

# (Super)Planckian Physics and Inflation

Matthew Kleban

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# What I'm **not** talking about

- Inflation expands the universe by a factor  $> \sim e^{60}$  in linear scale
- This inflates initially subPlanckian wavelengths to large, observable scales
- But the true physics of inflation is close to time-translation invariant (hence scale invariance), and high scale corrections are suppressed by powers of  $H/M < \sim 10^{-5}$

# Field ranges

- But there is another interesting superPlanckian issue
- If  $H/M > \sim \text{few} * 10^{-7}$ , the value of the inflaton field changed by a (reduced) Planck mass or more during the last  $\sim 60$  efolds of slow roll,  $\Delta\phi > \sim M$
- Such a large value of  $H$  also produces large tensor power via the usual mechanism of de Sitter quantum fluctuations
- If BICEP2 measured primordial gravity waves, the tensor/scalar power ratio  $r > \sim .1$
- Standard dS fluctuations give  $r \sim (10^5 H/M)^2$ , so  $r \sim .1$  implies  $\Delta\phi \sim 10 * M$
- Does this constitute superPlanckian physics, with large (and very interesting and very hard to compute) quantum gravity corrections?

Lyth, Turner

# Possible loophole?

- Before addressing that, one way out is if gravity waves can be generated in some non-standard way during inflation, so large  $r$  does not imply large  $H/M$
- Extra tensors can be produced, for instance, by coupling the inflaton to other fields, in such a way that while it rolls it produces stuff that then decays to or emits gravitons
- However, there is a no-go theorem being developed, the result of which is that any such mechanism produces more scalars than tensors, by a factor of  $\epsilon^{-2}$ , so that  $r \sim \epsilon^2$
- There are models that seem to evade this, but they are not beautiful (additional rolling scalar, many other fields)

Mirbabayi Senatore Silverstein Zaldarriaga

# SuperPlanckian vevs

- So,  $r > \sim .001$  probably implies a  $O(1)$  change in the inflaton vev in Planck units during inflation (and we will know if this is true in the near future)
- Of course, gravity doesn't couple (directly) to the vev, it couples to stress-energy
- So long as the energy density is below Planck, the gravitational coupling is weak

# Effective potential

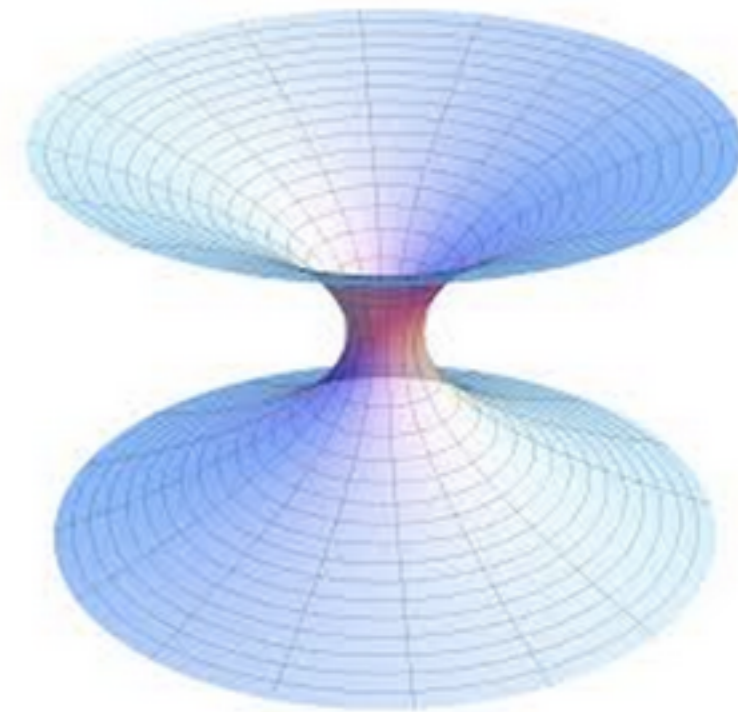
- Indeed, one can compute the effects of graviton loops on the effective potential for the inflaton  
Coleman-Weinberg, Smolin, Linde
- They produce corrections like  $V(1 + V/M^4 + V''/M^2 + \dots)$
- These are never large during inflation, when  $V/M^4$  and  $V''/V$  are small
- So, there is no problem with superPlanckian field ranges in perturbation theory
- However there is still a potential issue

# Hierarchy

- The mass of the inflaton must remain small for slow roll, but scalar masses get large corrections from loops
- Obvious candidate to protect the mass is an (approximate) shift symmetry  $\phi \rightarrow \phi + C$
- But we know that gravity breaks global symmetries (e.g. Hawking evaporation of black holes)
- So if not forbidden by symmetry, shouldn't we write all operators like  $\sum_p \phi^{p+4}/M^p$  with  $O(1)$  coefficients?
- That really would destroy inflation, because all terms become larger than  $\phi^4$  when  $\phi > M$  - so there are/were claims that large  $r$  is inconsistent, or at least incompatible with an EFT description of inflation

# Wormholes!

- Assuming they are absent in the classical theory, such terms must be generated non-perturbatively
- The known NP effects in gravity do indeed produce terms like that, but not with  $O(1)$  coefficients - instead they are multiplied by  $e^{-S}$ , where  $S$  is the action for some instanton that eats the global charge (a wormhole)
- Turns out that  $S$  (for wormholes) is extremely sensitive to threshold corrections at or below the Planck scale, as well as various curvature corrections to Einstein





# For example...

- Adding a Gauss-Bonnet term ( $R_{\dots}^2 - 4R_{..}^2 + R^2$ ) doesn't change Einstein's equations, because it's topological
- But it does change  $S$ , because a wormhole that eats global charge has a different topology than flat space
- In this way one can make  $S$  arbitrarily large without affecting experiments
- Another way is to have compact extra dimensions, where the wormhole throat never gets smaller than the extra dimension, making its action large, or by adding other higher curvature corrections

# Weakest force

- Another possible worry is that, at least in some models, superPlanckian field ranges end up corresponding to a force that is weaker than gravity
- However, the “gravity is the weakest force” argument only clearly has weight for gauge forces
- Furthermore there seem to be explicit models in string theory that accomplish this

Arkani-Hamed, Motl, Nicolis, Vafa

# Summary

- So far we have considered two distinct but related issues for inflation with superPlanckian field ranges
- Perturbative corrections to the mass, which can be controlled with a global symmetry
- Non-perturbative gravity corrections to the potential, which break that symmetry - but at least wormhole corrections can be made small in several ways (of course, it's still possible some other NP effects are more important)
- So superPlanckian vevs probe some aspects of non-perturbative quantum gravity, but not very directly, and it doesn't seem difficult to make the effects very small if you are not in pure Einstein

# Models?

- In fact we have various models that produce superPlanckian vevs without any apparent issue

- Monodromy inflation

Silverstein, Westphal, McAllister, Flauger...

- Unwinding inflation

D'Amico, Gobbetti, MK, Schillo

- Extranatural inflation

Arkani-Hamed, Cheng, Creminelli, Randall

- 4D effective models of Kaloper/Lawrence/Sorbo

# What about naturalness?

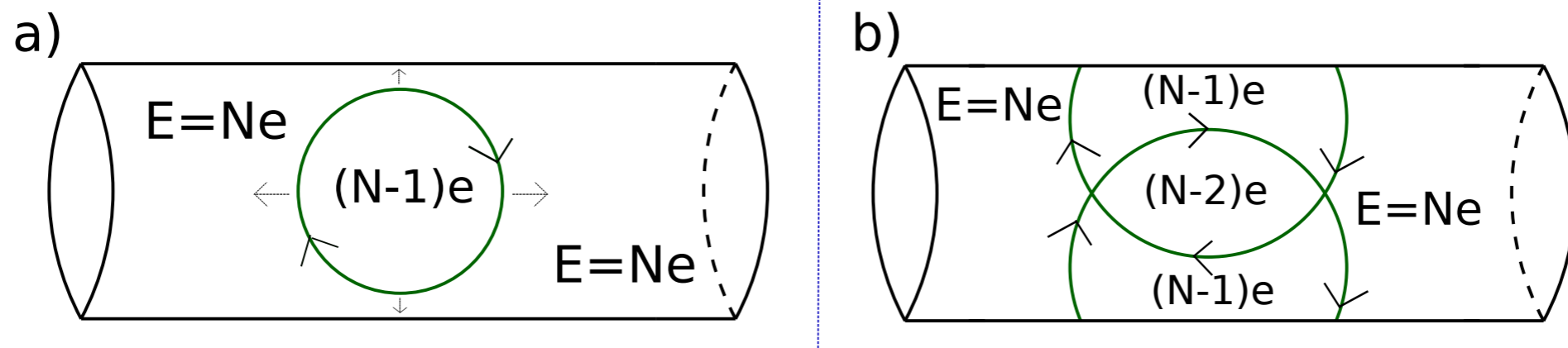
- Are large field models less natural than small field models?
- If there is a landscape, we are in a bubble that formed in a first-order phase transition from a metastable parent phase, and then underwent slow roll inflation of some sort
- These transitions can generically initiate unwinding inflation (“creation myth”)
- Because the inflationary vacuum energy is that of the parent phase, it is naturally large, hence large field range
- Small field inflation might require much more tuning

# Unwinding

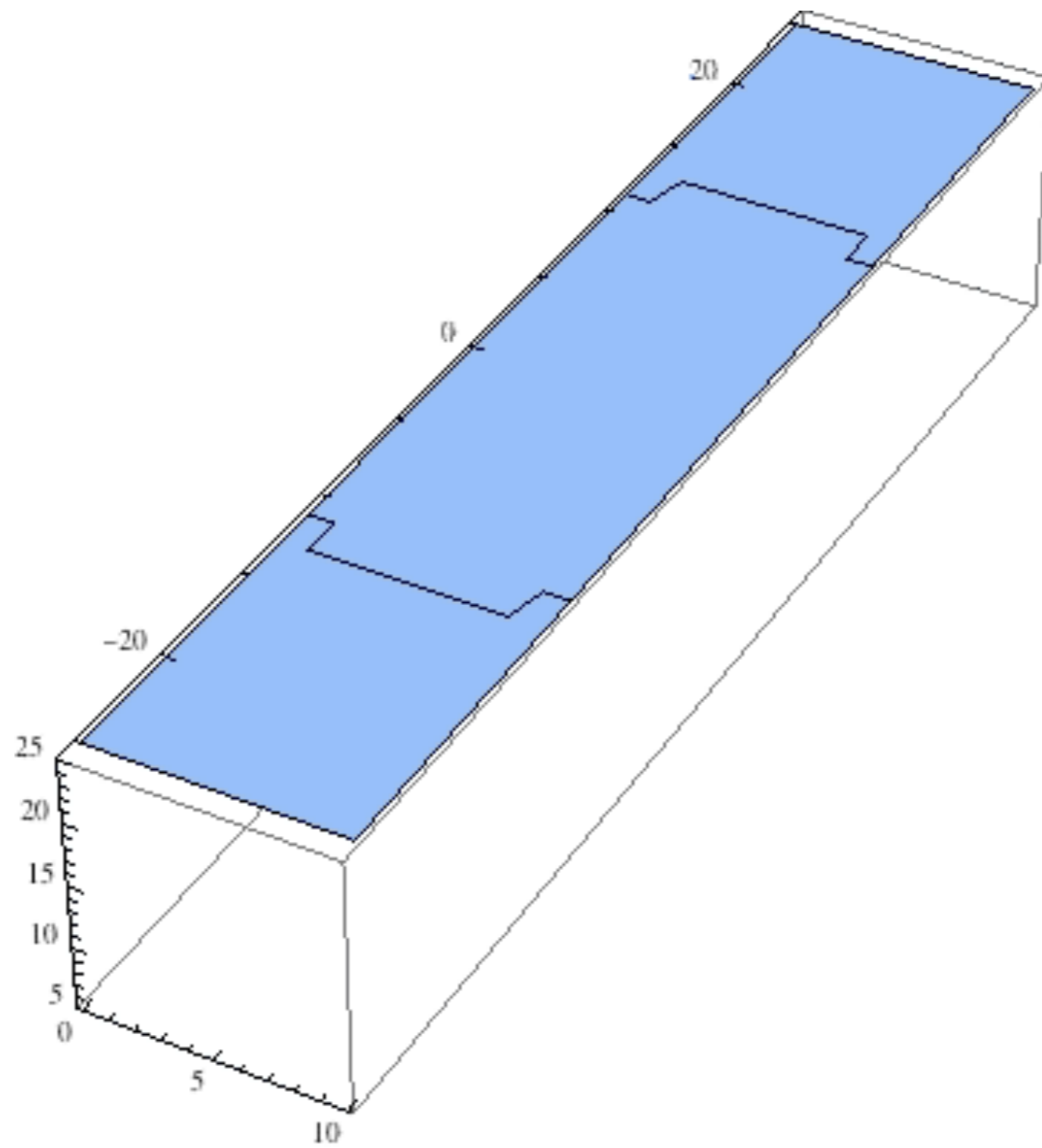
- Consider a spacetime of the form  $dS_4 \times M$ , where all moduli of  $M$  are stabilized (these had better exist, or we are in trouble!)
- There are several possible contributions to the vacuum energy of the  $dS$ , one of which is a flux  $F_p$  with  $p > 4$  that fills the  $dS$  and threads  $M$
- Any spacetime like this is at best metastable, and one decay mode is to discharge one unit of  $F$  flux ( $F$  is quantized) via nucleation of a bubble of charged brane (higher form flux analog of Schwinger pair production)

# Flux cascade

- This can initiate a cascade as the bubble expands around the compact directions - many units of flux discharge



- Result is a homogeneous and isotropic **open** FRW cosmology dominated by gradually decreasing vacuum energy - slow roll (open) inflation, where the inflaton is the radius of the bubble
- In a sense this is a version of old inflation, since it is the false vacuum energy density that dominates during inflation





# Old inflation rejuvenated

- Inflation ends when the flux is discharged and the brane annihilates with itself, reheating the universe
- $\Delta\phi=M$  (for the canonical 4D inflaton  $\phi$ ) corresponds to the discharge of a few units of flux (a few wraps around)
- Doesn't necessarily disturb the stabilization mechanism much, although it can and lead to flattening a la monodromy
- Most generic prediction (at current understanding) is that there is either detectable equilateral NG, or tensors, or both

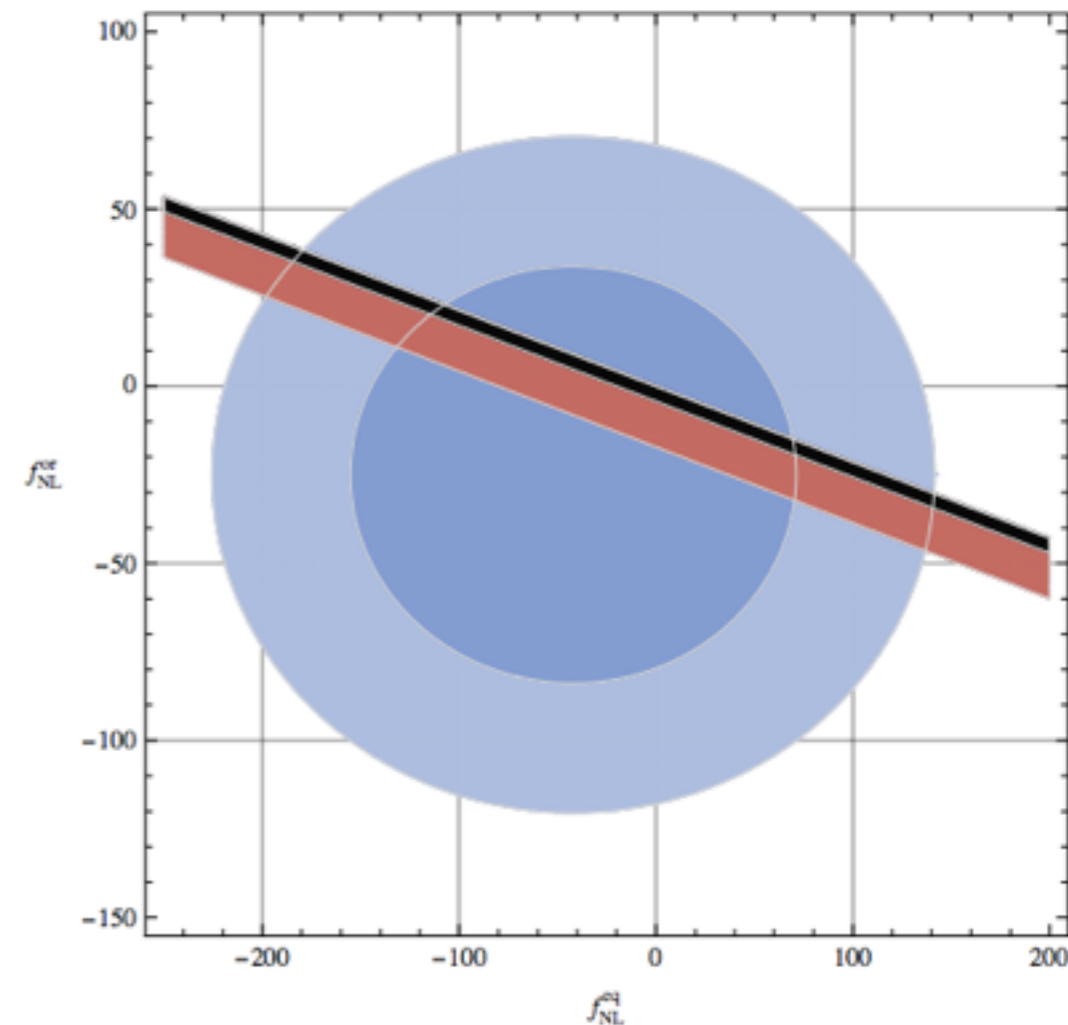
Brown

# Reheating?

- An interesting situation arises near the end of inflation
- The branes may “prematurely” annihilate in a Hubble-sized region before discharging all units of flux
- This possibility is realized by a random distribution of regions with different values of the flux at reheating
- I think these regions will collide and reduce the flux to zero everywhere, but there will be large perturbations on Hubble scales, and maybe gravitational waves

# Does large $r$ imply Gaussianity?

- Naively, large  $r$  makes detectable NG unlikely, because the easiest way to achieve large NG is to make  $c_s$  small ( $f_{\text{NL}} \sim c_s^{-2}$ ), but this suppresses tensors ( $r = 16\epsilon c_s$ )
- But there is an operator  $(\dot{\pi})^3$  in the cubic effective theory, and its coefficient is not constrained by large  $r$  and does not generate small  $c_s$
- This operator predicts a very specific shape for NG, which can be large



# High scales

- Anomalies in the CMB data are enhanced by large  $r$ , and hint at short inflation (Cora's talk)
- Various relics from the pre-inflationary state might be detectable
- I think **spatial curvature** should be our next target
  - if positive, it falsifies the landscape (or at least a parent vacuum) and slow roll eternal inflation in our immediate past
  - if negative, consistent with birth by tunneling, and rules out SREI
  - Only remaining very-large-scale observable that is far from cosmic variance limit

# What's next?

- If large  $r$  is confirmed, it's a great situation for both theorists and observers
- SuperPlanckian field range is possible and controllable, but “on the edge of respectability”
- Non-perturbative quantum gravity effects are potentially detectable
- Very useful hint