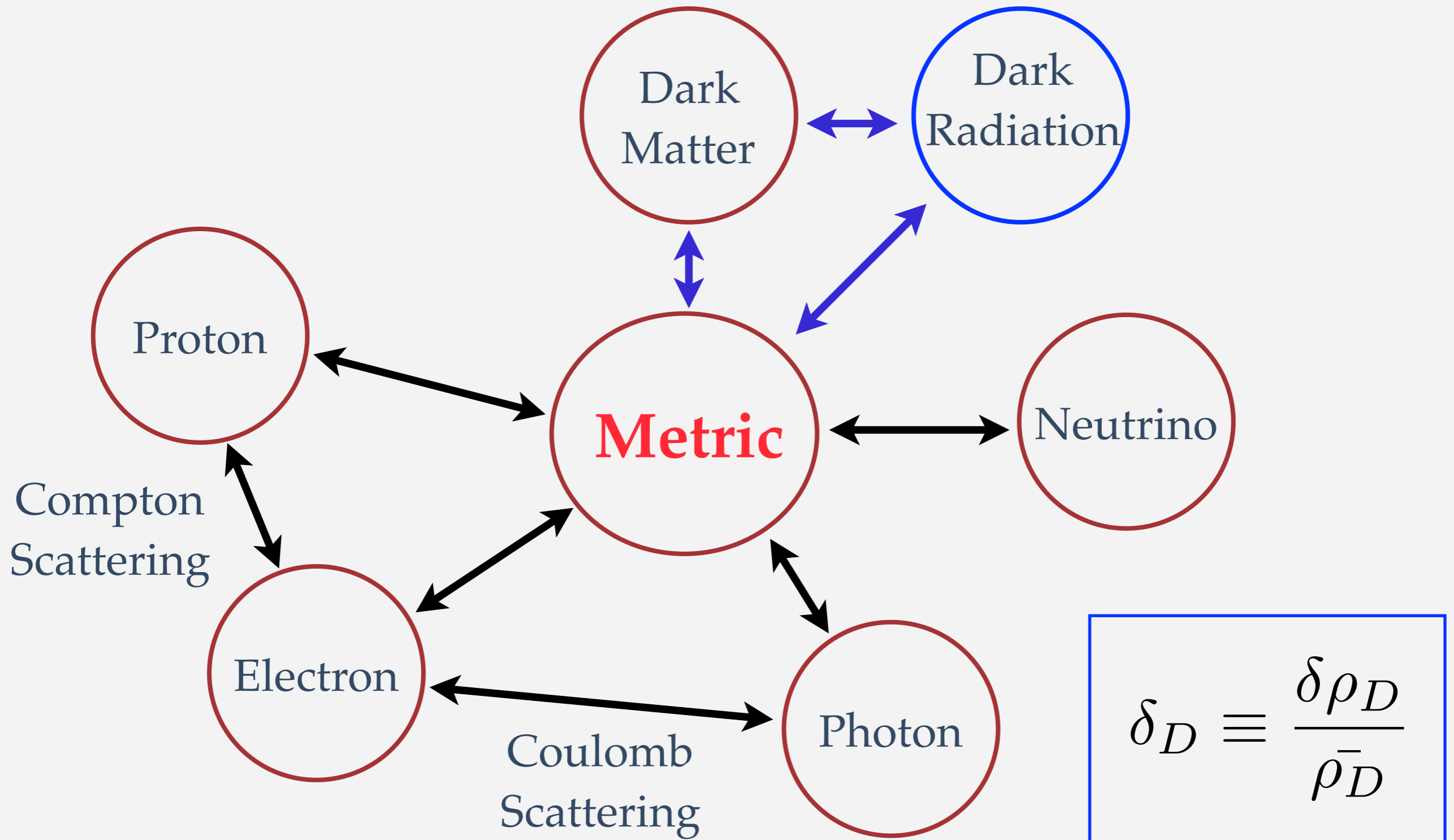


Structure Formation with Acoustic DM



A cartoon produced by the BOSS project showing the spheres of baryons around the initial dark matter clumps

Evolution of the Large Scale Structures



Boltzmann Equation in Conformal Newtonian Gauge

$$\dot{\delta}_D = -\theta_D + 3\dot{\psi} \quad \left(\frac{d\rho}{dt} = -\rho \nabla \cdot \vec{v} \right)$$

$$\dot{\theta}_D = -\frac{\dot{a}}{a}\theta_D + k^2\psi + \underline{a\Gamma(\theta_R - \theta_D)} \quad \Gamma \equiv \frac{1}{\langle p_D^2 \rangle} \frac{d\langle \delta p_D^2 \rangle}{dt}$$

$$\dot{\delta}_R = -\frac{4}{3}\theta_R + 4\dot{\psi}$$

$$\dot{\theta}_R = \frac{k^2}{4}\delta_R + k^2\psi + \underline{Ra\Gamma(\theta_D - \theta_R)}$$

$$\theta_s \equiv \partial_i v_s^i$$

Velocity Divergent

Metric Perturbation

$$ds^2 = a^2(\tau) [-(1 + 2\psi)d\tau^2 + (1 - 2\phi)\delta_{ij}dx^i dx^j]$$

No free-streaming particle => $\phi = \psi$

Boltzmann Equation in Conformal Newtonian Gauge

$$\dot{\delta}_D = -\theta_D + 3\dot{\psi} \quad \left(\frac{d\rho}{dt} = -\rho \nabla \cdot \vec{v} \right) \quad \Gamma \equiv \frac{1}{\langle p_D^2 \rangle} \frac{d\langle \delta p_D^2 \rangle}{dt}$$

$$\dot{\theta}_D = -\frac{\dot{a}}{a} \theta_D + k^2 \psi + \underline{a\Gamma(\theta_R - \theta_D)}$$

vanishes for cold DM,
similar expression for SM baryon

$$\dot{\delta}_R = -\frac{4}{3} \theta_R + 4\dot{\psi}$$

$$\dot{\theta}_R = \frac{k^2}{4} \delta_R + k^2 \psi + \underline{R a\Gamma(\theta_D - \theta_R)}$$

similar expression for
SM photon

Tightly coupled DM-DR (similar to the baryon-photon system):

$$\Gamma \gg H \quad \Rightarrow \quad a\Gamma \gg \tau^{-1}$$

In the tightly coupled DM-DR limit

We can simplify the evolution of DM perturbation

$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \phi$$

$$R \equiv \frac{3\rho_D}{4\rho_R}$$

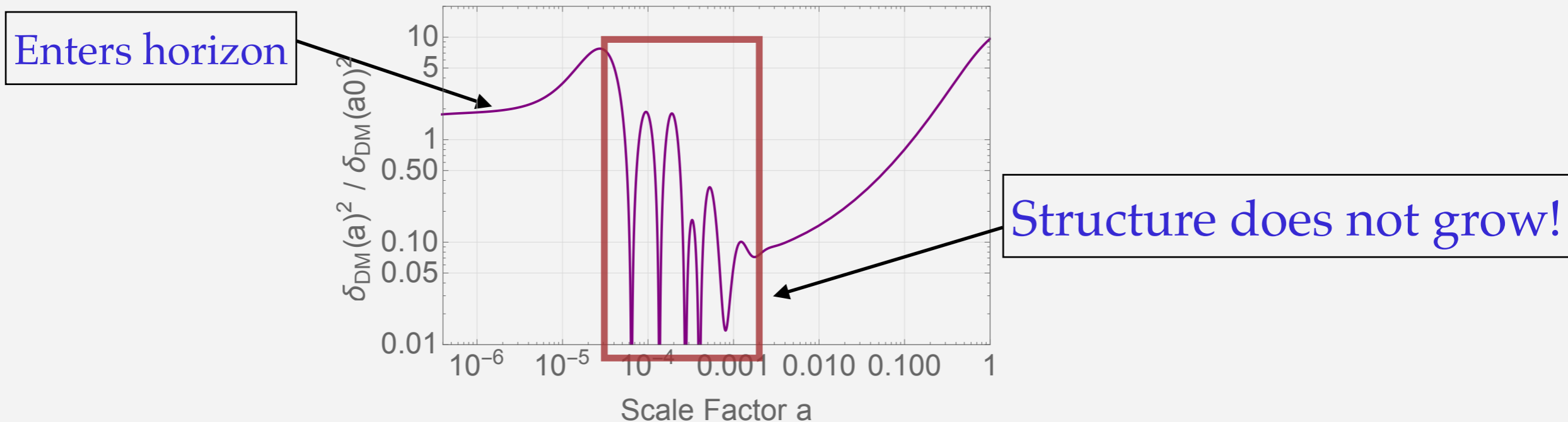
Parametrize the “mass”
of DM-DR fluid

Radiation Domination , $R \ll 1$

Density perturbation oscillates => No structure grows

$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \phi$$

The density perturbation oscillates as a harmonic oscillator!
Same physics as the baryon acoustic oscillation

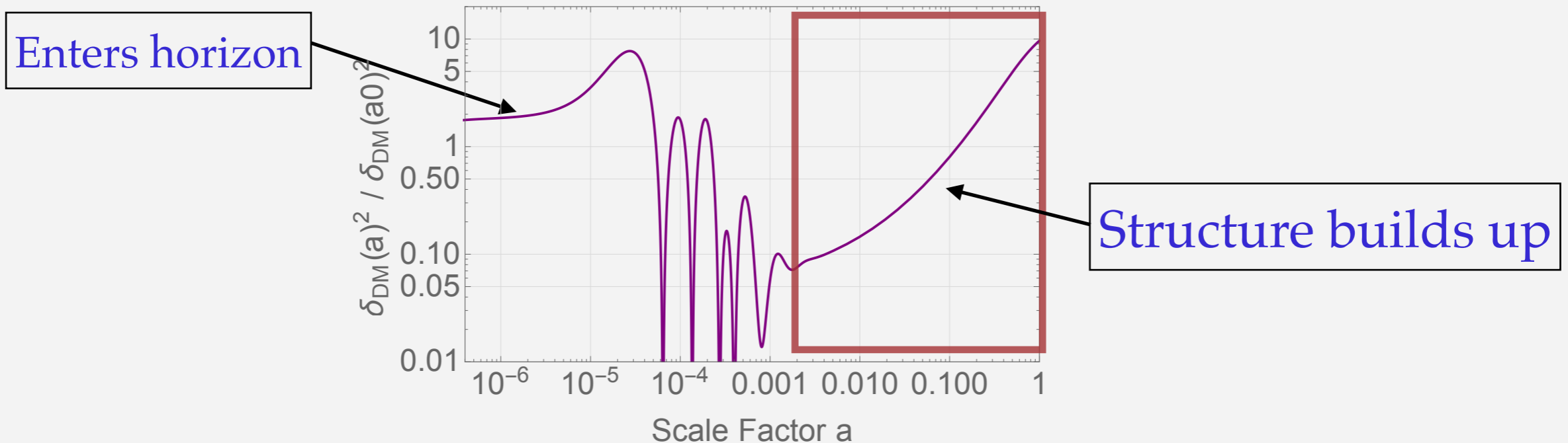


Matter Domination , $R \gg 1$

No oscillation => Linear growth

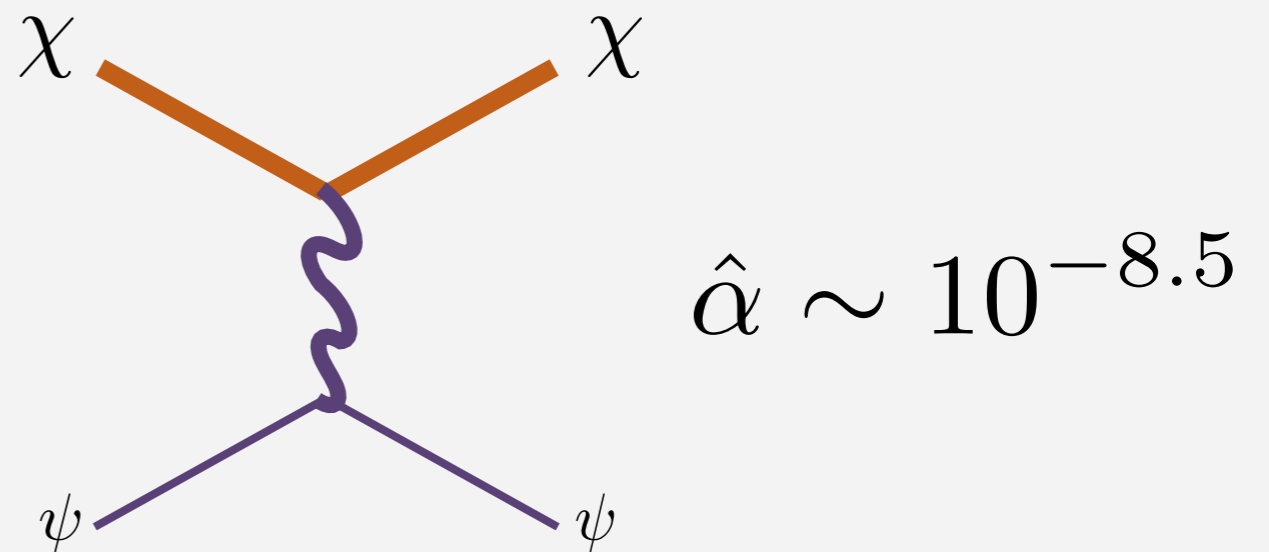
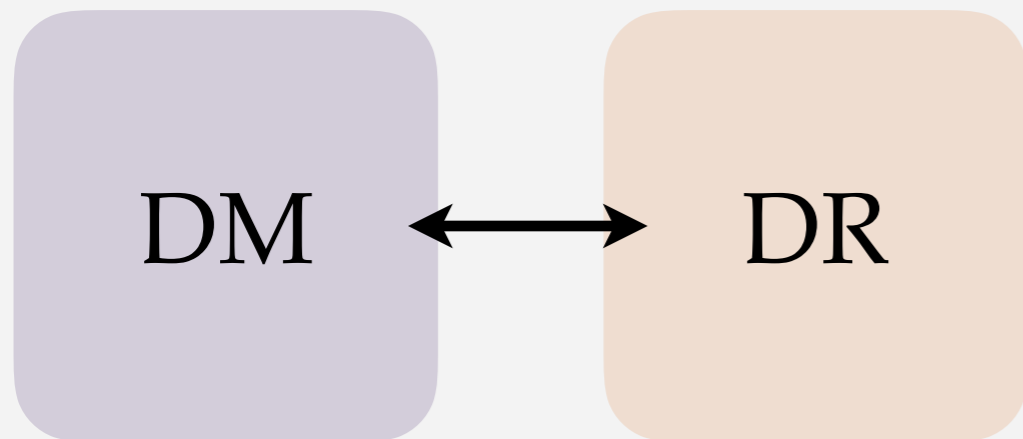
$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \phi$$

No oscillation, no damping from the DR scattering
Same structure formation as cold DM



If all the DM particles oscillate with DR

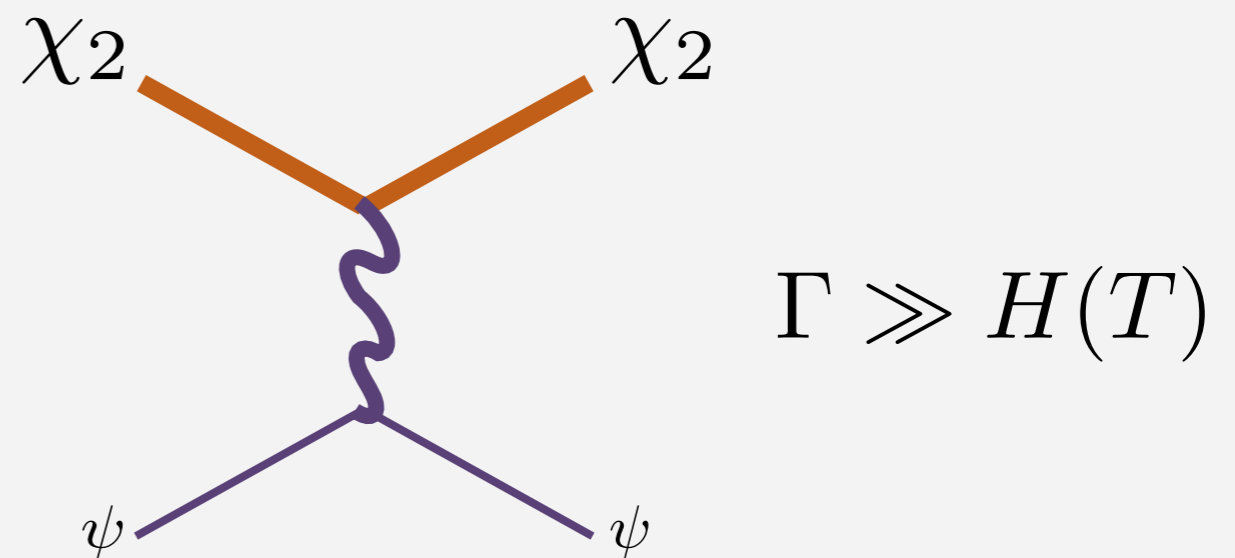
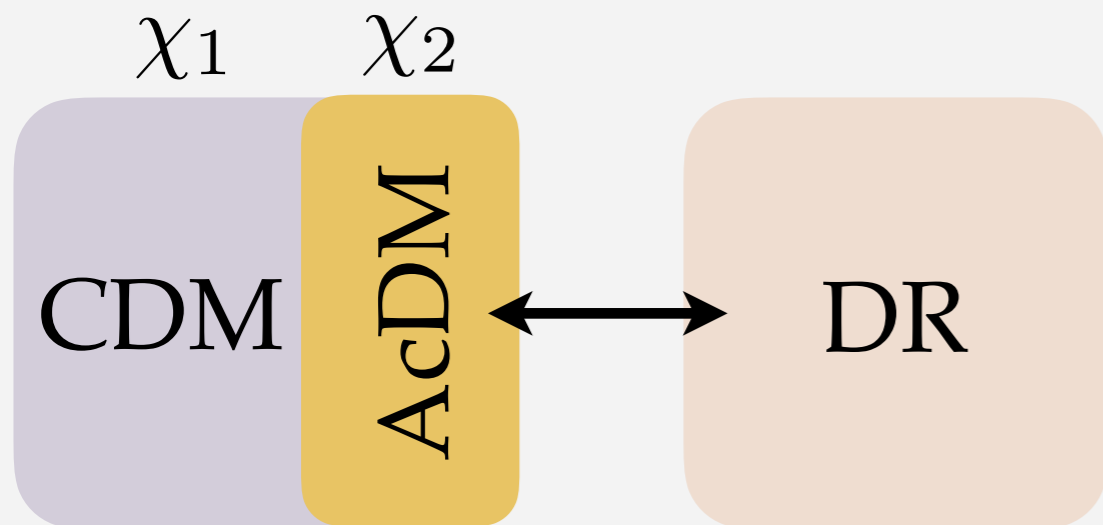
We need a small DM coupling for the right σ_8 suppression



Manuel A. Buen-Abad, Gustavo Marques-Tavares, and Martin Schmaltz (2015)

Julien Lesgourgues, Gustavo Marques-Tavares, and Martin Schmaltz (2015)

How about only a fraction of DM particles having the acoustic oscillation?

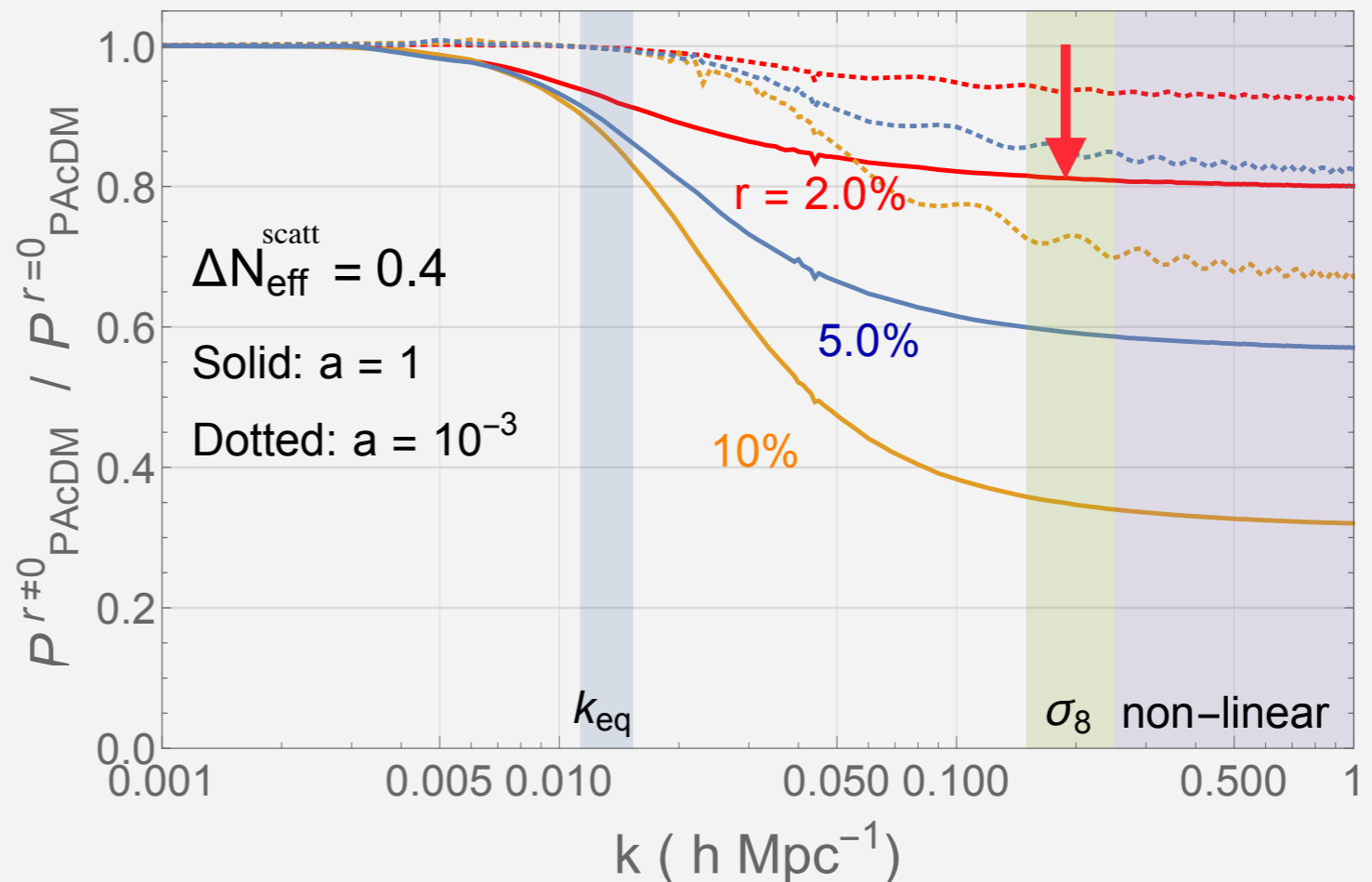


$$\Gamma \gg H(T)$$

If only a small component DM has the acoustic oscillation, we can allow DM-DR to be tightly coupled (remain equilibrium) and solve DM perturbation analytically

Solving Sigma8 problem with PAcDM

$$\left[\frac{\sigma_8(\Lambda\text{CDM})}{\sigma_8(\text{PAcDM})} \right]^2 \sim$$

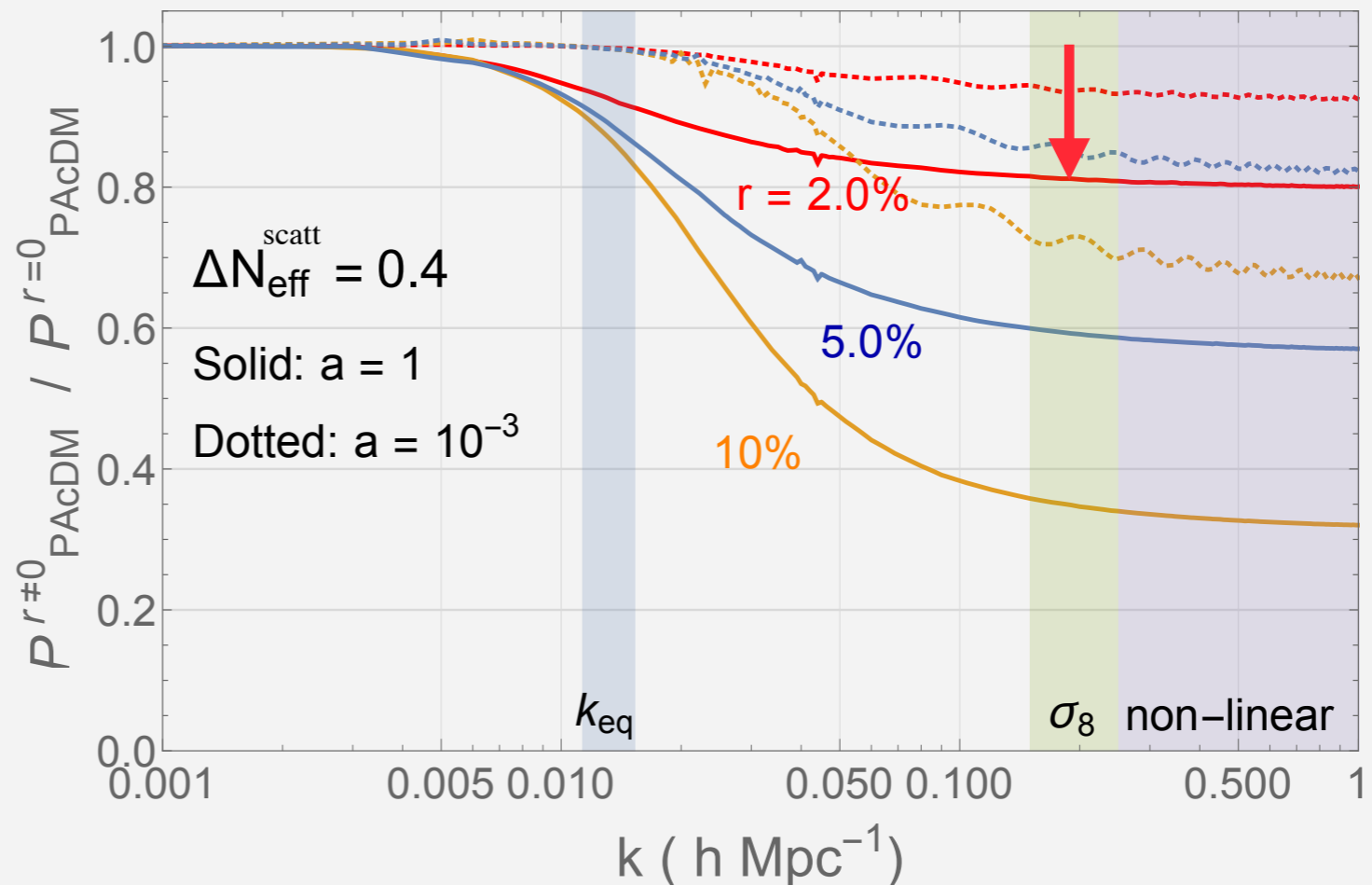


A simple analytical result

$$\frac{P(r)}{P(0)} \simeq (1 - 2r) \left(\frac{a}{a_{\text{eq}}} \right)^{-1.2r} \quad r \equiv \Omega_2 / \Omega_{\text{DM}}$$

Solving Sigma8 problem with PAcDM

$$\left[\frac{\sigma_8(\Lambda\text{CDM})}{\sigma_8(\text{PAcDM})} \right]^2 \sim$$



Need $\sim 2\%$ acoustic DM to solve the σ_8 problem

2% density is easy to obtained

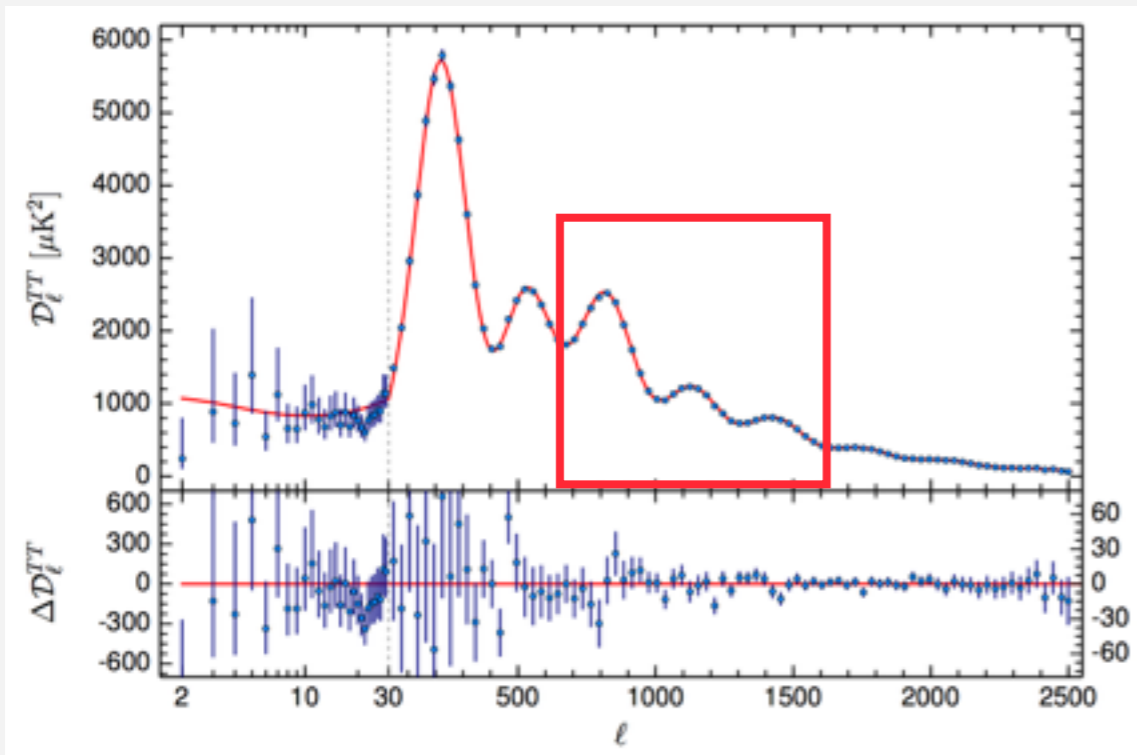
When both DM particles are WIMP-like and having thermal freeze out through a heavy mediator

$$\frac{\Omega_2}{\Omega_1} \simeq \left(\frac{m_2}{m_1} \right)^2$$

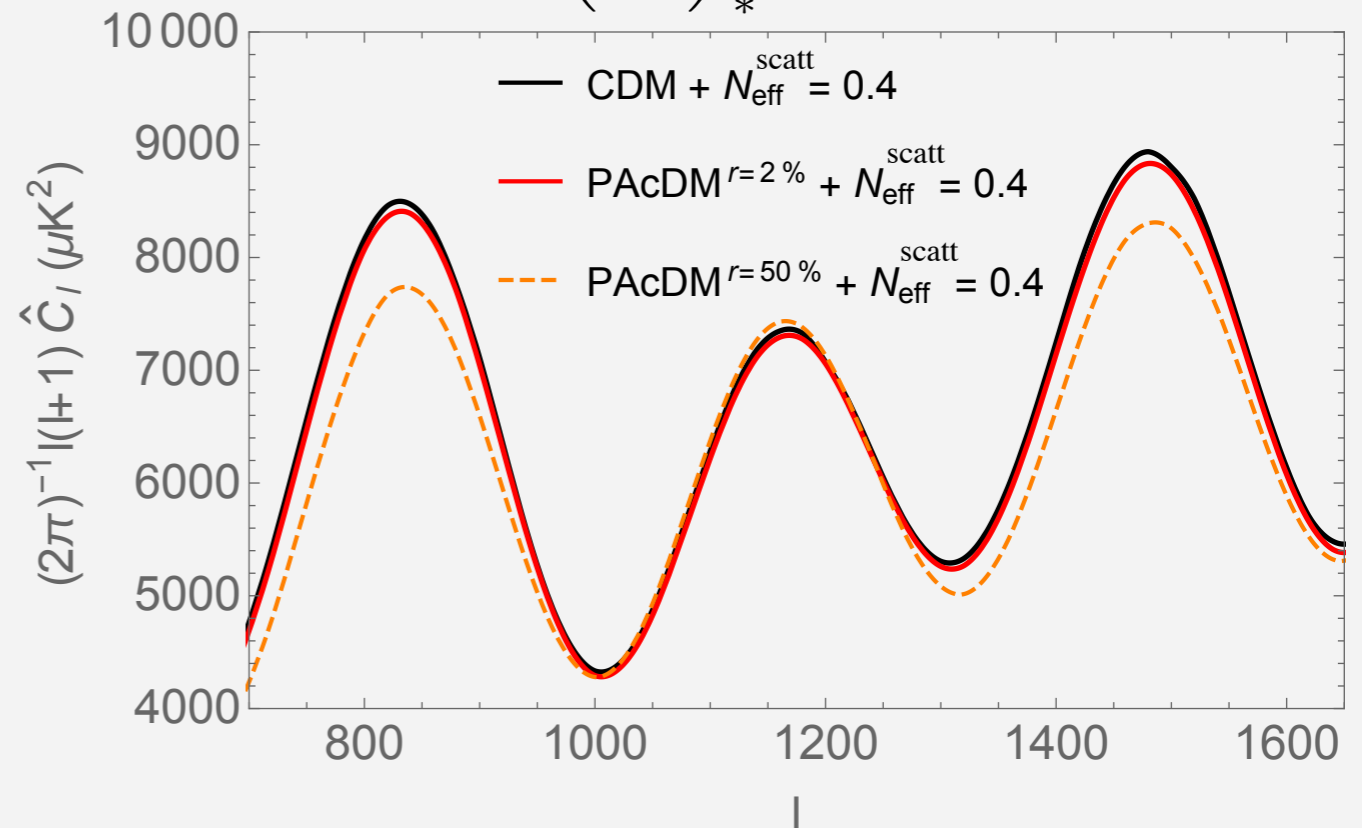
Only need $m_1 \simeq 7 m_2$ to obtain the 2% ratio
(assuming equal couplings)

Correction to the CMB spectrum

Planck 1502.01589



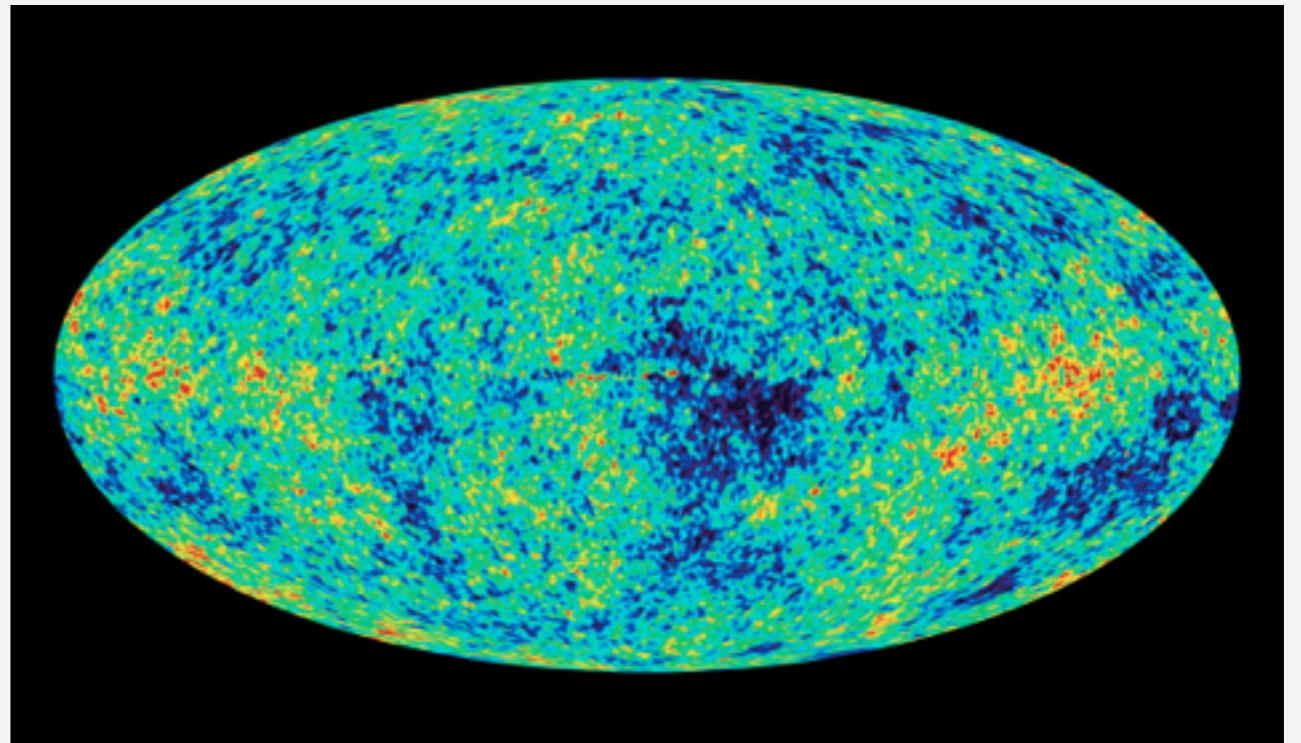
$$\left(\frac{\delta T}{T}\right)_* \equiv \frac{\delta\gamma}{4} + \psi$$



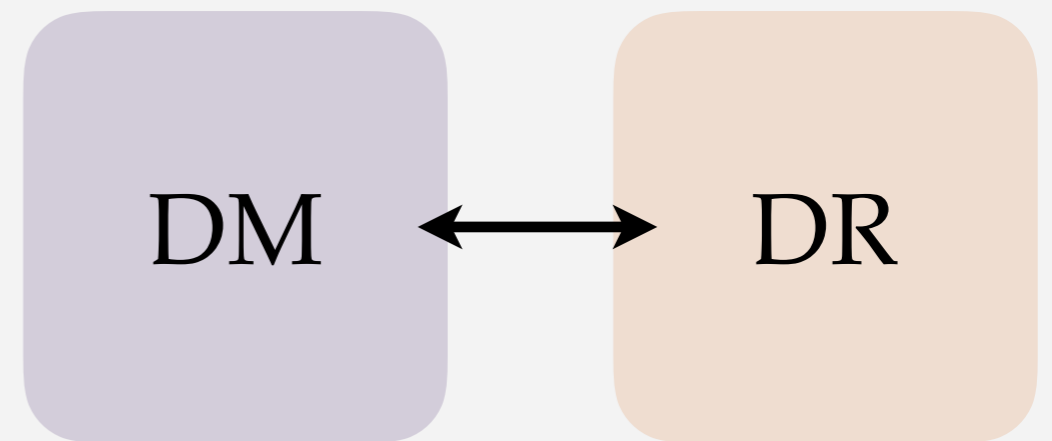
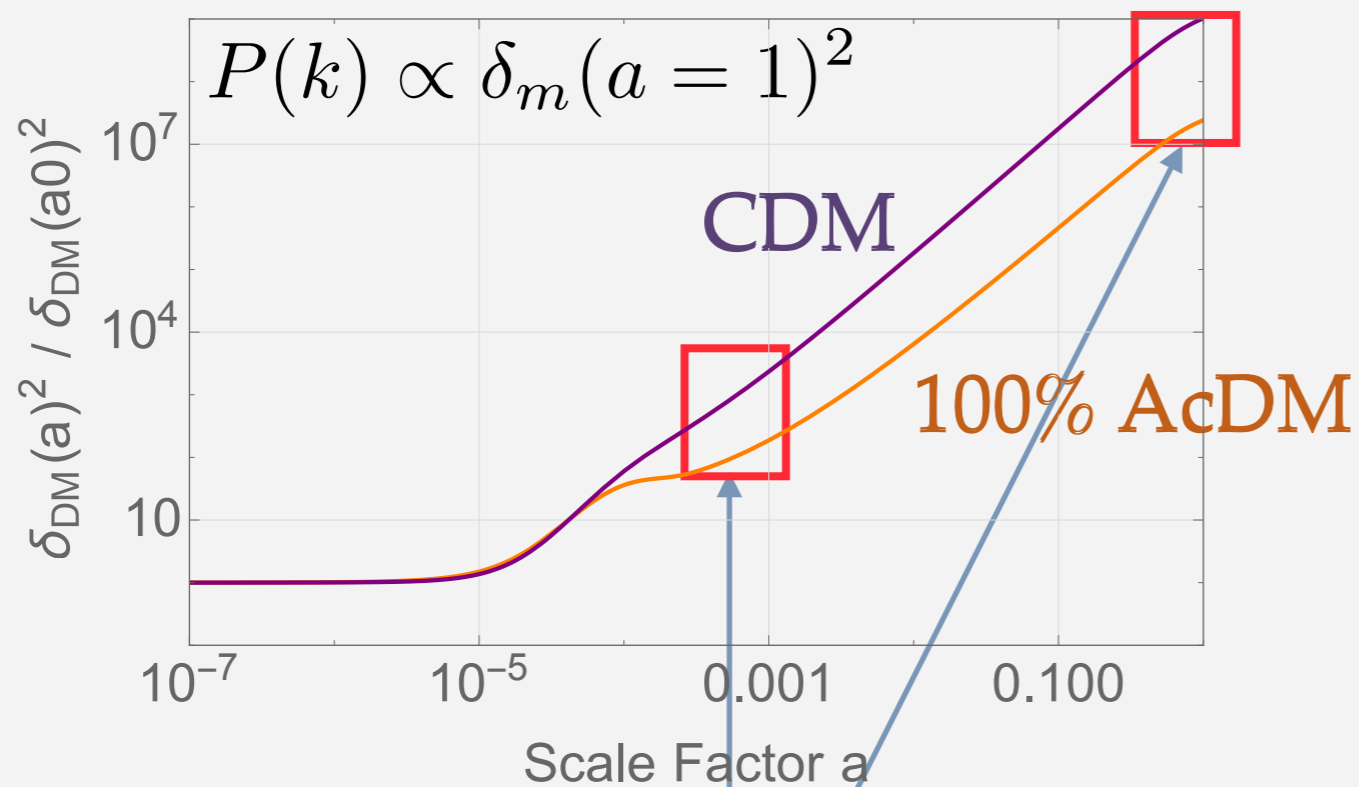
The pressure from dark fluid suppresses the compression peaks and enhances the expansion peaks

When $r = 2\%$, the correction to CMB is less than $\sim 2\%$, smaller than $> 5\%$ error bar in Planck result

Why is the CMB correction so small
in the Partially Acoustic DM case?

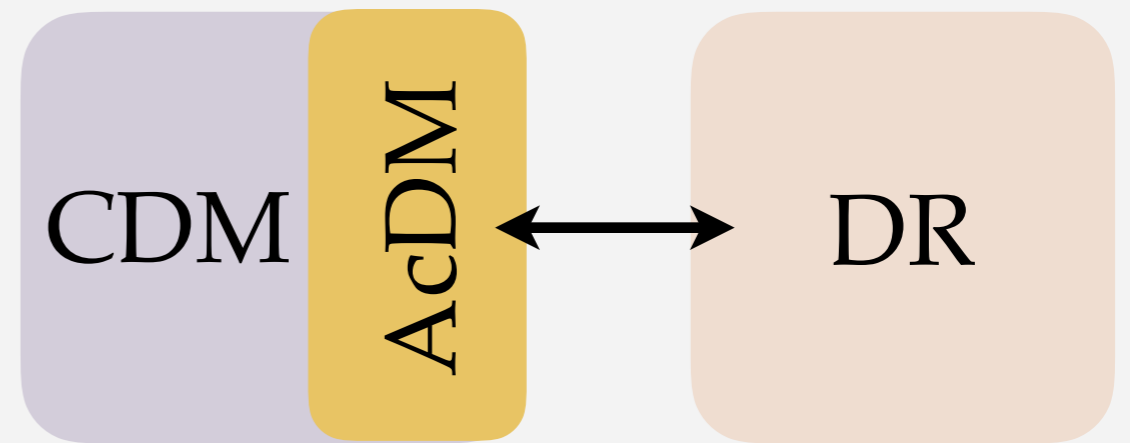
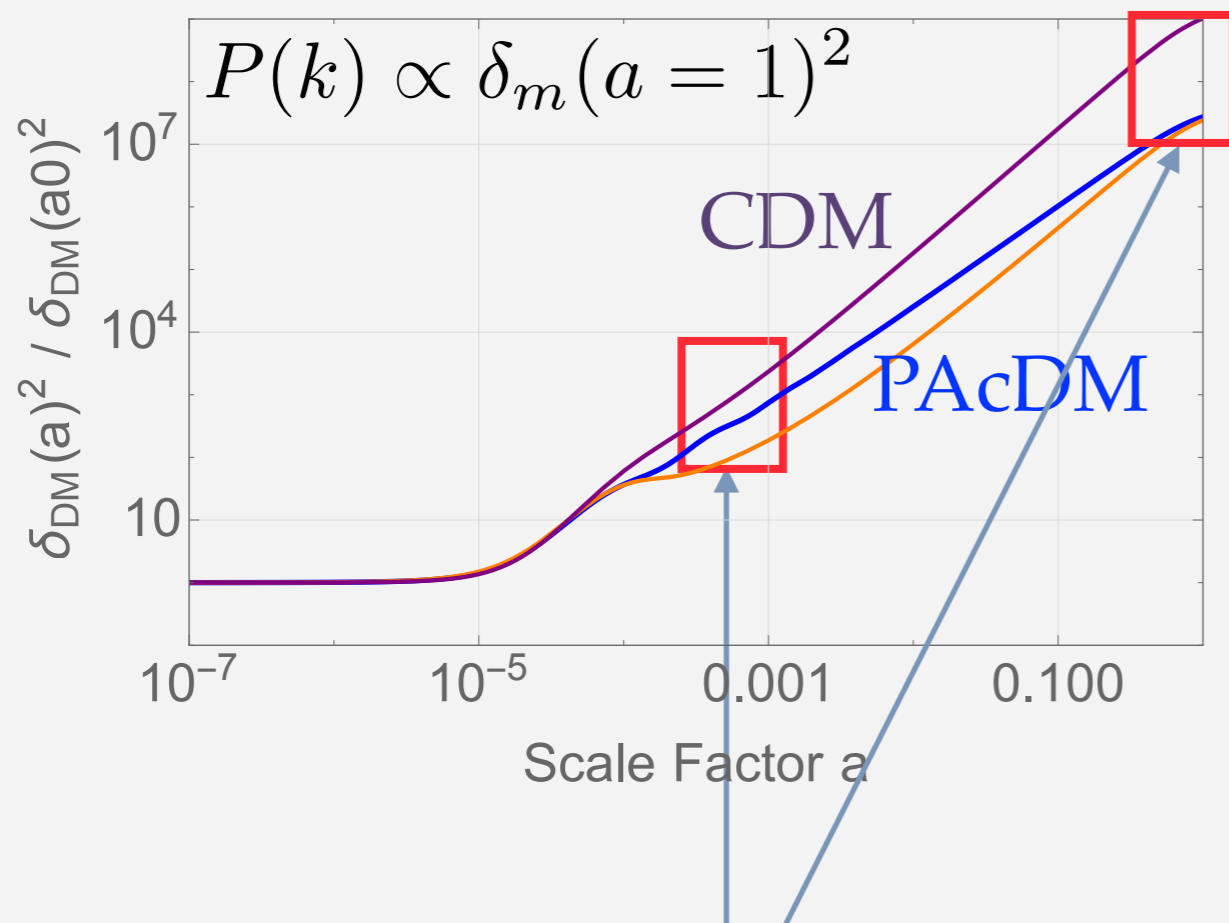


In the **fully** acoustic oscillation case



Similar damping between today / CMB time

In the **partially** acoustic oscillation case



$$\frac{P(r)}{P(0)} \simeq (1 - 2r) \left(\frac{a}{a_{eq}} \right)^{-1.2r}$$

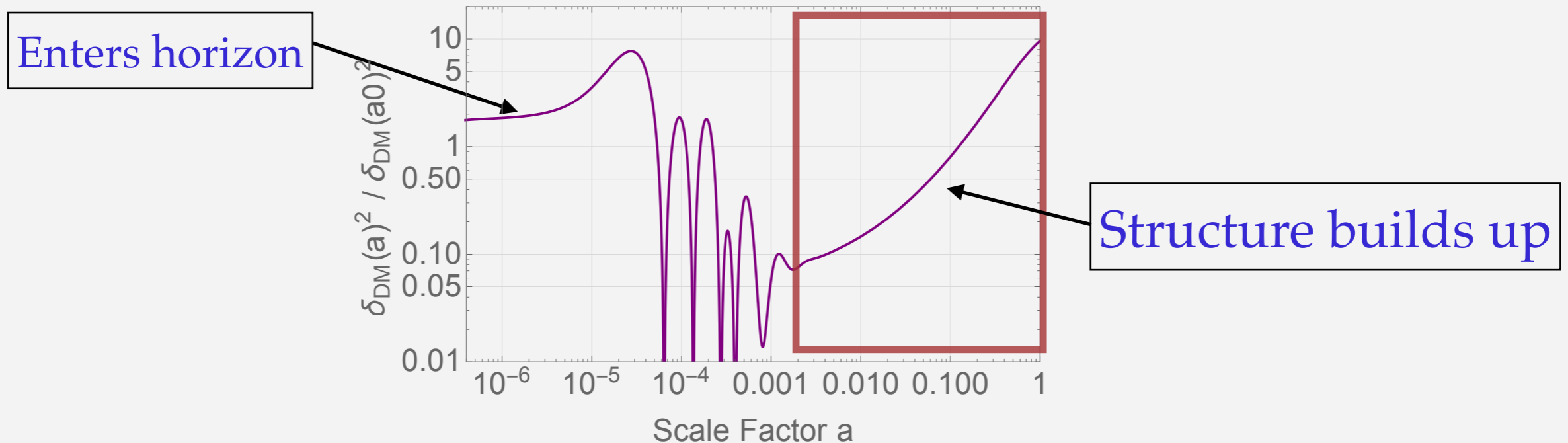
Structure grows **slower** comparing to CDM
 Smaller correction to the CMB spectrum

Matter Domination , $R \gg 1$

No oscillation => Linear growth

$$\ddot{\delta}_D + \frac{2}{\tau} \dot{\delta}_D = -k^2 \phi$$

No oscillation, no damping from the DR scattering
Same structure formation as cold DM



Matter Domination , $R \gg 1$

No oscillation => Linear growth

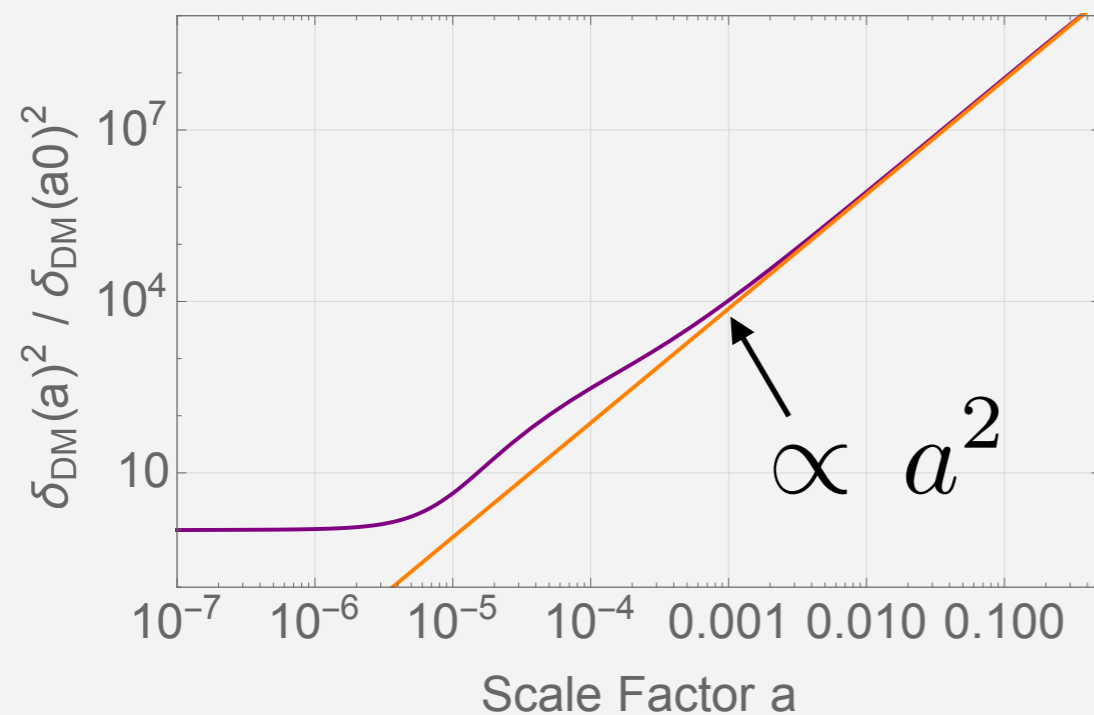
$$\ddot{\delta}_D + \frac{2}{\tau} \dot{\delta}_D = -k^2 \phi$$

$$\frac{k^2 \phi = -4\pi G a^2 \sum_s \delta_s \rho_s \simeq -4\pi G a^2 \delta_D \rho_D}{\text{Einstein equation}} = -\frac{6}{\tau} \delta_D$$

($\nabla^2 \phi = 4\pi G \rho$ Poisson's eq. for gravity)

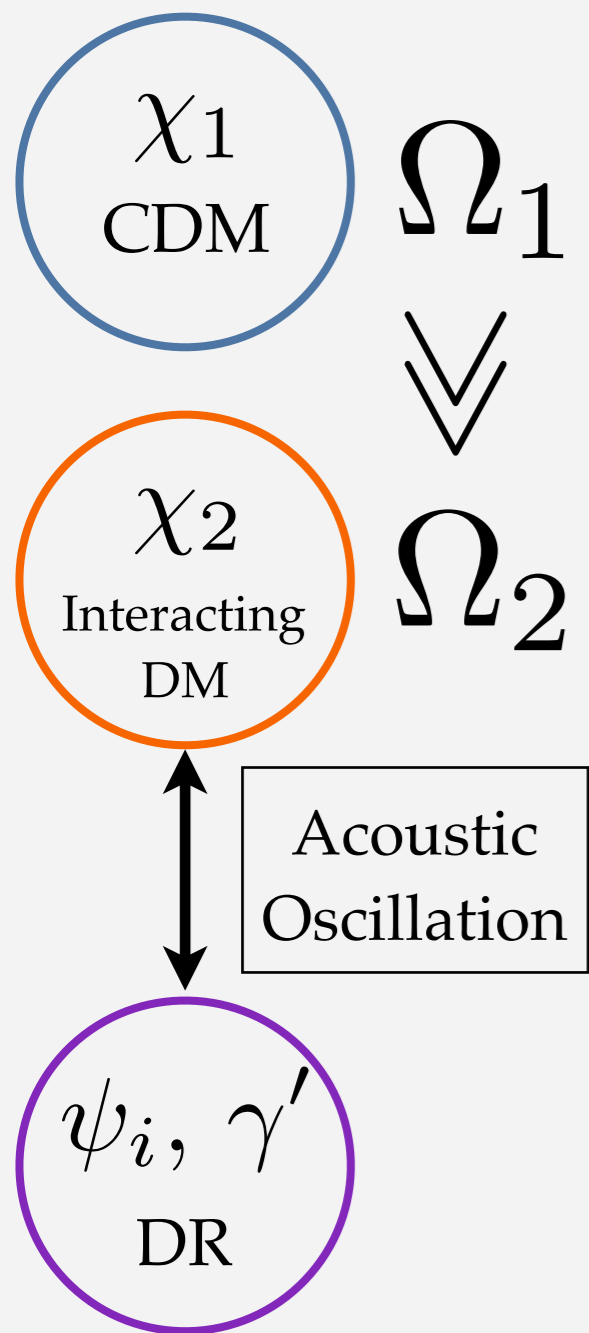
Linear growth of CDM

$$\ddot{\delta}_D + \frac{2}{\tau} \dot{\delta}_D = \frac{6}{\tau} \delta_D \quad \Rightarrow \quad \delta_D \propto \left(\frac{\tau}{\tau_{eq}} \right)^2 = \left(\frac{a}{a_{eq}} \right)^1$$



Density contrast in Cold DM case grows **linearly** in the deep matter-dominated era

In the partially acoustic case



Acoustic
Oscillation

$$\Rightarrow \delta_1 \gg \delta_2$$

DM density contrast is determined by χ_1

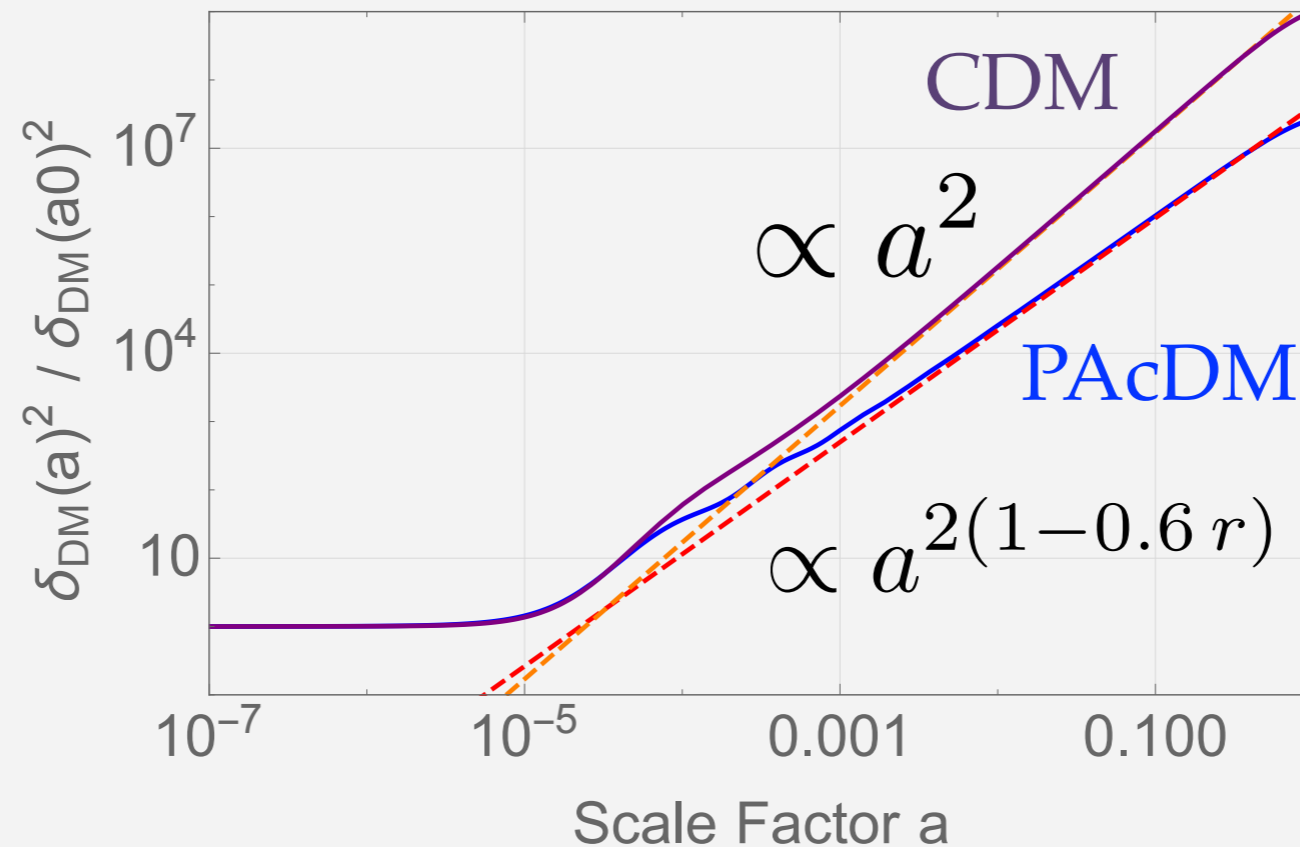
$$\ddot{\delta}_1 + \frac{2}{\tau} \dot{\delta}_1 = -k^2 \phi$$

$$k^2 \phi \simeq -4\pi G a^2 (\delta_1 \rho_1 + \delta_2 \rho_2)$$

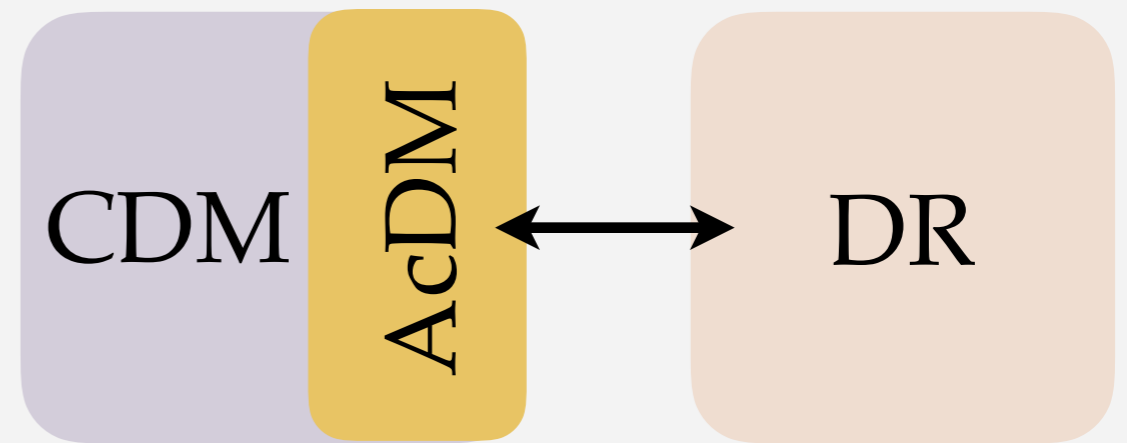
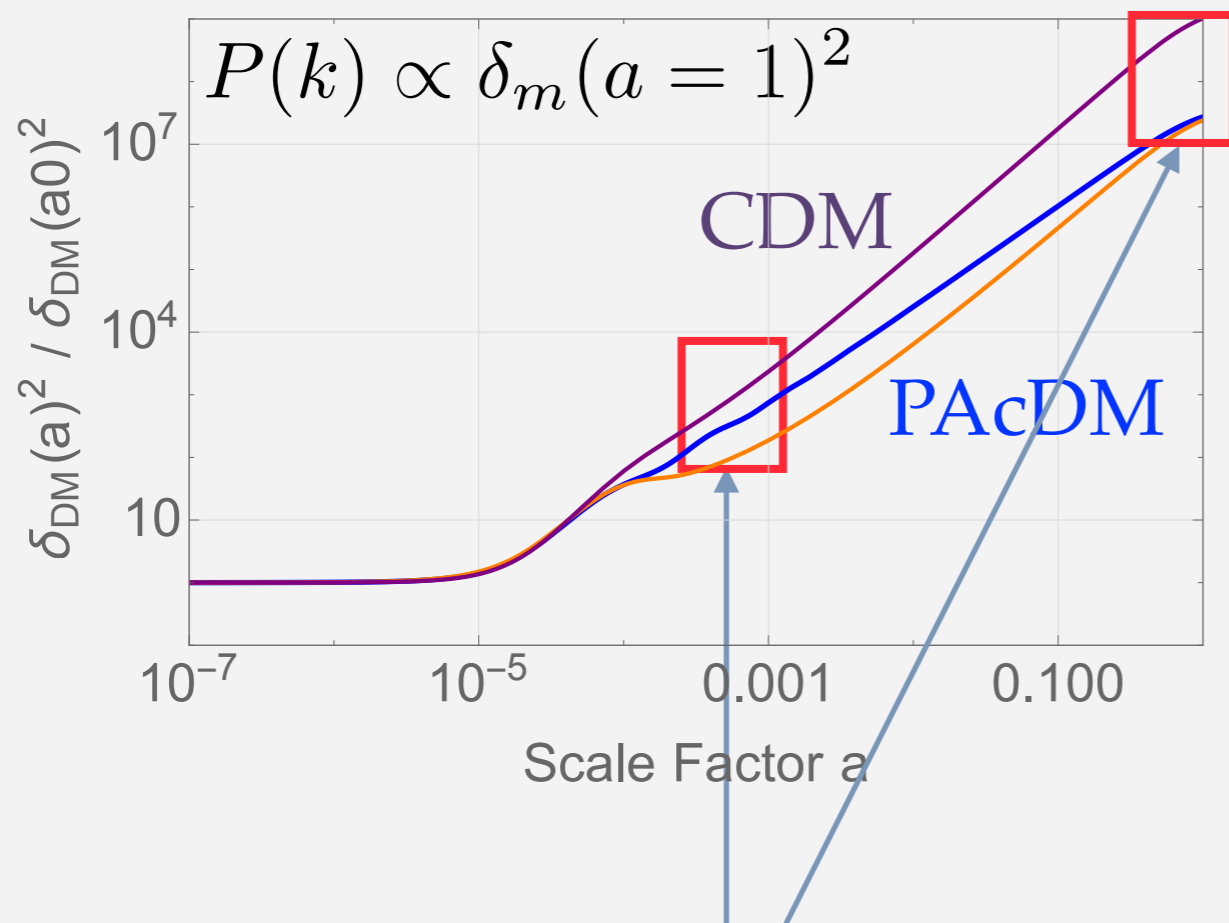
$$= -\frac{6}{\tau} (1 - r) \delta_1 \quad r \equiv \frac{\rho_2}{\rho_{DM}}$$

The < 100% CDM case

$$\ddot{\delta}_1 + \frac{2}{\tau} \dot{\delta}_1 = \frac{6}{\tau} (1 - r) \delta_1 \quad \Rightarrow \quad \delta_1 \propto \left(\frac{a}{a_{eq}} \right)^{1 - 0.6 r + \mathcal{O}(r^2)}$$



How about only a fraction of DM particles having the acoustic oscillation?



$$\frac{P(r)}{P(0)} \simeq (1 - 2r) \left(\frac{a}{a_{\text{eq}}} \right)^{-1.2r}$$

Structure grows **slower** comparing to CDM
 Smaller correction to the CMB spectrum

Conclusion

Large Scale Structure is sensitive to the dark sector dynamics

A smaller ratio of Cold DM

change the power-law growth of matter density spectrum

Acoustic Dark Oscillation

suppresses the matter power spectrum

Having Dark Radiation

change the expansion, different effects on CMB between
free-streaming / self-scattering

May also change the small scale structure

Working on it now, stay tuned!